

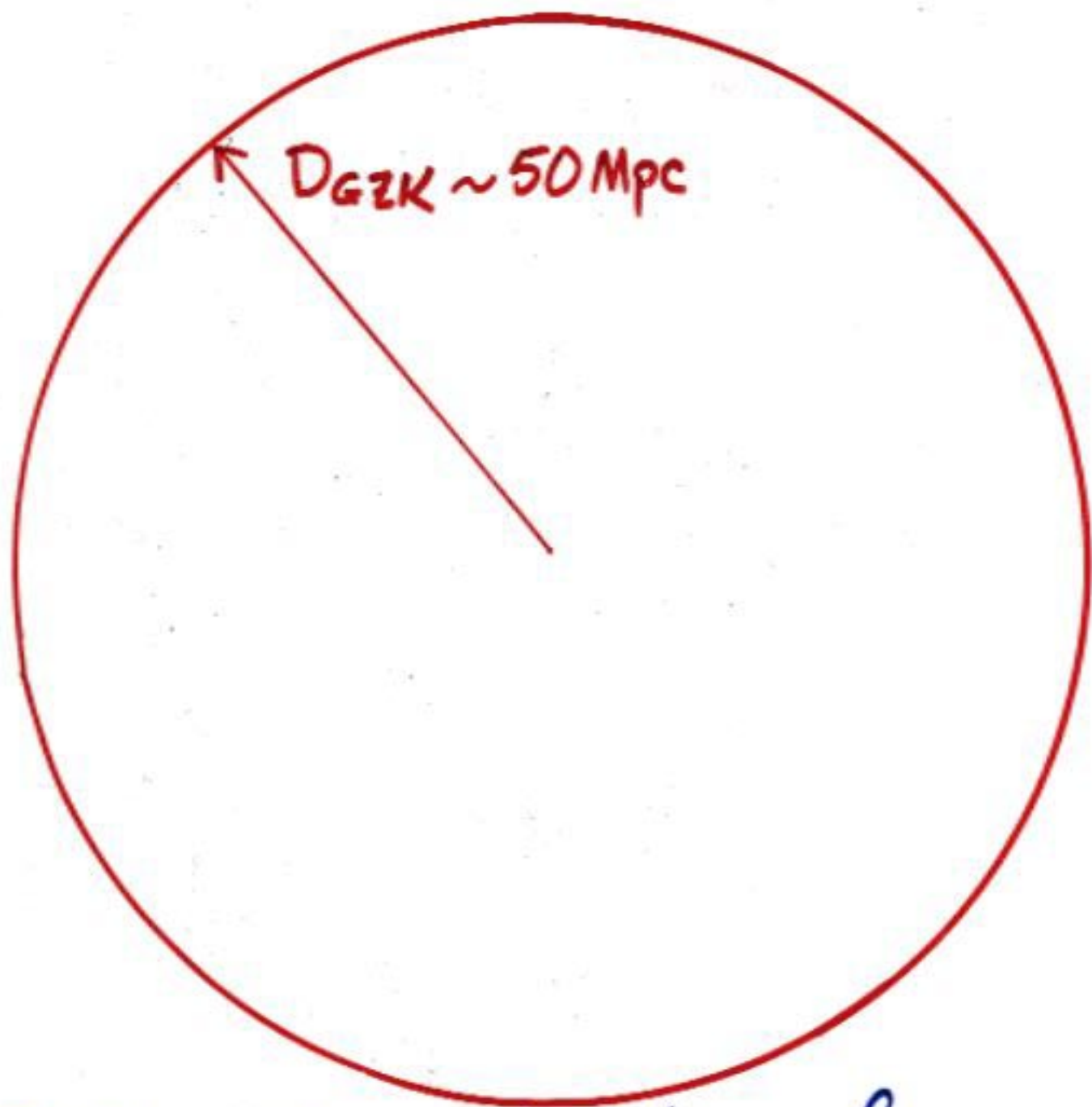
Cosmic Ray

NEUTRINOS

above  $10^{20}$  eV

[ and below  $10^{-3}$  eV ]

Tom Weiler

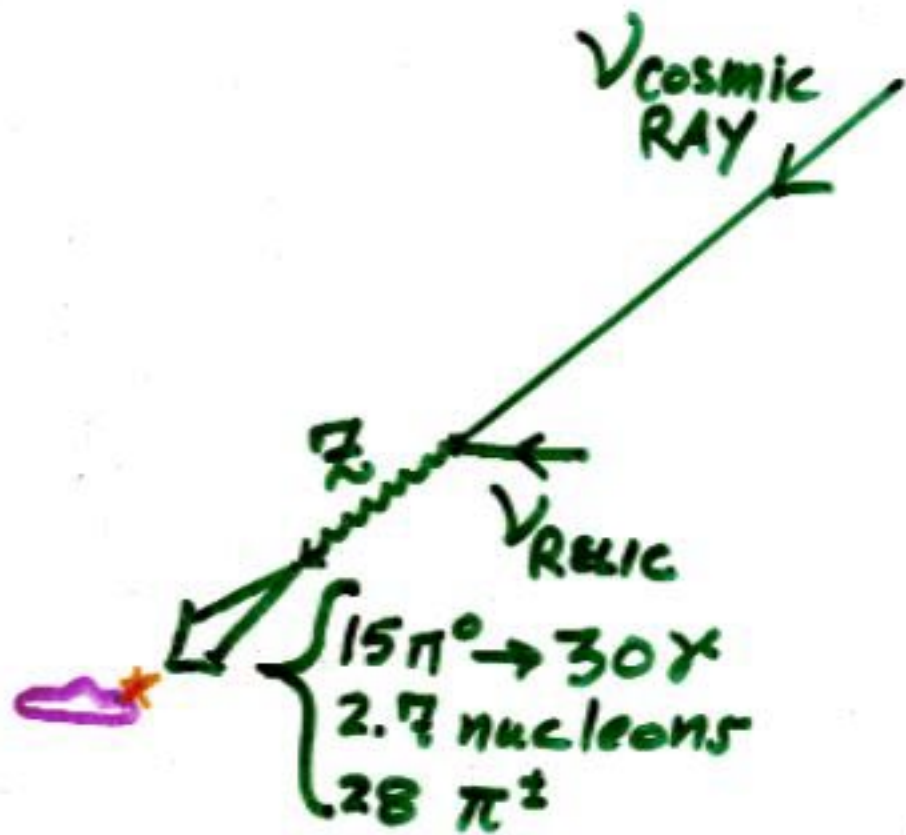


Find ~1% probability for  
resonant  $\nu \rightarrow \pi$ -burst within D<sub>GZK</sub>

TW, hep-ph/9710431

Astropart. Phys.

Also, Farqion, Mde, Salis  
astro-ph/9710029



# OUTLINE

## THE CVB

- WHAT IS IT? (overview)
- ATTEMPTS at DETECTION
- THE  $\nu_{CR} + \nu_{relic}$  annihilation hypothesis, for
  - Neutrino Spectroscopy
  - Super-GZK Cosmic Rays <sup>DATA</sup>
- More on the CVB (details)
- The future - Cosmic Ray viewpoint

# COSMIC $\nu$ BKG D

\* where particle physics/cosmology merge

Analogous to 2.73K CMB,

liberated at  $t \sim 500,000$  yrs

$$T \sim \text{eV}$$

$$z \sim 1100$$

[Planck]

when  $e^- + p^+$  recombined  $\rightarrow$  hydrogen  
(plasma) (transparent)



$\exists$  (Xpt'l challenge) C $\nu$ B

liberated at  $t \sim 1$  sec

$$T \sim \text{MeV}$$

when  $\Gamma_{\text{weak}} = \sigma_{\text{weak}} n$

fell below  $\Gamma_{\text{expansion}} = H$ .

