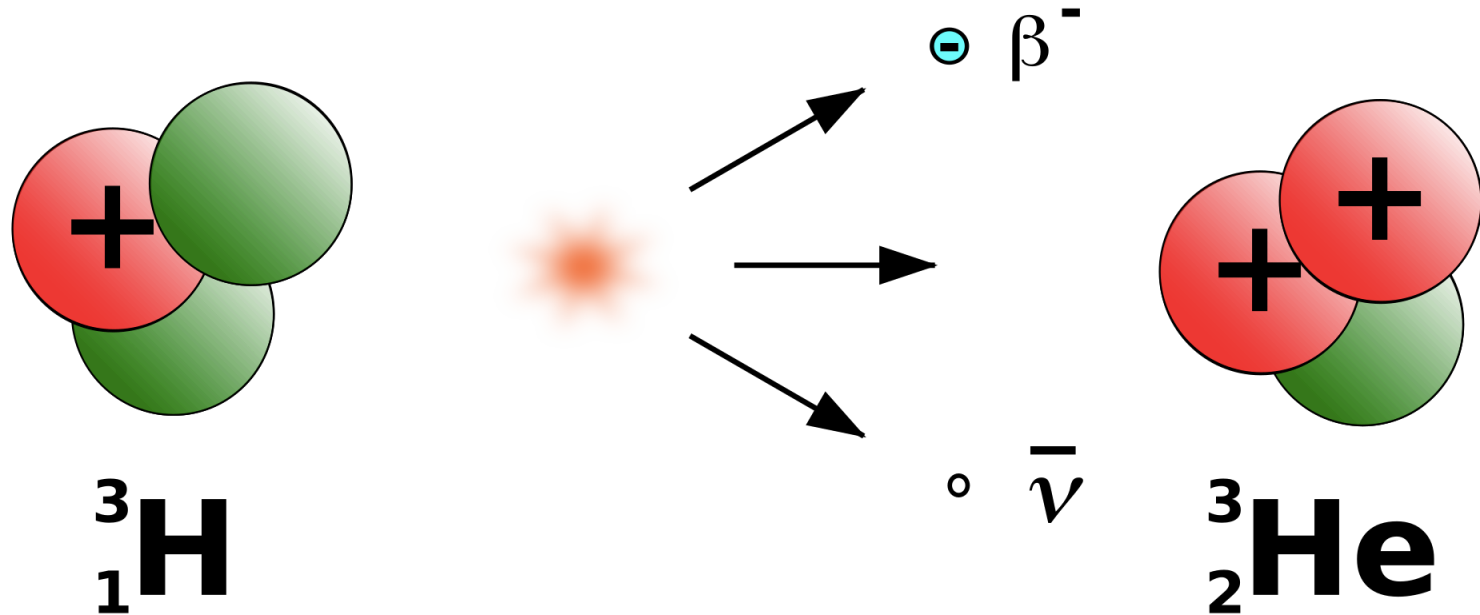


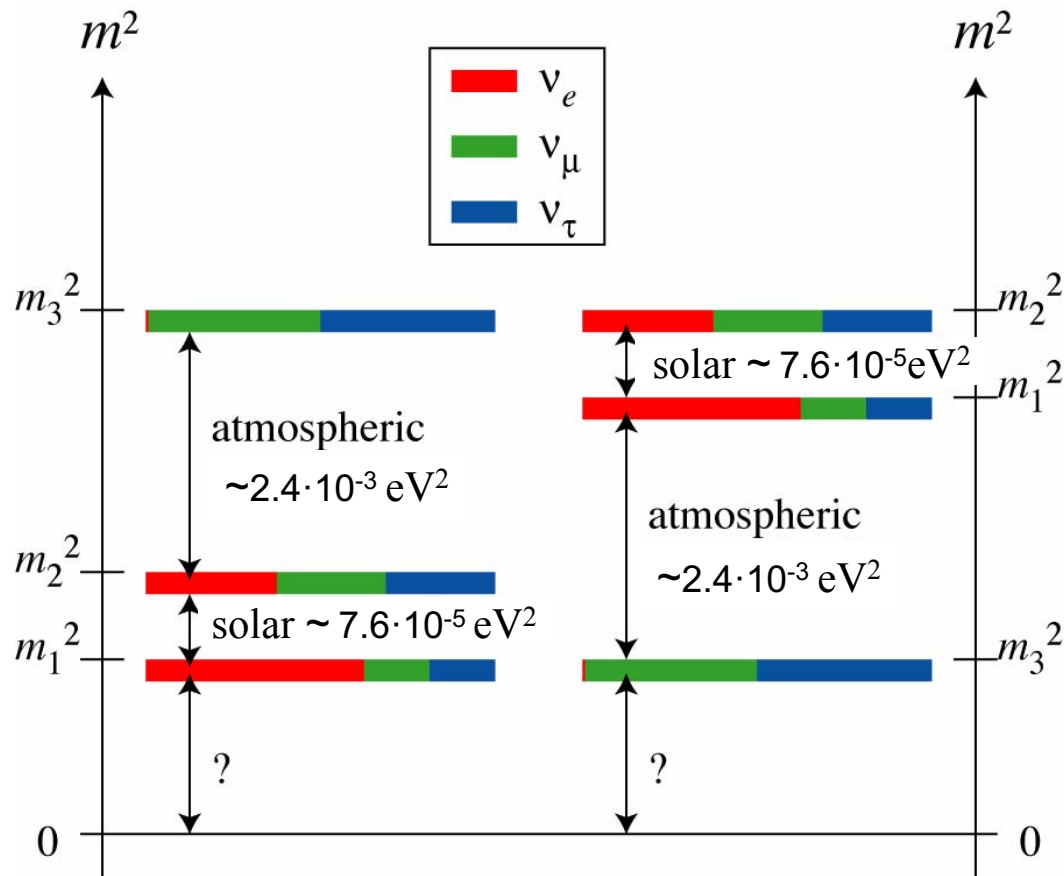
Neutrino Physics 2



*with thanks to Diana Parno,
John Wilkerson, and others,
from whom I borrowed liberally

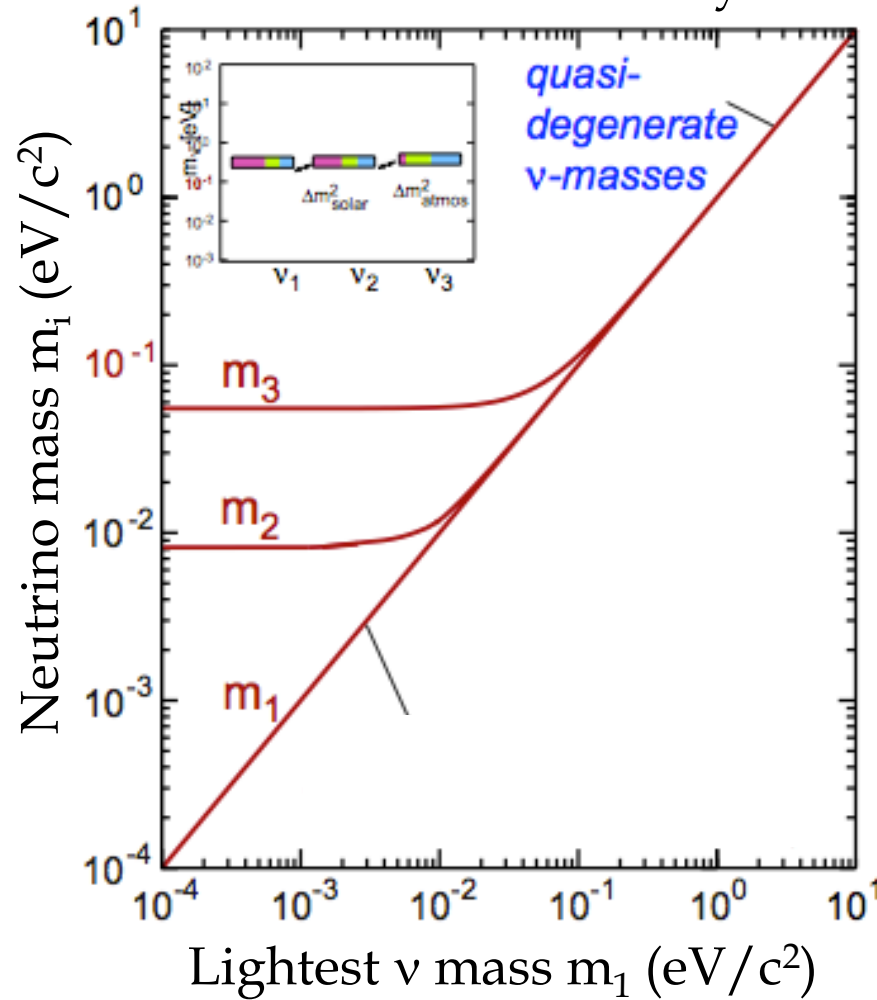
Michelle Dolinski
Drexel University
NNPSS, 11 July 2019

What we know about neutrinos



Absolute neutrino masses

Normal hierarchy



$\sim 2 \text{ eV}$

From tritium endpoint
(Maintz and Troitsk)