Normalizing to CMB,  $T_\nu \sim 1.95\text{K}$

$$z_\nu \sim 10^{10}$$



with  $T_\nu \sim 2\text{K}$ ,

$$\langle p_\nu \rangle \sim 3T \sim 6 \cdot 10^{-4} \text{eV}$$

$\Rightarrow$  Non Relativistic if  $m_\nu \geq 6 \cdot 10^{-4} \text{eV}$

$$\text{Since } \langle E_\nu \rangle \sim \frac{p^2}{2m} \sim \frac{10^{-7} \text{eV}}{(m_\nu/\text{eV})},$$

way subthreshold for CC,  
& NC provides no particle tag,

Searches are necessarily\*  
**INDIRECT!**

$$\text{Note also, } \langle N_\nu \rangle \equiv \frac{\langle p \rangle}{m} = \frac{6 \cdot 10^{-4}}{(m_\nu/\text{eV})}.$$

# Possible Neutrino Mass Spectrum:

From oscillations,  $\delta m_{jk}^2 \equiv m_j^2 - m_k^2$

$$\Rightarrow m_j \gtrsim \sqrt{\delta m_{jk}^2}$$

$$\sqrt{\delta m_{\text{LSND}}^2 \sim 0.3 - 2 \text{eV}^2} \sim 0.5 \text{ to } 1.5 \text{eV}$$

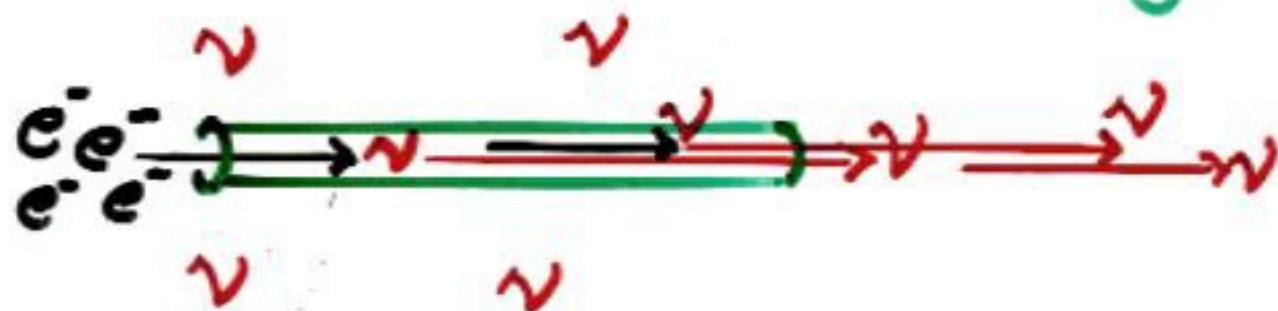
$$\sqrt{\delta m_{\text{atm}}^2 \sim 10^{-3 \pm 0.5}} \sim 10^{-1} \text{ to } 10^{-2} \text{eV}$$

$$\sqrt{\delta m_{\text{sun}}^2 \sim 10^{-5} \text{ or } 10^{-10}} \sim 3 \times 10^{-3} \text{ or } 10^{-5} \text{eV}$$

$\Rightarrow$   $\begin{array}{c} \text{---} \\ | \\ \text{---} \end{array}$  preferred over  $\begin{array}{c} | \\ \text{---} \end{array}$   
(BGG)

$\Rightarrow$  2  $\nu$ 's @  $m_\nu \gtrsim \text{eV}$ .

# Direct Detection via TeV Lineac Scattering



$$\begin{aligned}
 \text{Production Rate} &= I_{acc} \frac{L_{acc}}{\lambda_{e\nu}} \\
 &= \left(\frac{I}{\text{Amp}}\right) \left(\frac{L_{acc}}{\text{km}}\right) \left(\frac{n_\nu}{300 \text{ cm}^{-3}}\right) \left(\frac{\sigma_{e\nu}}{10^{-44} \text{ cm}^2} \sim \frac{E_e}{\text{TeV}} \frac{m_\nu}{\text{eV}}\right) \\
 &\quad \times 10^{10} / \text{yr}
 \end{aligned}$$

## ⊗ Detection Probability

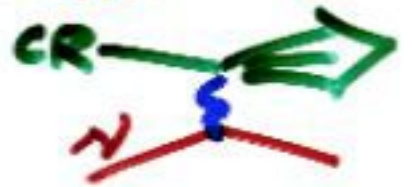
$$\theta_\nu \sim \frac{L_{det}}{\lambda_{\nu N}} \sim 10^{12} \left(\frac{\text{TeV}}{\text{eV}}\right) \left(\frac{m_\nu}{\text{eV}}\right) \quad \text{Light Pipe}$$

$$= \frac{L_{Det}}{\lambda_{\nu N}} \sim \left(\frac{L_{Det}}{\text{km}}\right) \left(\frac{n_N}{10^{26} \text{ cm}^{-3}}\right) \left(\frac{\sigma_{\nu N}}{10^{-36} \text{ cm}^2} \sim \frac{E_\nu}{\text{TeV}}\right) \cdot 10^{14}$$



# TRANSPARENCY

of the CνB



$$\lambda_{CR} = \frac{1}{n \sigma} \quad , \quad \sigma_{\text{weak}} \sim \frac{G_F^2 s}{\pi} \sim \frac{2}{\pi} G_F^2 m_\nu E_{CR}$$

$$= \frac{1}{\Omega_\nu} \left( \frac{10^{23} \text{ eV}}{E_{CR}} \right) D_H$$

with  $D_H = H_0^{-1} = 5 h_{65}^{-1} \text{ Gpc}$ .

Also,  $\sigma_{\text{weak}}$  saturates at  $\frac{G_F^2 M_W^2}{\pi}$ ,  
(mod ln's)

$$\Rightarrow \lambda_{CR} \gtrsim 10^4 \left( \frac{50 \text{ cm}^{-3}}{m_\nu} \right) D_H$$

But



is **RESONANT**

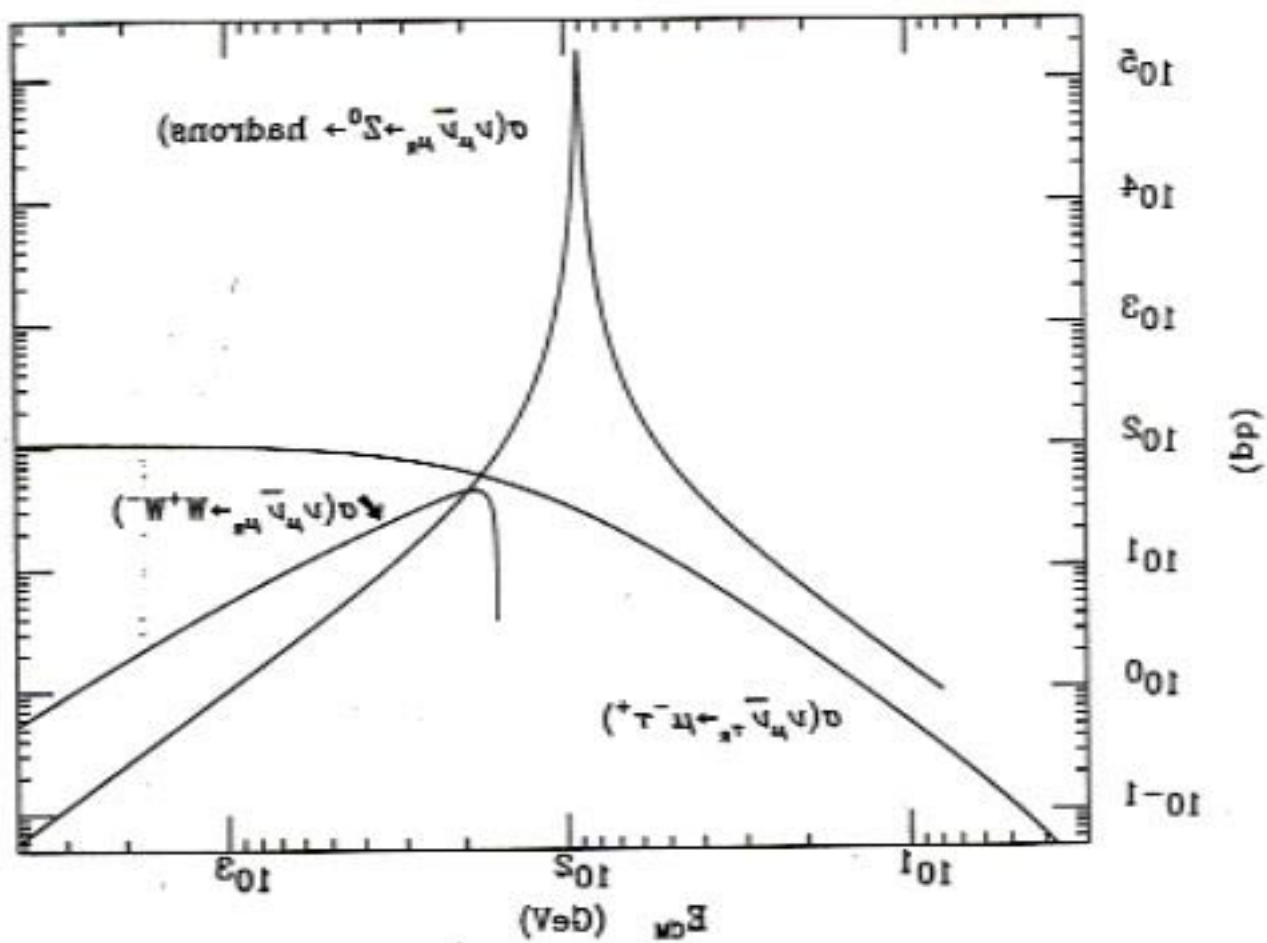
$$\langle \sigma \rangle \equiv \int \frac{ds}{M_Z^2} \sigma(s) = 2\sqrt{2} \pi G_F$$



$$\lambda_{\nu\nu} = 30 \left( \frac{50 \text{ cm}^{-3}}{n_\nu} \right) D_H$$

$$\therefore P(\nu_{\text{relic}} \bar{\nu}_{\text{relic}} \text{-annihilate}) dx$$

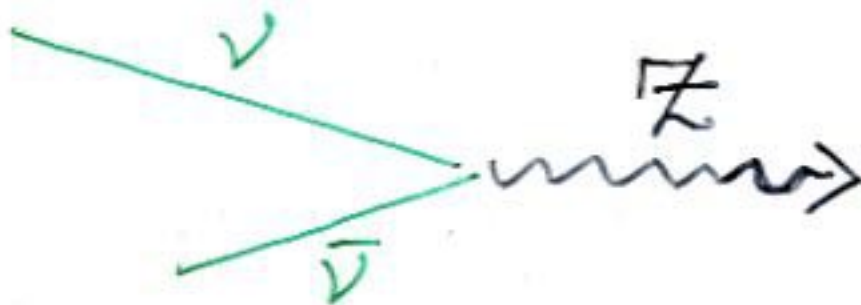
$$= \frac{dx}{\lambda} = 3\% \left( \frac{n_\nu}{50 \text{ cm}^{-3}} \right) \frac{dx}{D_H}$$



# Relic Neutrino (CνB) Spectroscopy

Imagine Cosmic Ray  $\nu$  annihilation  
on CνB:

at the  $Z$ -resonance!

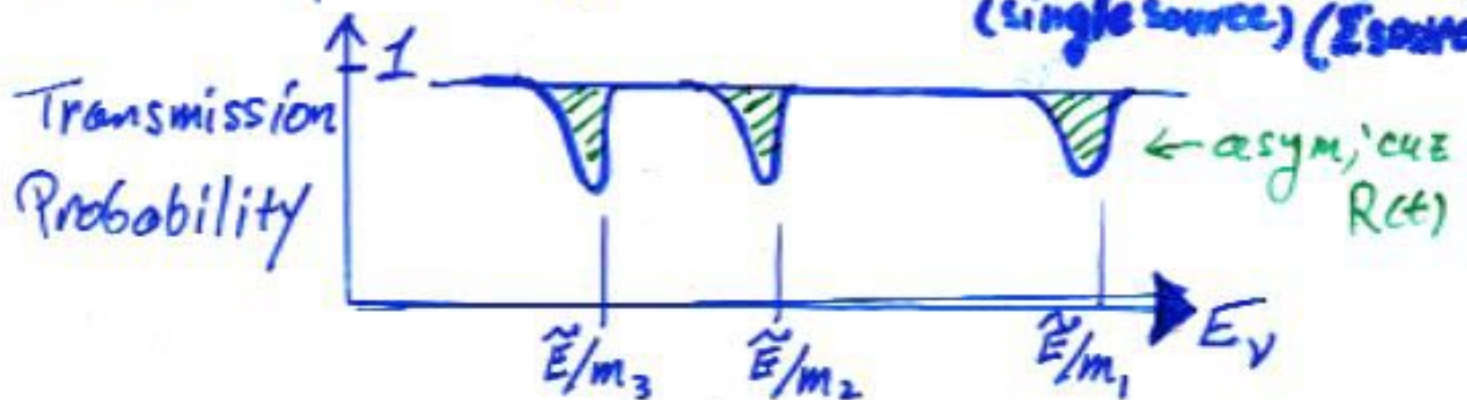


$$s = \begin{cases} M_Z^2 & \text{in CMS} \\ 2E_\nu m_\nu & \text{in LAB} \end{cases}$$

$$\Rightarrow E_{\text{res}} = \frac{M_Z^2}{2m_\nu} = \frac{4 \times 10^{21} \text{ eV}}{(m_\nu/\text{eV})} \equiv \frac{\tilde{E}}{(m_\nu/\text{eV})};$$

$$\Rightarrow \frac{\Delta E}{E_R} = \frac{2\Gamma_Z}{M_Z} \sim 6\%;$$

$\Rightarrow$  Absorption Spectrum (TW, 1982-84)  
PRL Ap. J.  
(single source) (Isotopes)



For CR-V's from AGN &  
GRB

$$z_{\text{source}} \gtrsim 1,$$

$\Rightarrow$

★  $\frac{dx}{D_H} \sim 1$

★  $n_V(z) = (1+z)^3 n_V(\text{today})$

★ Absorption Range is  $\frac{E_R}{1+z} \leq E_V \leq E_R$

$\Rightarrow P_{\text{annihilate}} \sim 3 \left(\frac{n_V}{50}\right) \% \times \left\{ \begin{array}{l} \text{Cosmological} \\ \text{Enhancement} \end{array} \right\}$