$\sim 0.1 \text{ eV}$

From $0\nu\beta\beta$ if ν is Majorana

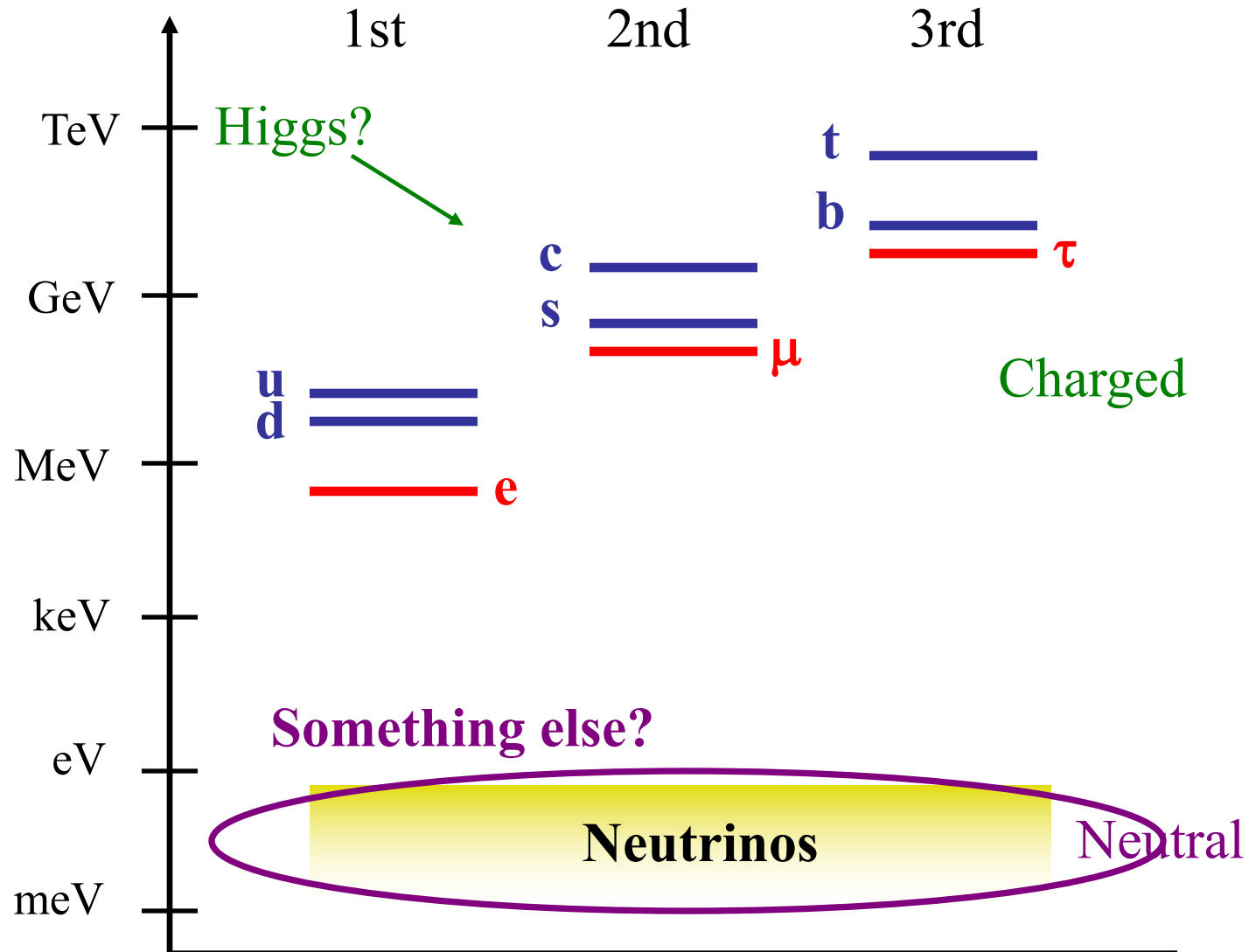
$\sim 20 \text{ eV}$

$\Sigma \sim 0.2 \text{ eV}$

From Cosmology

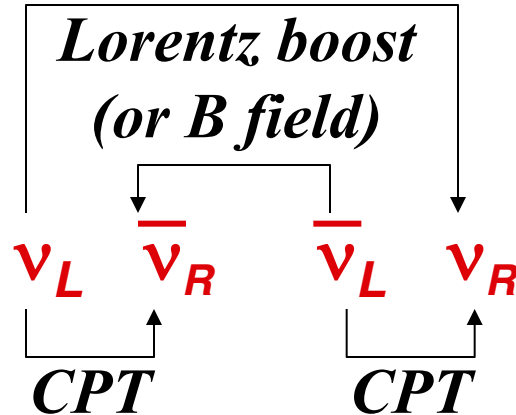
Time of flight from SN1987A
(PDG 2002)

Neutrino mass puzzle



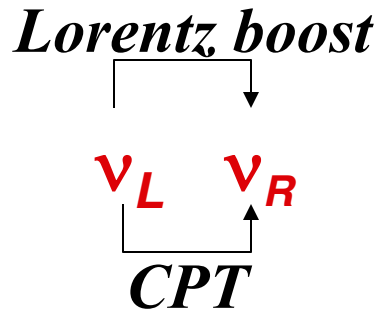
Massive neutrinos

“Dirac” neutrinos



“Majorana” neutrinos

No lepton number conservation!



The two descriptions are distinct and distinguishable only if $m_\nu \neq 0$.
Lecture 3 is how to tell the difference experimentally.

Seesaw mechanism

Experimentally, it is an open question whether neutrinos are Majorana or Dirac, but Majorana neutrinos are strongly preferred by theorists. Seesaw mechanism can be used to explain small neutrino masses (see 2019 PDG). Type I seesaw mechanism (Gell-Mann, Ramond, Slansky and Yanagida, 1979):

$$\mathcal{L}_{Y,M}(x) = \left(\lambda_{il} \overline{N_{iR}}(x) \Phi^\dagger(x) \psi_{lL}(x) + \text{h.c.} \right) - \frac{1}{2} M_i \overline{N_i}(x) N_i(x)$$

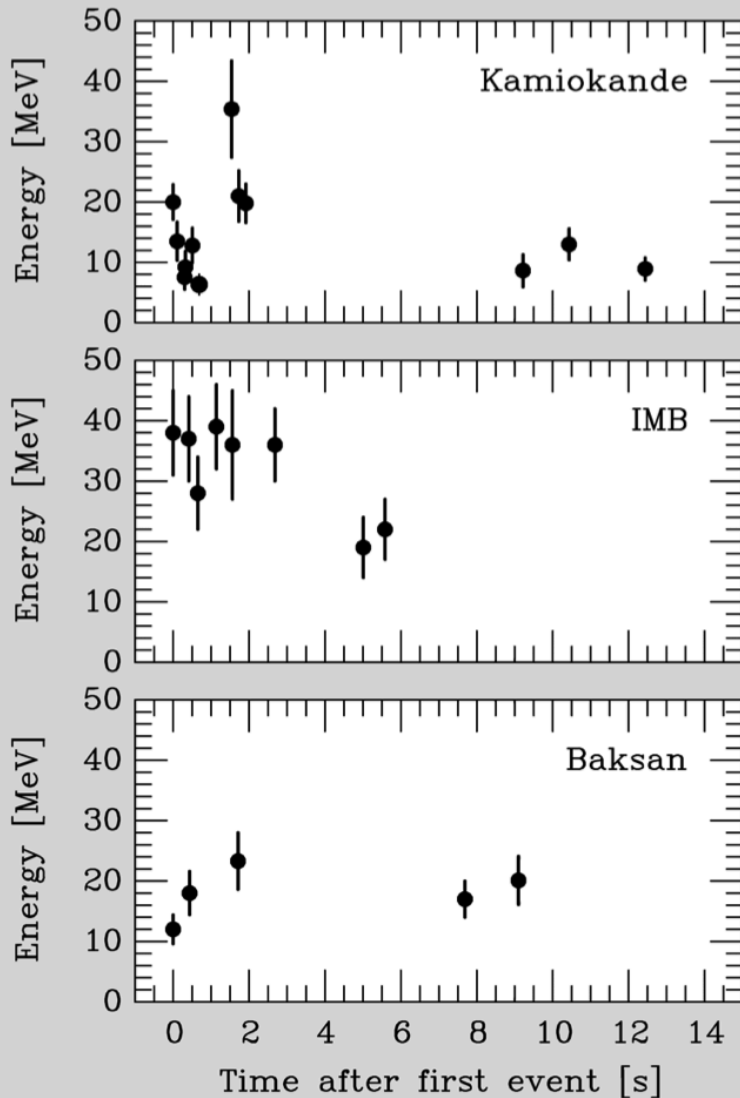
$$m_{ll}^{LL} \cong -(m^D)_{l'j}^T M_j^{-1} m_{jl}^D = -v^2 (\lambda)_{l'j}^T M_j^{-1} \lambda_{jl}$$

$$M_\nu = \begin{bmatrix} 0 & m_D \\ m_D & m_R \end{bmatrix}$$

$$Z^T M_\nu Z = D_\nu = \begin{bmatrix} m_1 & 0 \\ 0 & m_2 \end{bmatrix} \approx \begin{bmatrix} m_D^2 / m_R & 0 \\ 0 & m_R \end{bmatrix}$$

How to weigh a neutrino?

Time of flight



Neutrino events from supernova 1987a (Large Magellanic Cloud) were detected in KamiokaNDE, IMB, and Baksan observatories.

With a model for neutrino production, it is possible to look for smearing due to neutrino mass. Early analyses gave limits ~ 20 eV.

Improved supernova modeling and Bayesian statistical approaches do better:

< 5.7 eV @ 95% C.L.

Loredo and Lamb, *PRD* 65 (2002)