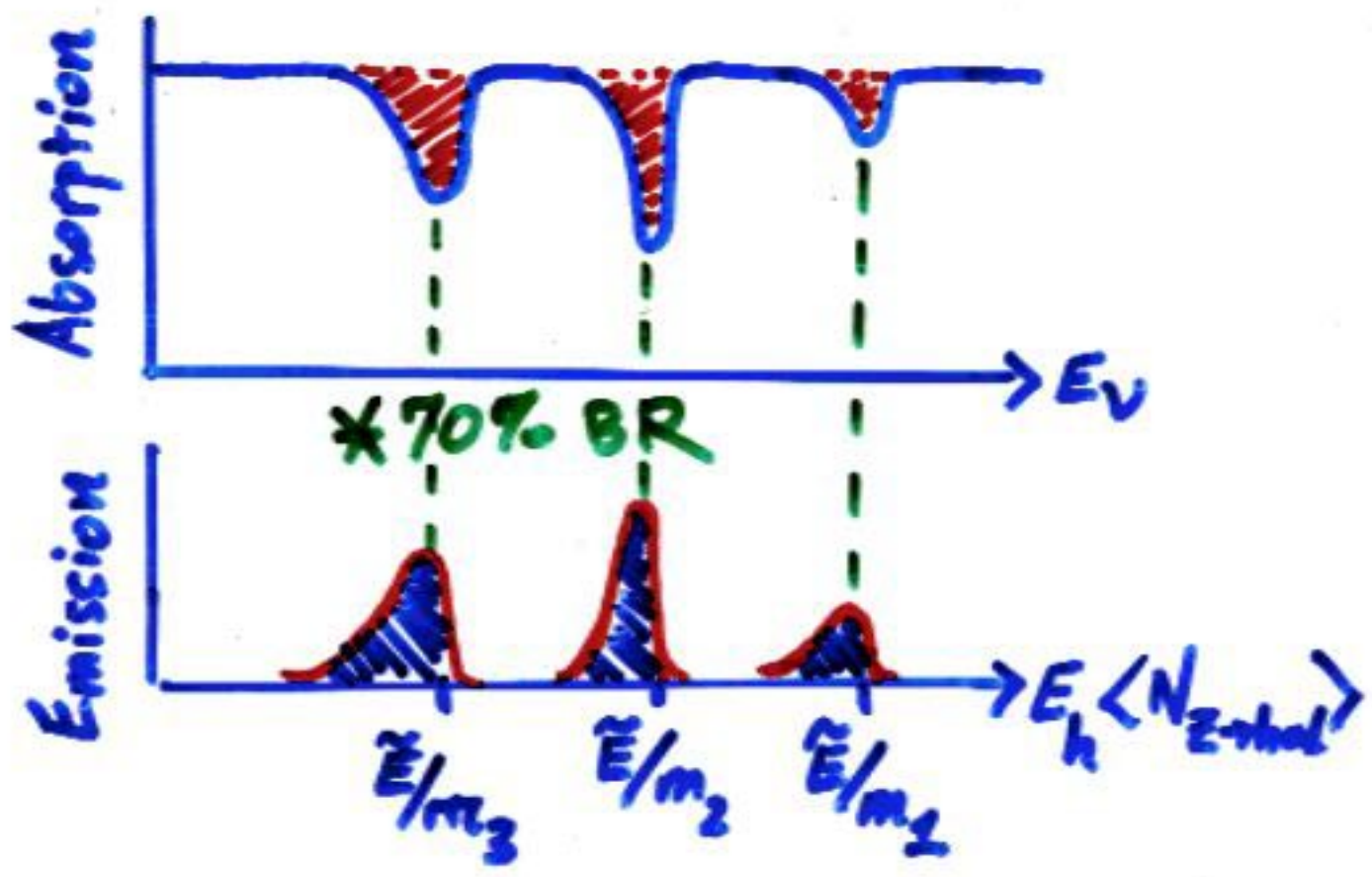
$\rightarrow \underline{10 \text{ to } 40\%} \left(\frac{n_V}{50}\right) \times |W_j|^2, \text{ for } z \sim 3.$   
depends on  $\Omega_b$ ; can be  $\mathcal{O}(5)$ !

Recent Idea { Oct '97; T.J.W.;  
Fargion, Mele, Salis

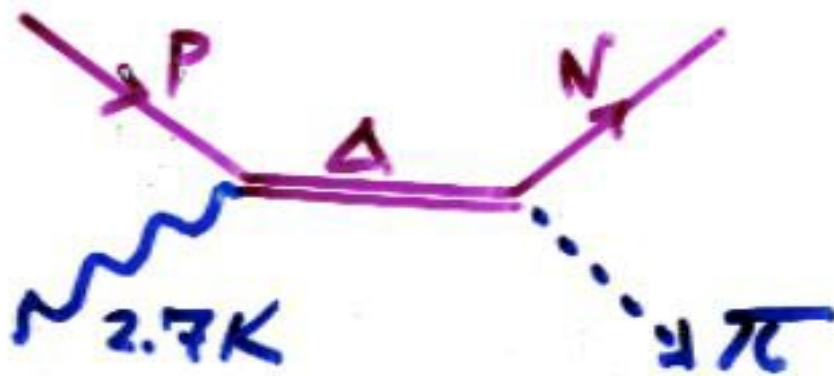
Instead of "ABSORPTION SPECTROSCOPY"  
[need hi stats]

Do "EMISSION SPECTROSCOPY"  
★ Low statistics fine  $\xrightarrow{70\%}$  {3 nucleons  
[30%]}

★ Zero Bkgd above GZK cutoff.

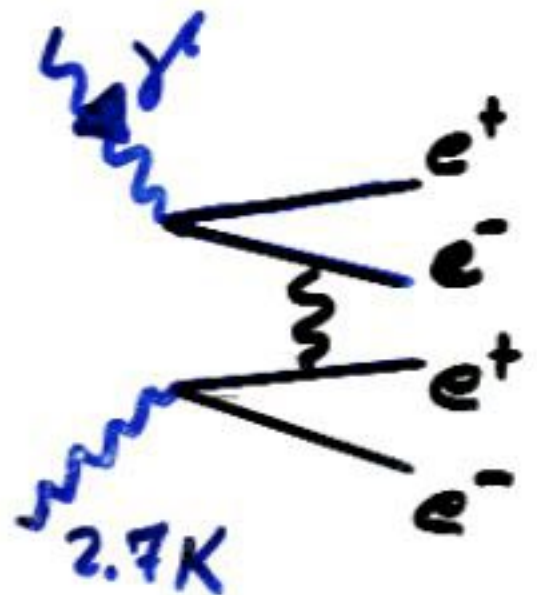
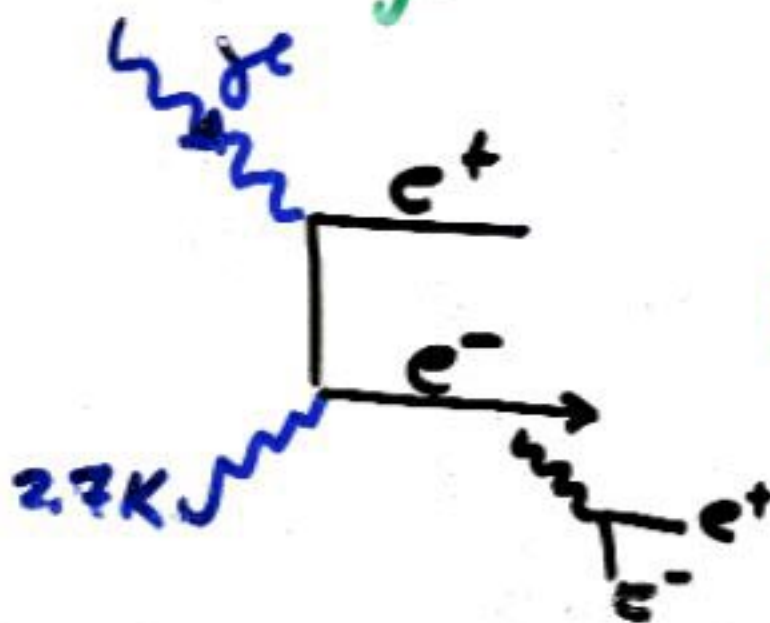


# The GREISEN-ZATSEPIN-KUZMIN (GEZK) CUTOFF



Costs  $\frac{\Delta E}{E} \sim 20\%$  per 6 Mpc.

Similarly,



## AGASA (Japan)

$E > 10^{19} \text{ eV}$	581	
$> 4 \cdot 10^{19}$	47	$D_{50\%} = 130 \text{ Mpc}$
$> 10^{20}$	6	19 Mpc

## World Sample:

$> 10^{20}$	13*	19 Mpc
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\* including John Linsley/Volcano Ranch  
1963



AGASA, PRL 81, 1163 (98)

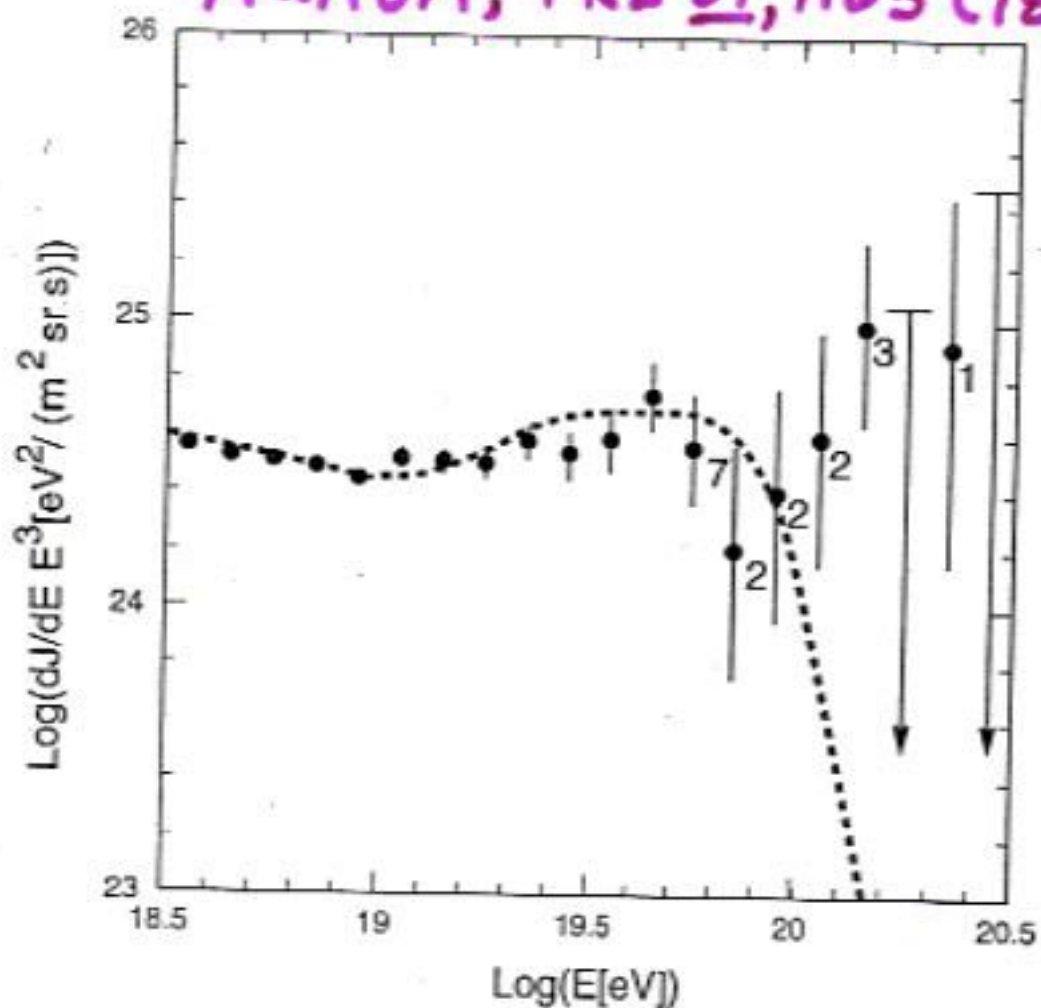


FIG. 2. Energy spectrum observed with AGASA. The vertical axis is multiplied by  $E^3$ . Error bars represent the Poisson upper and lower limits at 68% and arrows are 90% C.I. upper limits. Numbers attached to points show the number of events in each energy bin. The dashed curve represents the spectrum expected for extragalactic sources distributed uniformly in the Universe, taking account of the energy determination error [11].

## FIGURES

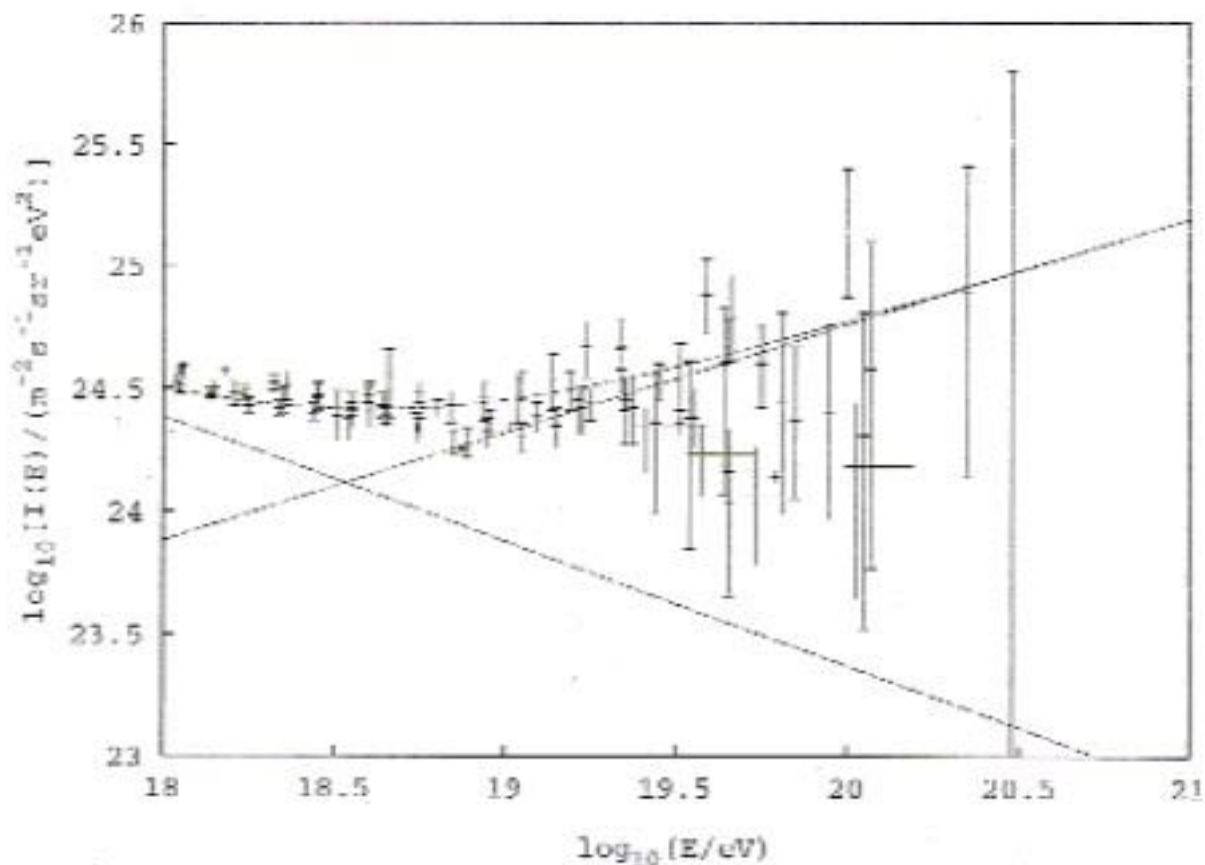


FIG. 1. The high energy cosmic ray spectrum beyond the 'ankle'. (Note that the differential flux has been multiplied by  $E^2$ .) The data shown are from AGASA, stereo Fly's Eye, Haverah Park and Yakutsk and are AGASA-normalized [8]. The highest energy monocular Fly's Eye event at  $3 \times 10^{20}$  eV is also shown. A fit to the spectrum from the superposition of a steeply falling and a flatter power law (dashed lines) is indicated [2].

↳ Fly's Eye PRL

## AGASA sees

3 pairs and 1 triplet

within  $\theta_{\text{resolution}} \sim 2.5^\circ$

$P(\text{chance}) < 1^\circ$

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Highly Significant:

★ Cosmic  $\vec{B}$  bends charged-particles

★ Bend is  $E$ -dependent

No Bending  $\Rightarrow$

- close source [unlikely]
- no  $\vec{B}$  [untenable]
- $Q = 0$

No GZK Cutoff  $\Rightarrow$

- close source
- $Q = 0$ , mag. moment  $\sim 0$

**$\gamma$ 's ARE PROPAGATING PARTICLE ?!**

## Correlation between Compact Radio Quasars and Ultra-High Energy Cosmic Rays

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(June 17, 1998)

### Abstract

Some proposals to account for the highest energy cosmic rays predict that they should point to their sources. [ We study the five highest energy events ( $E > 10^{20}$  eV) and find they are all aligned with compact, radio-loud quasars. The probability that these alignments are coincidental is 0.005, given the accuracy of the position measurements and the rarity of such sources. The source quasars have redshifts between 0.3 and 2.2. ] If the correlation pointed out here is confirmed by further data, the primary must be a new hadron or one produced by a novel mechanism. (large  $E_\nu$ ).