Decay kinematics

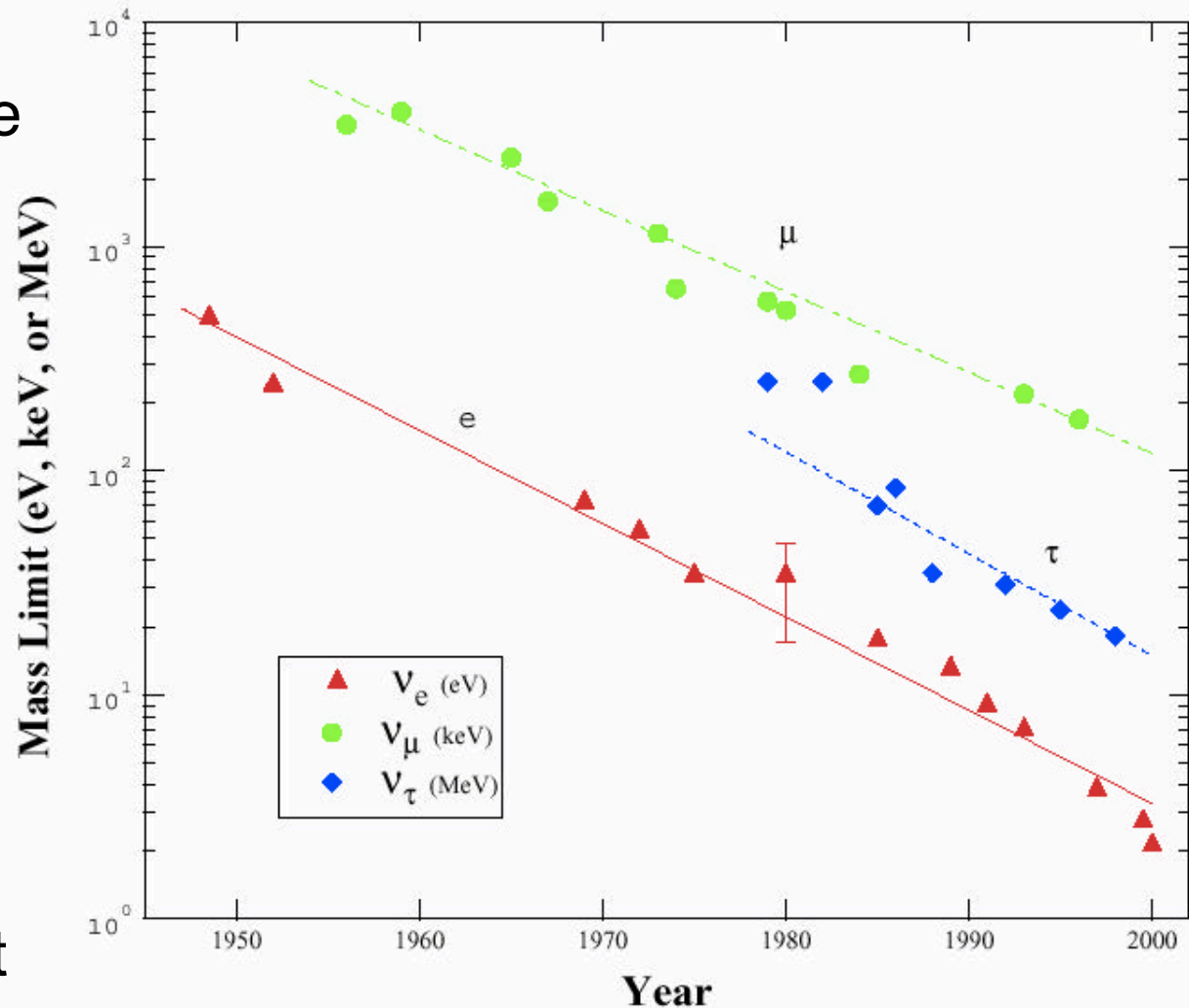
Look at the impact of non-zero ν mass on the following decays.

ν_e : beta decay

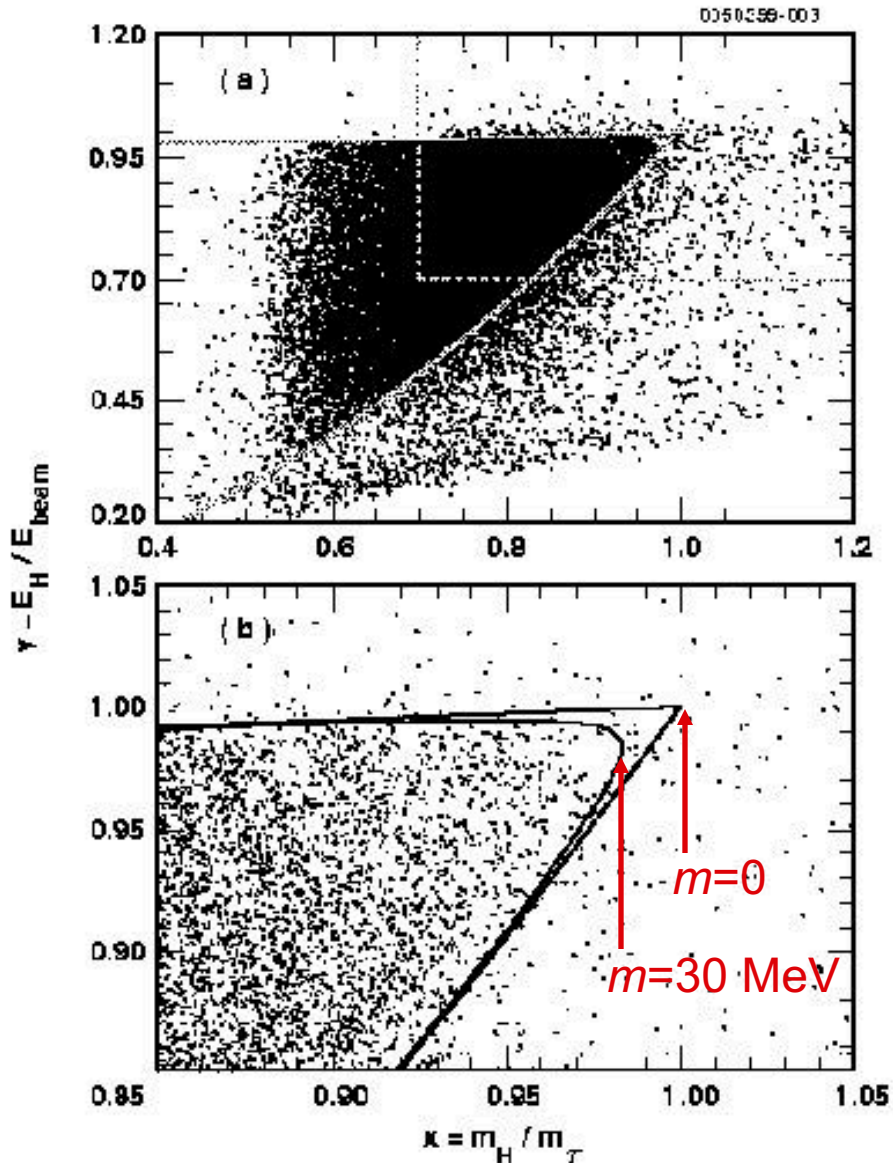
ν_μ : pion decay*

ν_τ : tau decay*

*thanks to Mike Shaevitz for next two slides, 2002 lectures at Lake Louise School



τ decay



Current best limit from studies of the kinematics of τ decays.

$$\tau^- \rightarrow 2\pi^- \pi^+ \nu_\tau$$

$$\tau^- \rightarrow 3\pi^- 2\pi^+ (\pi^0) \nu_\tau$$

• Fit to scaled visible energy vs. scaled invariant mass. Best limit

$<18.2 \text{ MeV @ 95\% C.L.}$
Aleph, EPJ C2 395 1998

Pion decay

Current best limit from studies of the kinematics of $\pi \rightarrow \mu \nu_\mu$ decay.

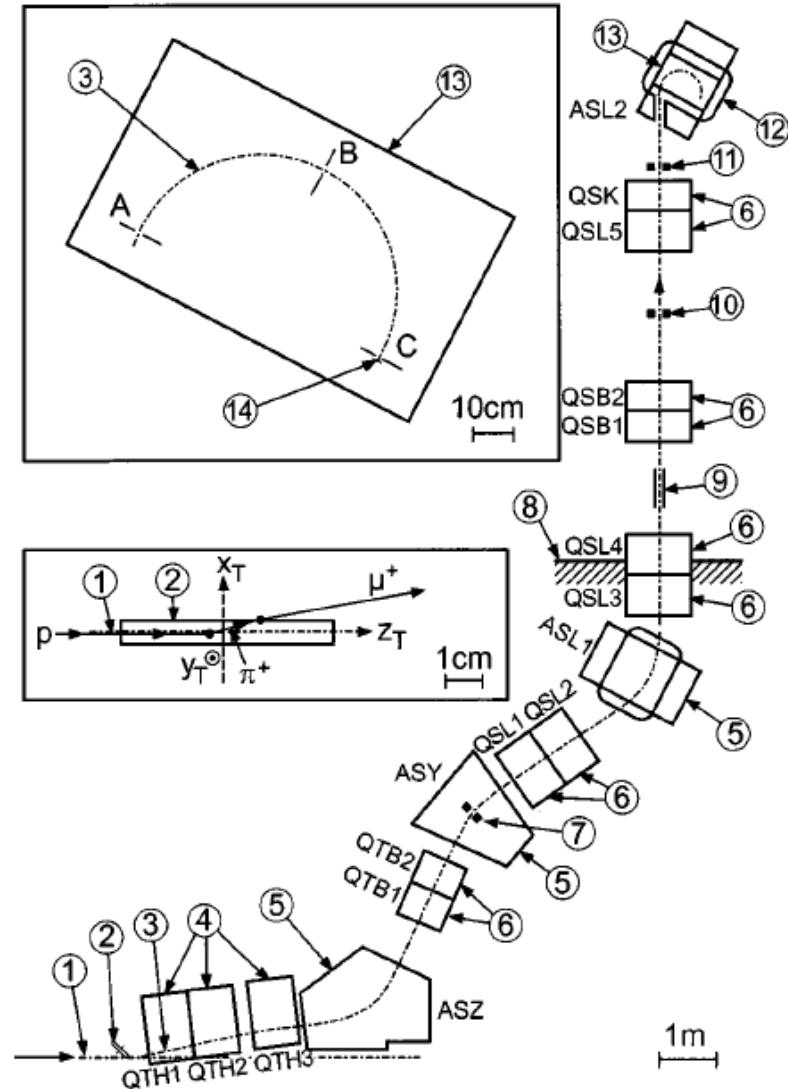
$$p_\mu^2 + m_\mu^2 = (m_\pi^2 + m_\mu^2 - m_\nu^2)^2 / 4m_\pi^2$$

- Pion decay in flight is limited in practice by momentum resolution.
- Pion decay at rest is limited by pion mass uncertainty. This currently gives the best limits from PSI

<170 keV @ 95% C.L.