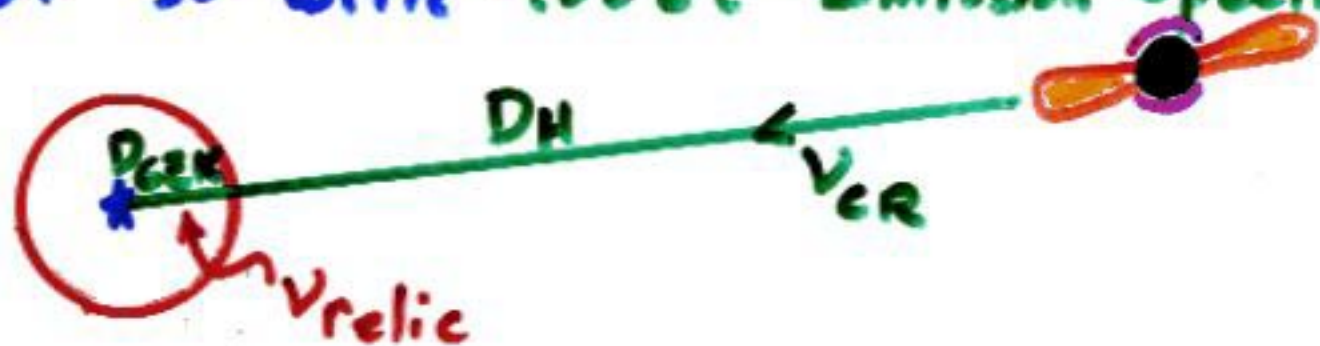
{pacs{}}

Typeset using REVTeX

With Absorption Spectroscopy,

high opacity is desirable.

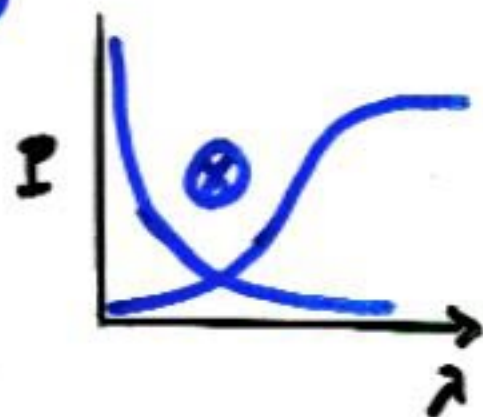
Not so with "local" Emission Spectroscopy.



$$P = P(\text{annihilate at } < D_{GZK})$$

$$\times P(\text{survive at } D_H)$$

$$\approx \frac{D_{GZK}}{\lambda} e^{-D_H/\lambda}$$



Exponential suppresses

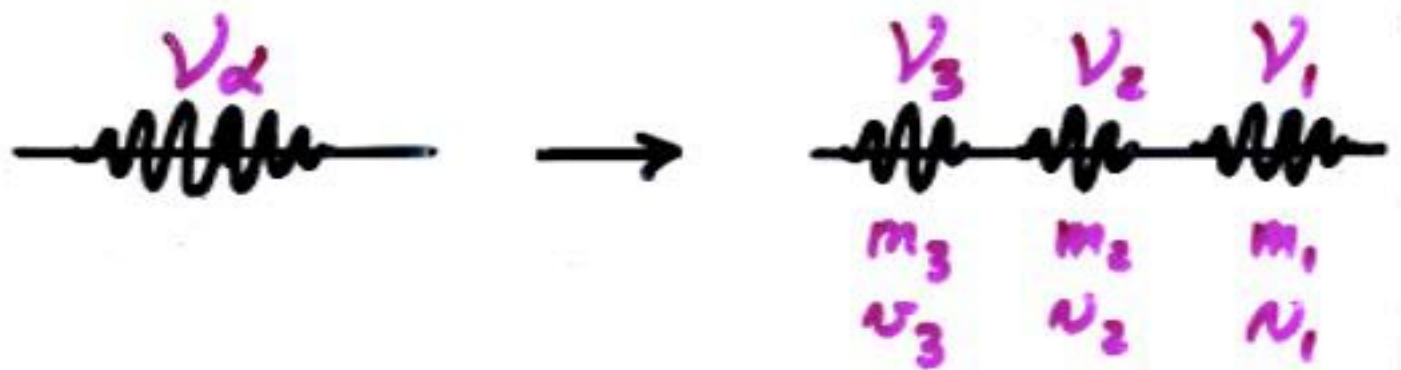
$\lambda < D_H$ , so find

$$P_{max} = P(\lambda = D_H) = \frac{D_{GZK}}{e D_H}$$

$$\sim \frac{1}{e} \frac{50 \text{ Mpc}}{5 \text{ Gpc}} \sim 0.3\%$$

# CYB = flavor

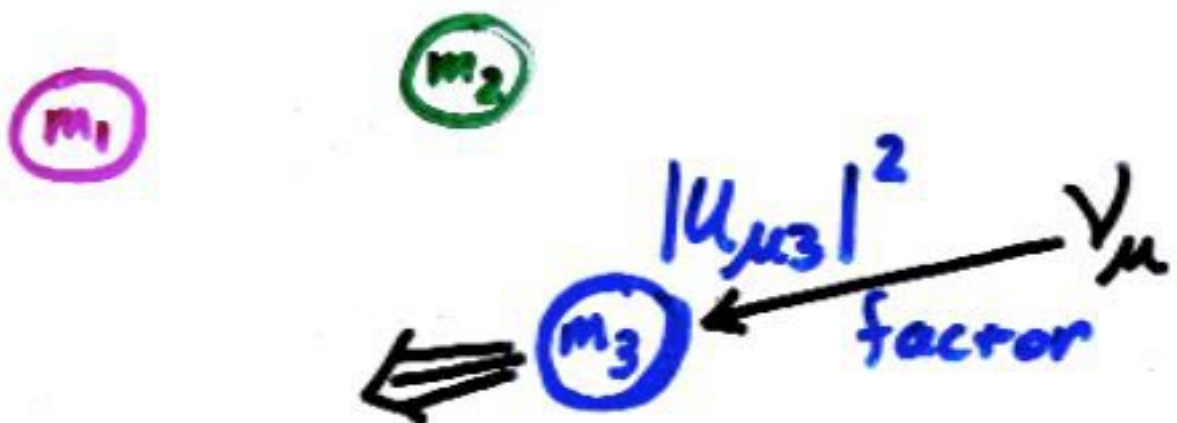
Decoherence after 12-15 Gyr



with  $\nu_j \sim \frac{3T}{m_j}$

$$P(\nu_j) = \sum_\alpha |U_{\alpha j}|^2 P(\nu_\alpha)$$

$$\underline{\underline{T \sim e_8^M}} \quad \frac{1}{N} \sum_\alpha |U_{\alpha j}|^2 = \frac{1}{N} \quad \forall j.$$



# Dirac vs. Majorana

When relativistic,  
100% polarized:



↓ red-shift by  $z_0 \sim 10^{10}$

Nonrelativistic today,  
Unpolarized:



$$\text{ie } \begin{pmatrix} \nu_L \\ \bar{\nu}_R \end{pmatrix} \xrightarrow{\text{red-shift}} \begin{pmatrix} \boxed{50\% \nu_L + 50\% \nu_R} \\ \boxed{50\% \bar{\nu}_R + 50\% \bar{\nu}_L} \end{pmatrix}$$

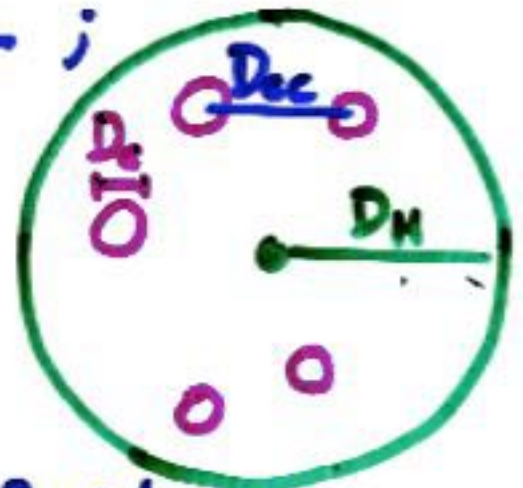
active      sterile if  $\nu_D$   
active if  $\nu_M$

∴ Dirac CνB half as active as  
Majorana CνB.

# Enhancement from Clustering:

$$\xi \equiv \frac{n_V}{\langle n_V \rangle} = \frac{V_{uni}}{N_c V_c};$$

$$N_c = \frac{V_{uni}}{L_{cc}^3};$$



$$\Rightarrow \xi \sim \left( \frac{L_{cc}}{L_c} \right)^3 \xrightarrow{\substack{\text{S-cluster} \\ L_c \sim 20 \text{ Mpc} \\ L_{cc} \sim 100 \text{ Mpc}}} \sim 100$$

$$* \text{ Path Length} \sim \frac{L_c}{D_H} \cdot \frac{20 \text{ Mpc}}{5000 \text{ Mpc}} \sim 4 \cdot 10^{-3};$$

for  $h = 0.6$

$$\Rightarrow \text{"Local" Rate} \sim \frac{L_{cc}^3}{L_c^2 D_H} \sim 0.4$$

$$\Rightarrow P(\text{z-burst within S-cluster}) \sim 3\%$$



# LEPTON ASYMMETRY

increases  $\nu$ -DENSITY

$$|n_\nu - n_{\bar{\nu}}| = 0.025 [\pi^2 \xi + \xi^3]$$

$$\xi \gtrsim 1 \quad n_\nu + n_{\bar{\nu}}$$

From BBN : from Structure Formation,

$$-0.06 \lesssim \xi_{\nu_e} \lesssim 1.1$$

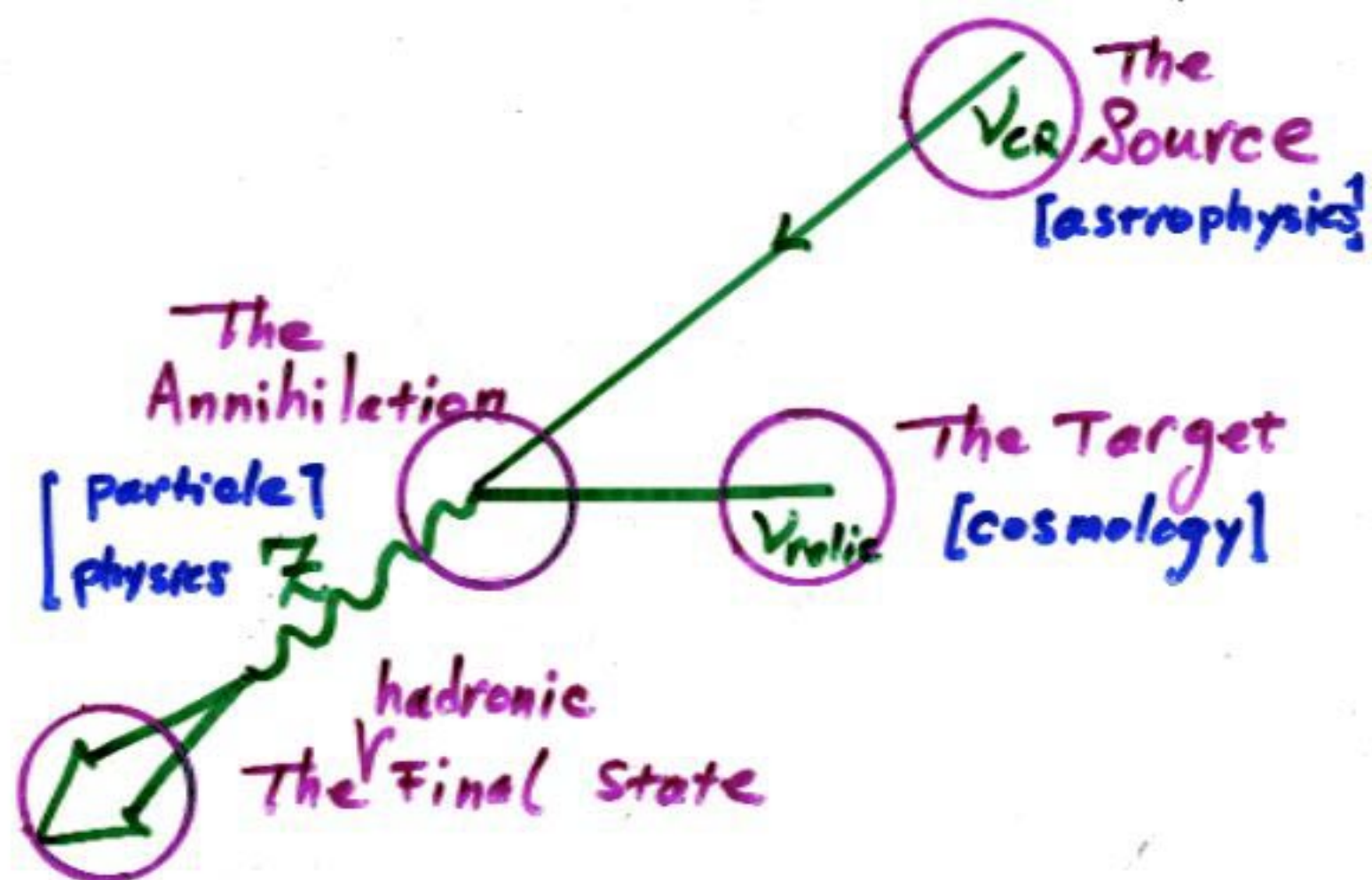
$$|\xi_{\nu_{\mu, \tau}}| \lesssim 6.9$$

Gelmini, Kusenko (99) :

$P$  (Local annihilation) optimized (ie  $\lambda \sim D_H$ )

with  $\xi = 5$

$$\leftrightarrow n_\nu = 30 \times n_\nu(\xi=0).$$



Flux of "Z-bursts" produced in distance  $D$  is

$$F_Z(D) = \int dE F_V(E, 0) [1 - e^{-\tau_{\text{ann}}(E) S(D)}]$$

where relic column density is

$$S(D) = \int_0^D n_{V_{\text{relic}}}(x) dx,$$

$$\approx \underbrace{E_R F_V(E_R, 0)}_{\text{source}} \underbrace{\left[ \int \frac{dE}{E_R} \sigma_{\text{ann}}(E) \right]}_{\text{annihilation}} \underbrace{S(D)}_{\text{target}}$$