Assamagan et al., *PRD* (1996)

*Proposals exist to get this down to ~8 keV



Beta decay

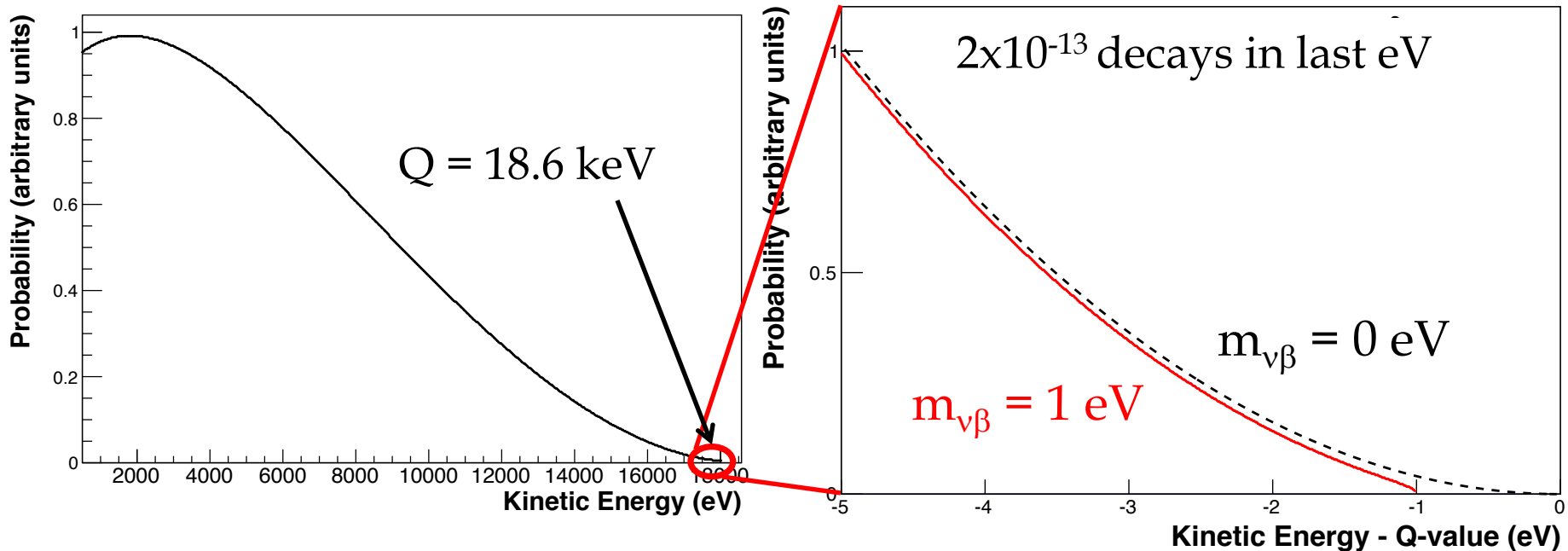
Taking into account ν mass eigenstates, the original spectrum

$$dN(E) = K|M|^2 F(Z,R,E) p_e E (E_0 - E) \{(E_0 - E)^2 - m_{\nu_e}^2 c^4\}^{1/2} dE$$

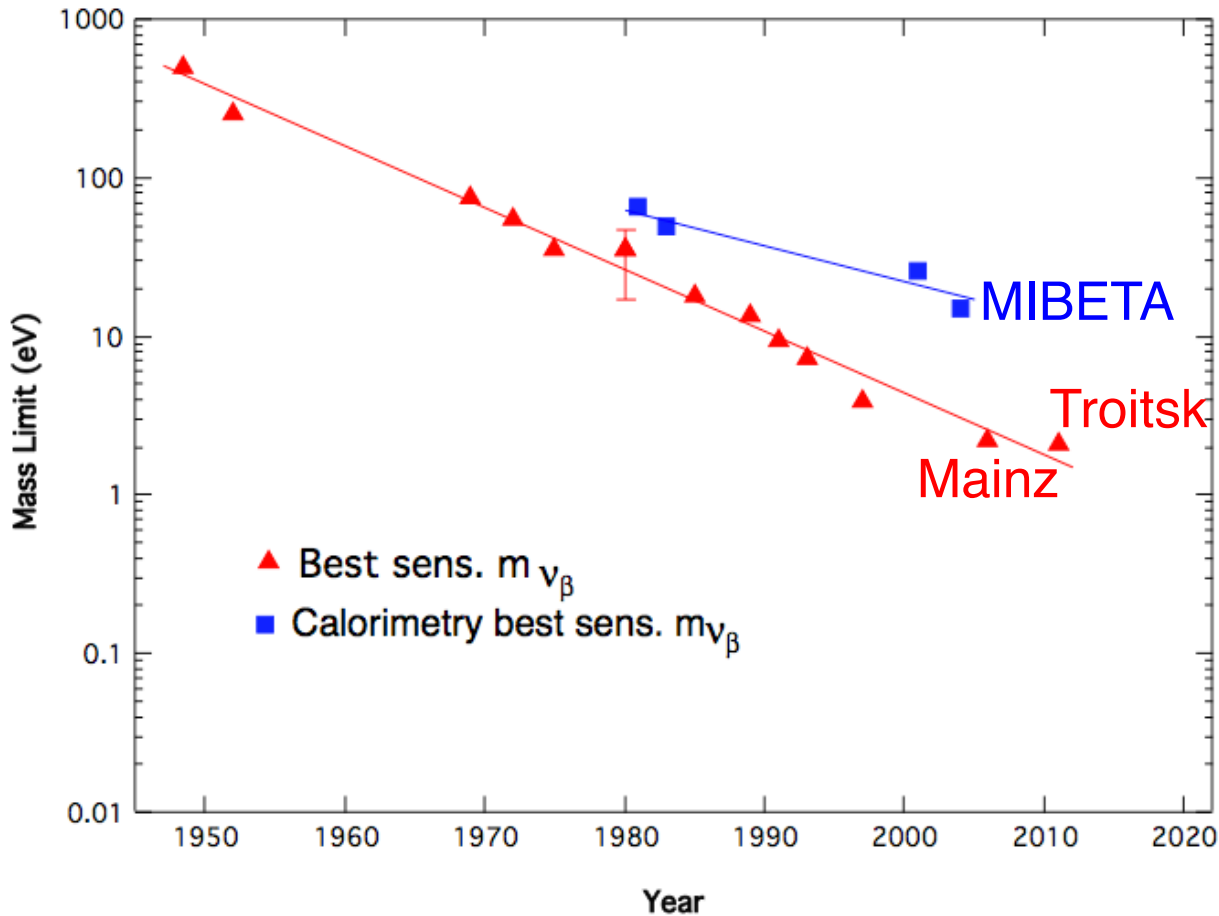
becomes

$$dN(E) = K|M|^2 F(Z,R,E) p_e E (E_0 - E) \sum_i |U_{ei}|^2 \{(E_0 - E)^2 - m_{\nu_i}^2 c^4\}^{1/2} dE$$

For 3 ν mass spectrum, with degenerate states, the beta spectrum simplifies to an “effective mass” : $m_\beta = \sum |U_{ei}|^2 m_{\nu_i}^2$



Beta decay limits



^{187}Re

$Q = 2.47 \text{ keV}$

$t_{1/2} = 4.5 \times 10^9 \text{ years}$

Forbidden

^3H (tritium)