# NEUTRINO FLUX ISSUE

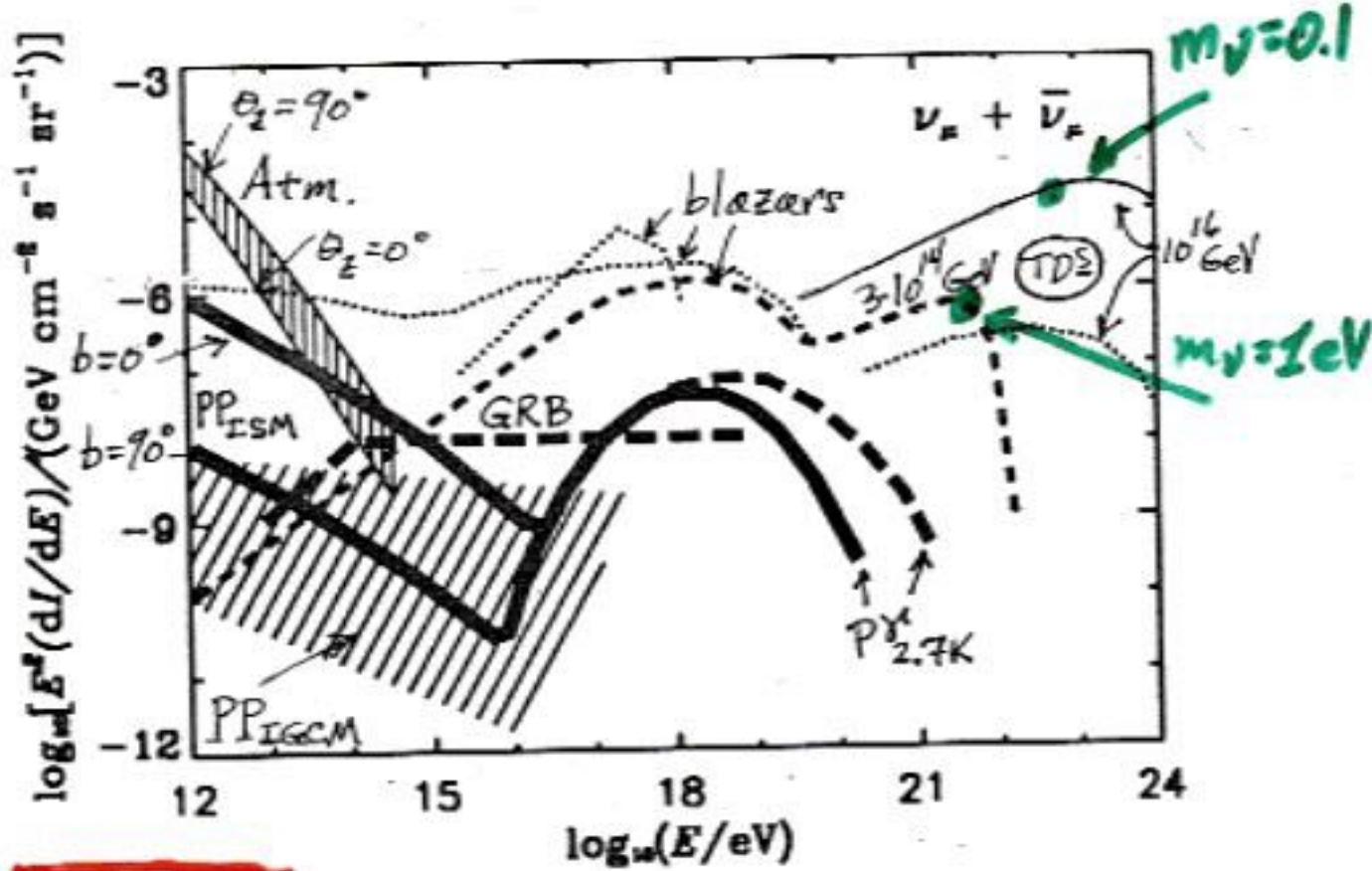
$$F_{\text{observed}} (\geq E_{68k})$$

$$\sim \underbrace{\text{Prob}(\nu \rightarrow Z)}_{\sim 17\%} \times E_R = F_\nu(E_R) \times \underbrace{\langle N \rangle_Z}_{30}$$

$\uparrow$   
 $4 \cdot 10^{21} \text{ eV}/\text{Mpc}$

$$\Rightarrow F_\nu (\geq E_R \sim 10^{22}) \sim 3 \cdot F_{\text{obs}} (\geq 10^{20})$$

Need  $F_\nu(E_R) \sim \frac{F_{\text{observed}}(E > E_{\text{GZK}})}{E_R P(\nu \rightarrow \bar{\nu}) \langle N_{\nu\gamma} \rangle}$



**Protheroe**

Grand Unified Neutrino Spectrum – a personal opinion about the predicted neutrino intensities: thick solid lines – certain; long dashed lines – almost certain; short dashed lines – speculative; dotted lines – highly speculative.

Coming:

? ! \* ? \* \* !

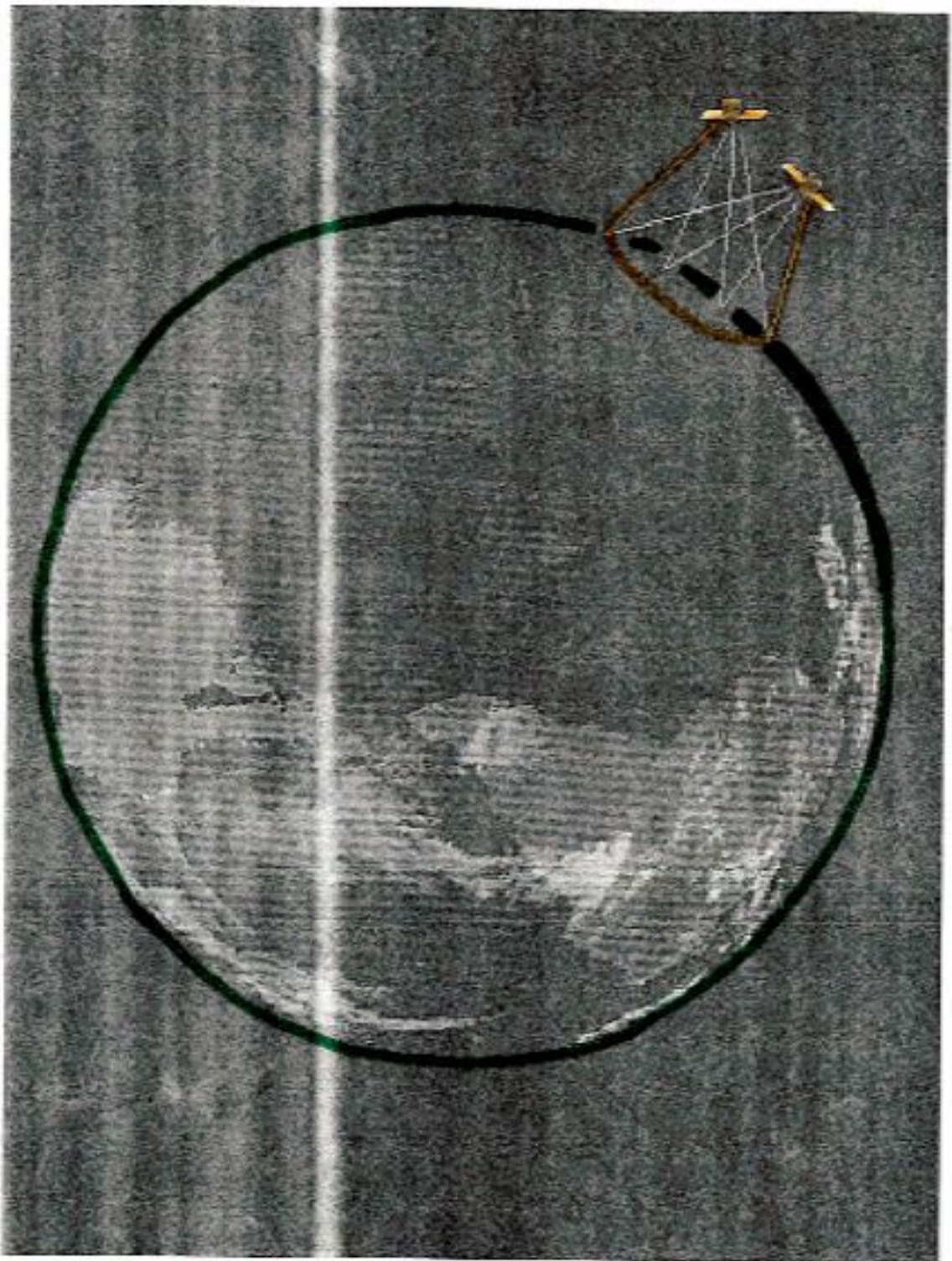
	AGASA	Hi Res	Telescope Array + Auger	OWL Air Watch
Area*	100	$10^3$	$10^4$	$10^5 - 10^6$
No. Evts > $10^{20}$ eV	1	10	100	$> 10^3$

\*  $\text{km}^2 \cdot \text{steradian}$

\*  $\text{yr}^{-1}$

D. Cline: if  $F_{\nu}(E_{\nu})$  exists &  $\Theta_{\nu 2}$  is max,  
then upcoming  $\nu_{\nu}$ 's!

**The Mission Concept for OWL Which Will Detect  
Cosmic Rays at the Highest Energies**



# Summary

[T.W. hep-ph/9710431]

If Nature provides

(i)  $m_\nu \sim \text{few} \times 10^{-12} \text{ eV}$ ,

(ii)  $F_\nu \text{ at } E = \frac{4 \times 10^{21} \text{ eV}}{m_\nu}$ ,

then,

$$\text{get } F_{\text{obs}}(> E_{\text{GZK}}) \sim \underbrace{P(\nu \rightarrow \bar{\nu})}_{\sim 1\%} \cdot \underbrace{E F(E_R)}_{R^2 \nu R} \cdot \underbrace{\langle N \rangle_E}_{\sim 30}$$

with  $\nu$ -clustering,  $10^{-4}$  without;  
up to 0.37. with Lepton Asym.

Potential payoff is HUGE -  
inferring relic  $\gamma$  background  
(and  $\nu$ -mass)

validating Big Bang back to  $t = 1 \text{ sec!}$