$Q = 18.6 \text{ keV}$

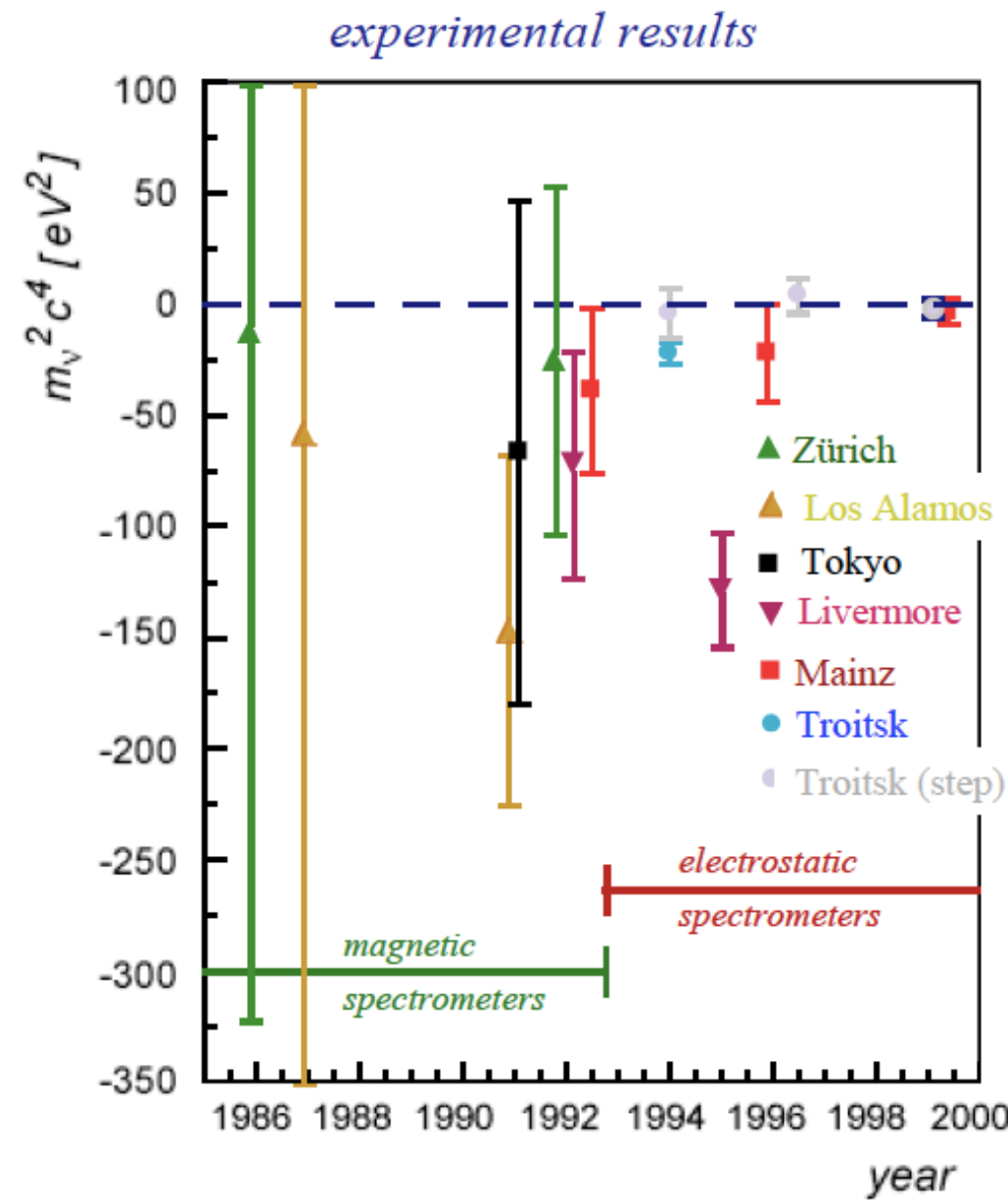
$t_{1/2} = 12.3 \text{ years}$

Super-allowed

Figure from J. Wilkerson, Neutrino 2012

Existing tritium results

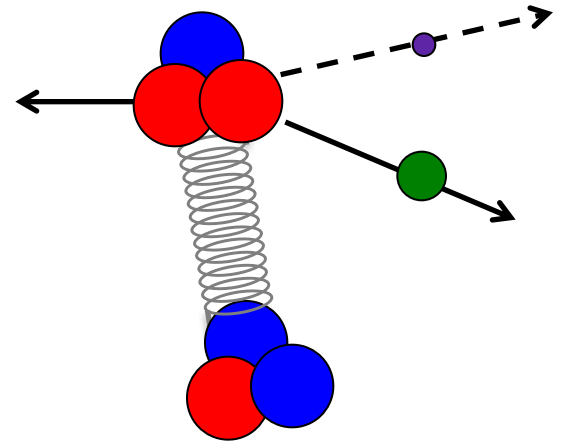
Location	Source	Instrument	Result (m_ν)
ITEP	T_2 in complex molecule	magn. spectrometer (Tret'yakov)	17-40 eV
Zürich	T_2 - source impl. on carrier	magn. spectrometer (Tret'yakov)	< 11.7 eV
Los Alamos	gaseous T_2 - source	magn. spectrometer (Tret'yakov)	< 9.3 eV
Tokyo	T - source	magn. spectrometer (Tret'yakov)	< 13.1 eV
Livermore	gaseous T_2 - source	magn. spectrometer (Tret'yakov)	
Mainz (1994-today)	frozen T_2 - source	electrostat. spectrometer	< 2.2 eV
Troitsk (1994-today)	gaseous T_2 - source	electrostat. spectrometer	(< 2.5 eV)



Tritium gas sources

Gas sources give the best results, but we're limited to using molecular tritium.

- Electronic excitations in T atoms
- Excitations in T₂ gas
 - Electronic: 20 eV
 - Vibrational: ~0.1 eV
 - Rotational: ~0.01 eV
- Beta spectrum depends on excitation energies V_k and probabilities P_k
- KATRIN needs 1% uncertainties on final state distribution.

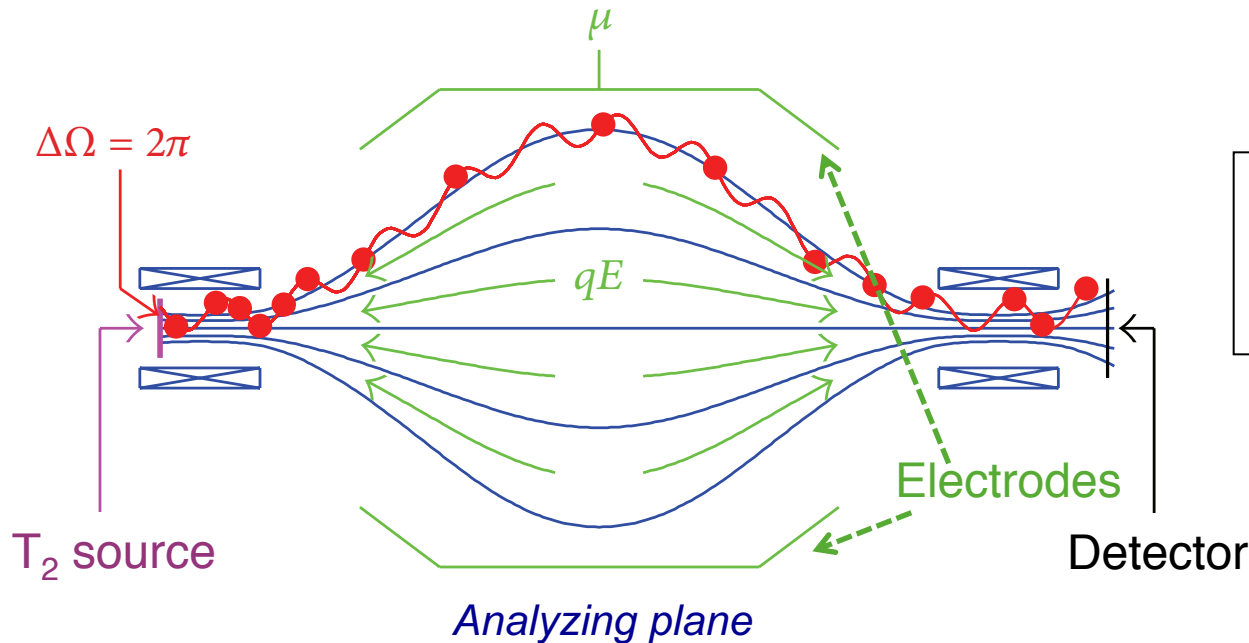


$$\frac{dN}{dE_e} = \frac{G_F^2 m_e^5 \cos^2 \theta_C}{2\pi^3 \hbar^7} |M_{\text{nuc}}|^2 F(Z, E_e) p_e E_e \times \sum_{i,k} |U_{ei}|^2 P_k(E_{\text{max}} - E_e - V_k) \times \sqrt{(E_{\text{max}} - E_e - V_k)^2 - m_{\nu i}^2} \times \Theta(E_{\text{max}} - E_e - V_k - m_{\nu i})$$

MAC-E filter

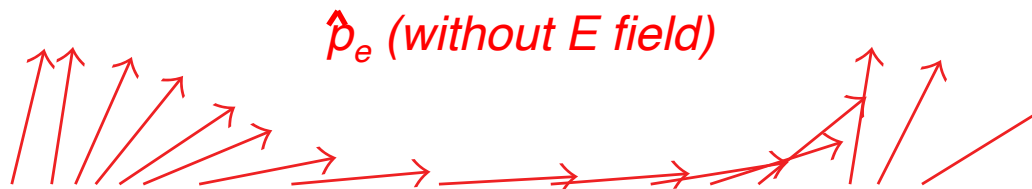
Magnetic Adiabatic Collimation and Electrostatic filter

The MAC-E filter allows measurement of integral spectrum with an adjustable threshold. Only see the endpoint of the decay!



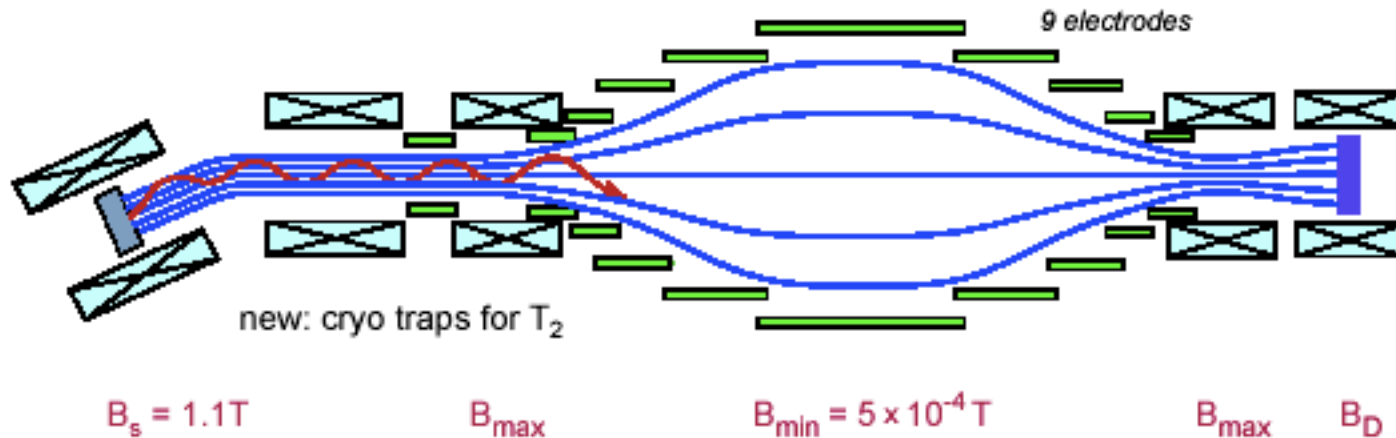
$$\mu = \frac{E_{\perp}}{B} = \text{const}$$

$$\frac{\Delta E}{E} = \frac{B_{\min}}{B_{\max}}$$



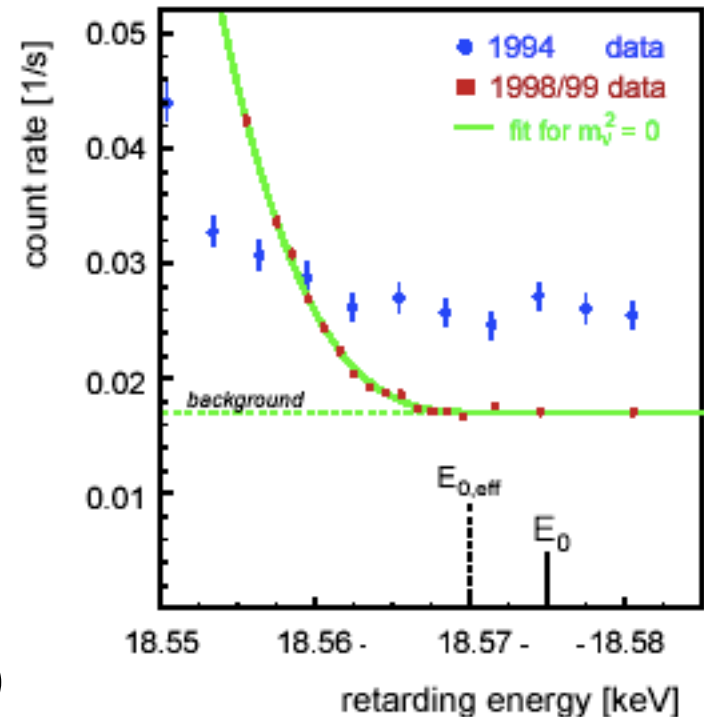
MAINZ

overall length source - detector ~6 m



- Quench condensed solid T_2 source
- Early results (1994) showed systematic effects, traced to source film roughening transition (fixed by lowering temperature)
- 1995-1997 significant background reduction, signal improvement
- Best limit

< 2.2 eV @ 90% C.L.



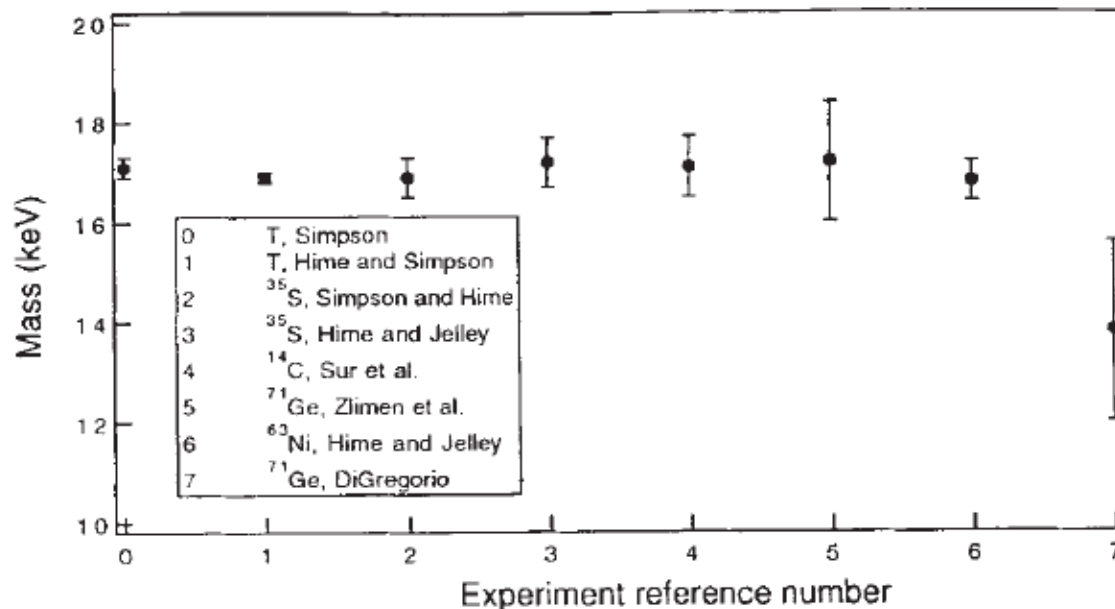
Speaking of anomalous results...

The rise and fall of the 17-keV neutrino

Douglas R. O. Morrison

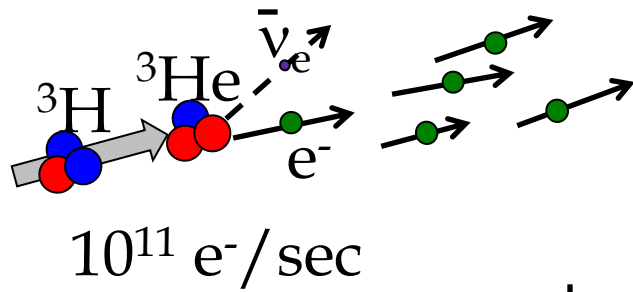
Nature 366, 29 - 32 (04 November 1993)

Experiments showing evidence for a heavy neutrino with a mass of 17 keV launched the new particle on an erratic eight-year career, during which it raised questions about the Standard Model of particle physics and about cosmological theories, stimulated many theoretical papers and pushed experimental techniques to their limit. Its demise provides grounds for faith in the efficacy of the scientific method.





KARlsruhe TRItium Neutrino experiment



Gaseous T_2 source

Electron transport

$10^3 \text{ e}^-/\text{sec}$

$1 \text{ e}^-/\text{sec}$

Detect β s

Analyze β energy

Monitor energy threshold



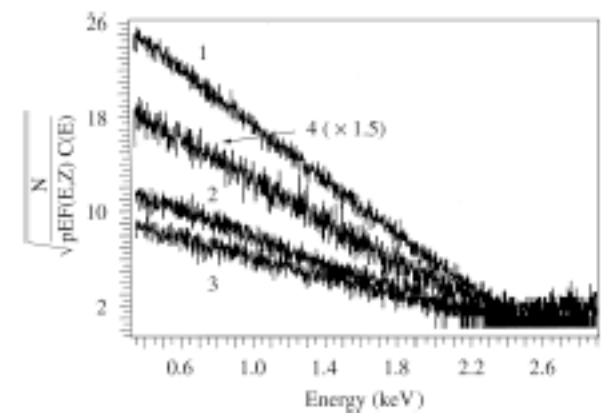
KATRIN outlook

- Intense T_2 source (10^{11} decays/second)
- Spectrum analysis with electromagnetic filter
- Design resolution 0.93 eV
- Design m_β sensitivity: **0.2 eV/c² at 90% C.L.**
- 2015 commissioning revealed high backgrounds from ^{210}Po
- Tritium commissioning began in May 2018! Stay tuned...

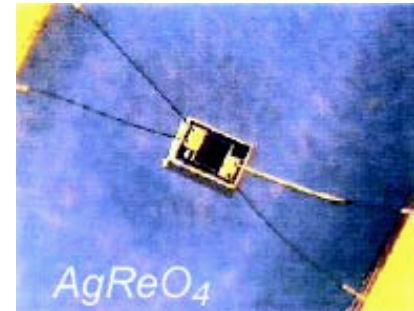


MIBETA and MARE

Bolometric measurements on ^{187}Re



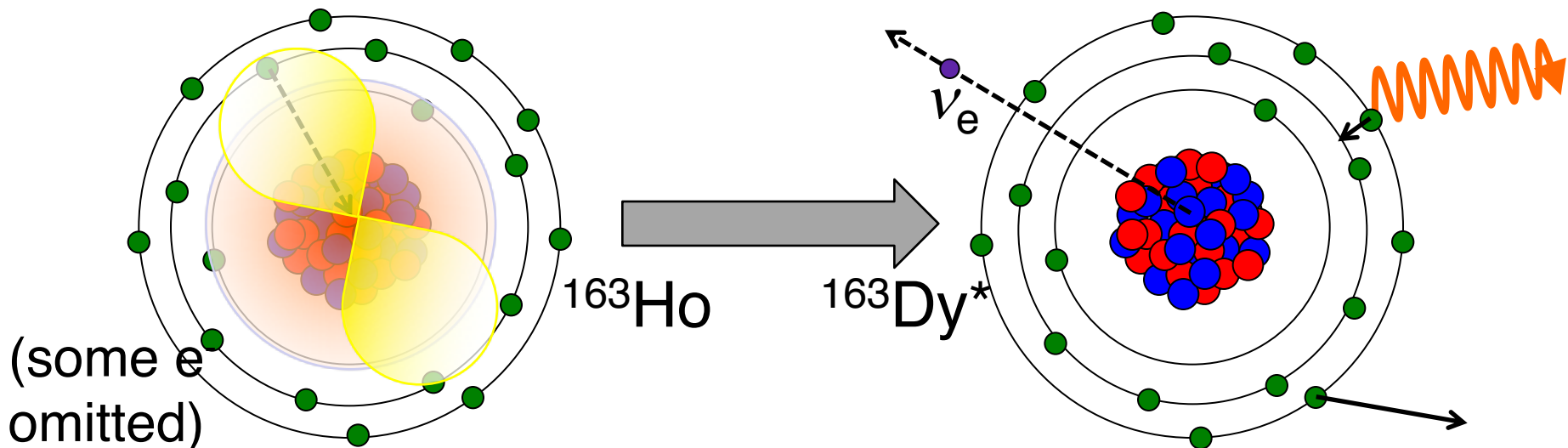
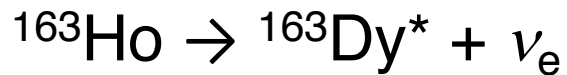
- ~ 15 eV sensitivity for MiBETA (2004)
- R&D by MARE collaboration
 - Metallic Re (superconducting)
 - Complex thermalization
 - Dielectric AgReO_4
 - Long response time
- Low specific activity (10^5 pixels!)



Community has moved on to ^{163}Ho

Nucciotti, arXiv:1511.00968

Holmium decay



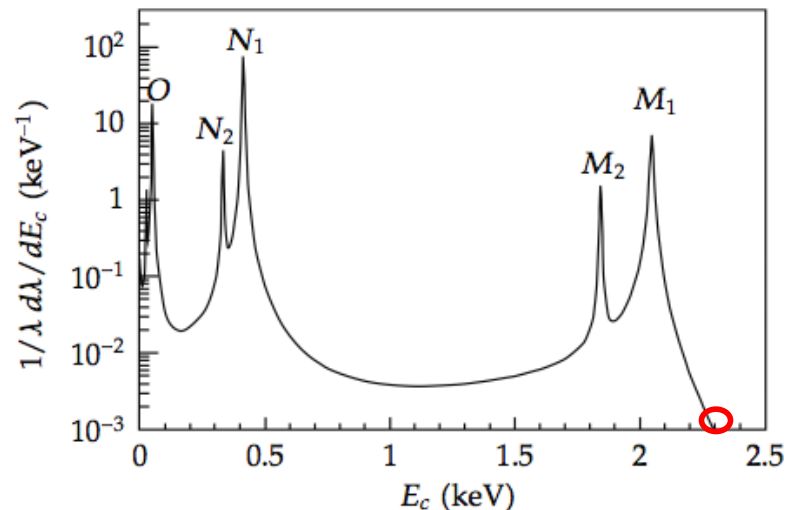
^{163}Ho

$Q = 2.83 \text{ keV}$

$t_{1/2} = 4750 \text{ years}$

*De Rújula and Lusignoli,
Phys. Lett. B **118** (1982) 429*

*Lusignoli and Vignati, Phys. Lett. B **697** (2011)*





Holmium experiments



Radioisotope



Thermometer

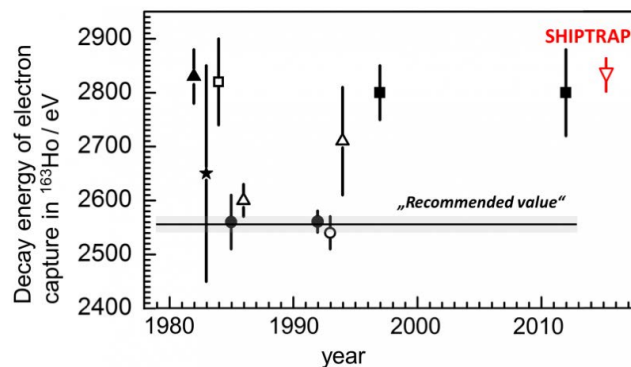
Thermal link

Heat sink
(sub-Kelvin)

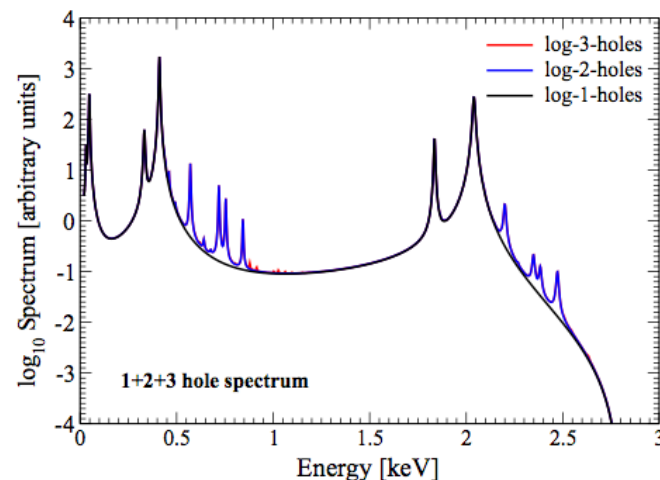
Holmium microcalorimetry, two competing experiments.

- HOLMES uses MKID sensor technology, ECHO uses MMCs

Big problems with the endpoint and theory



*Eliseev et al., PRL 115
(2015) 062501*



*Faessler et al., PRC
91 (2015) 064302*

A new approach:

PROJECT 8



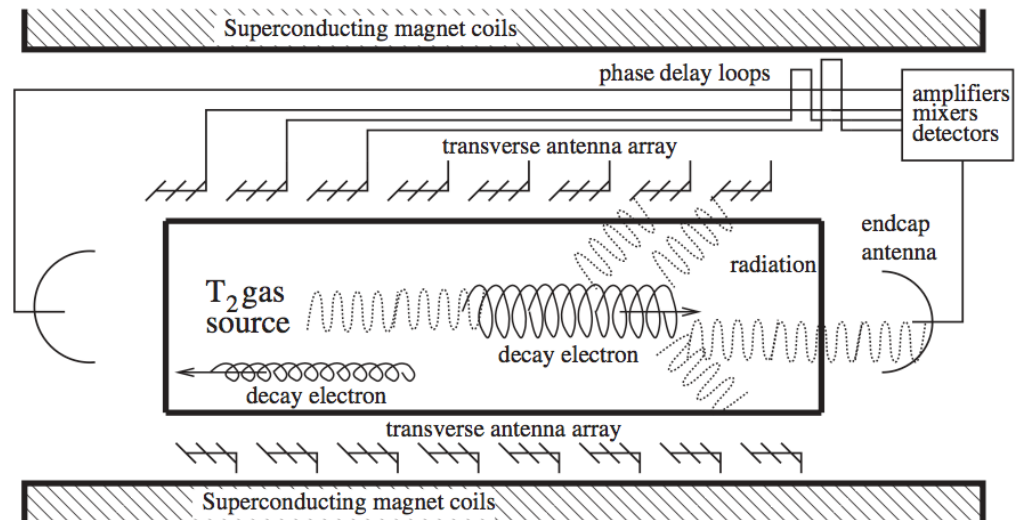
Never measure anything but frequency.

-Arthur Schawlow

An electron in a magnetic field will radiate at:

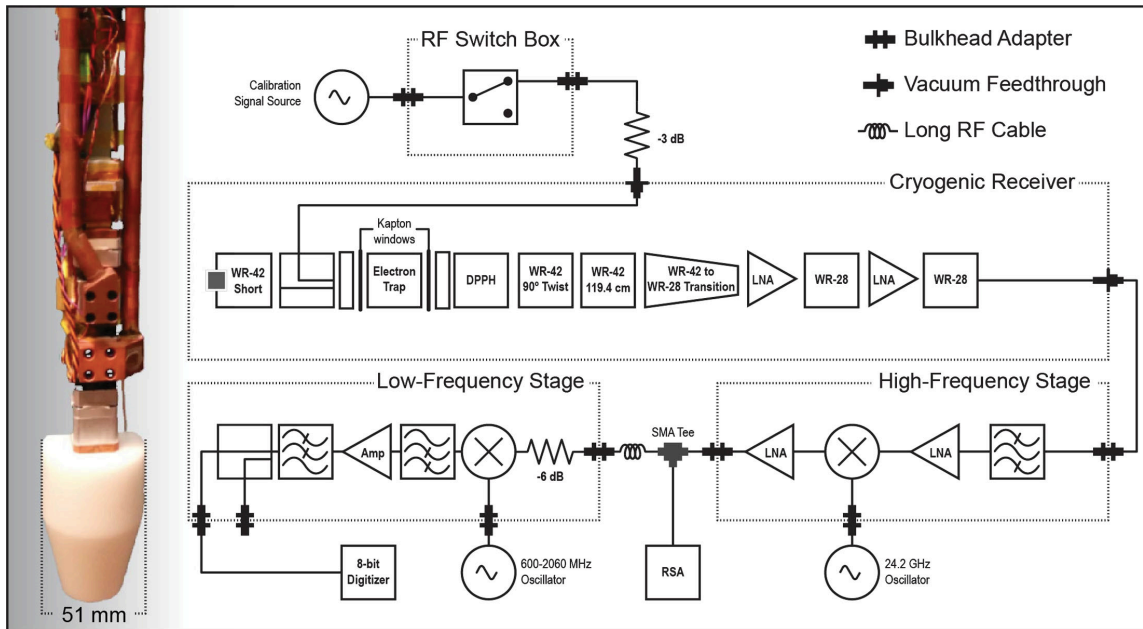
$$f_{\gamma} = \frac{f_c}{\gamma} = \frac{eB}{2\pi m_e} \frac{1}{1 + \frac{1}{c^2} E_{\beta}^2}$$

Measure entire beta spectrum at once:
Cyclotron Radiation Emission Spectroscopy (CRES)



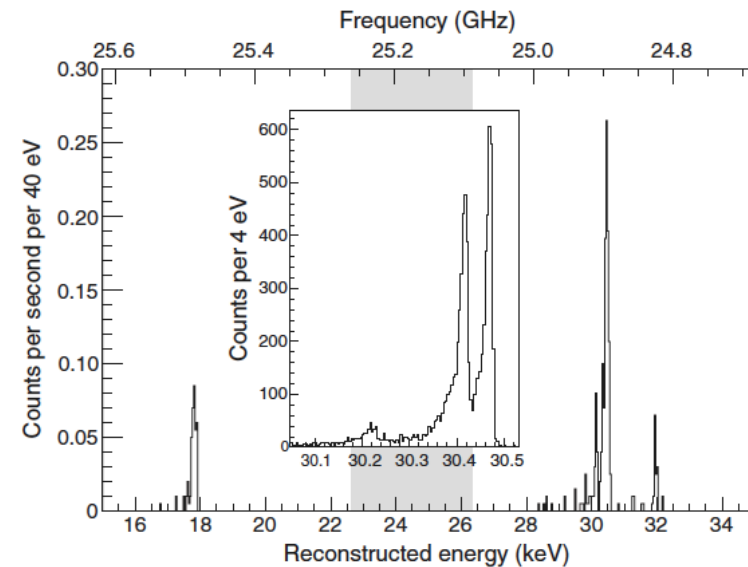
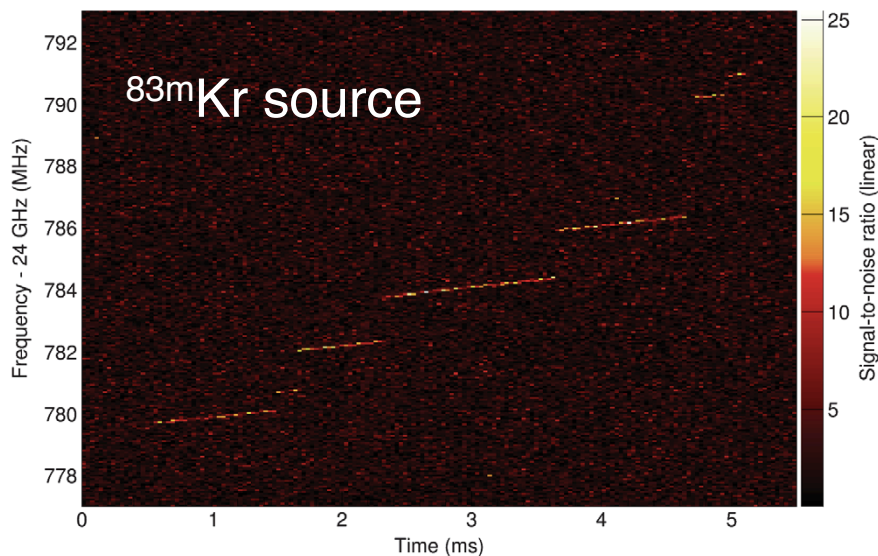
Monreal and Formaggio, PRD 80 (2009)

CRES single electron detection



PROJECT 8

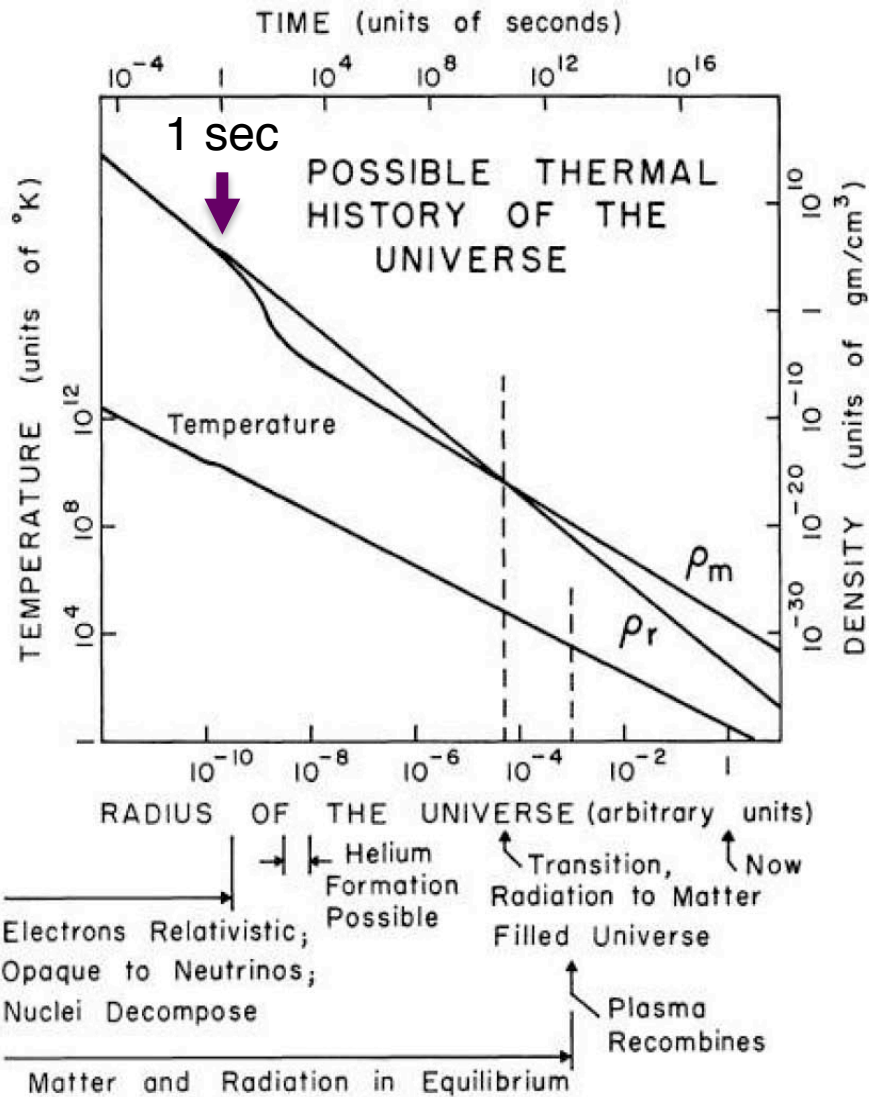
PRL 114, 162501 (2015)



Cosmic neutrino background

Next few slides stolen from a talk by Chris Tully at LNGS in 2017
for more on PTOLEMY, see [arXiv:1902.05508](https://arxiv.org/abs/1902.05508)

Cosmic neutrino background



$$n_\nu = \left(\frac{3}{4}\right)\left(\frac{4}{11}\right)n_\gamma = 112/\text{cm}^3$$

per neutrino species
(neutrino+antineutrino)

$$T_\nu(t) = \left(\frac{4}{11}\right)^{1/3} T_{CMB} \quad T_\nu \sim 1.95\text{K}$$

start of nucleosynthesis
 $n/p \sim 0.15 * 0.74 \sim 0.11$

$$\frac{\lambda(p \rightarrow n)}{\lambda(n \rightarrow p)} = e^{-Q/kT}$$

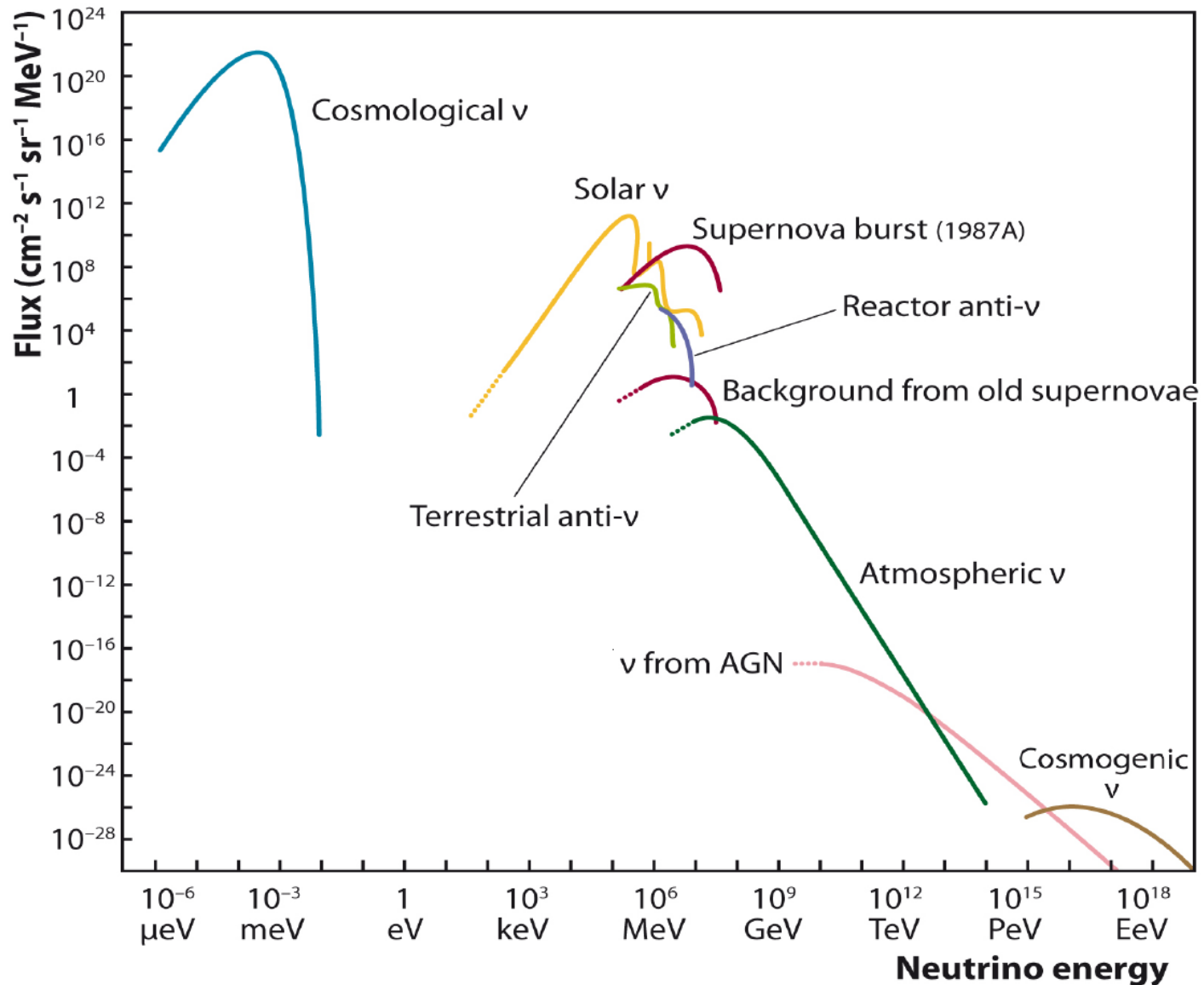
Relic velocity depends on mass

$$\langle v_{rms} \rangle \propto T/m_\nu \text{ instead of } \propto \sqrt{T/m_\nu}$$

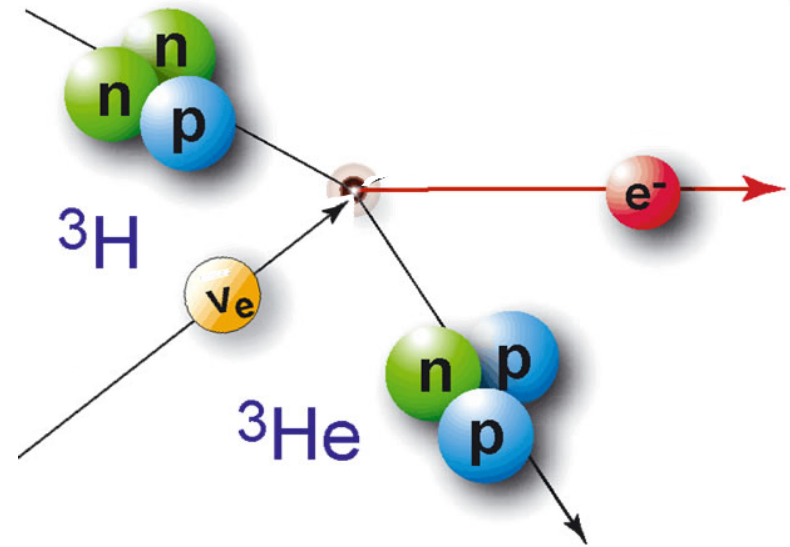
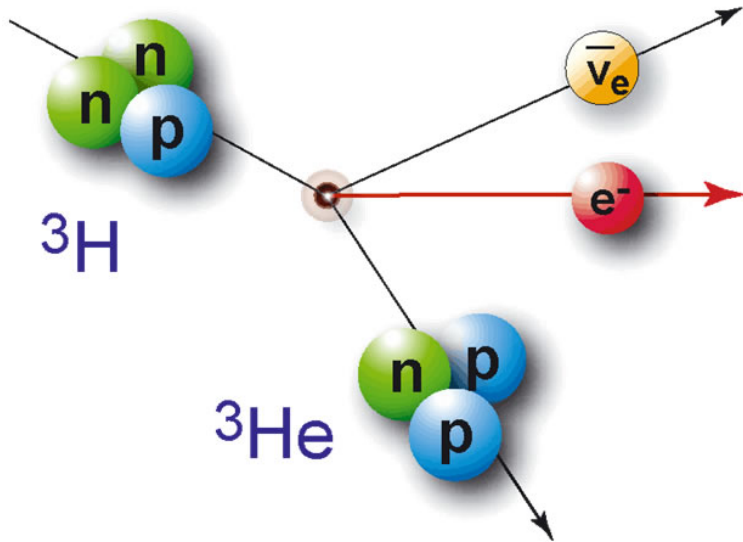
$$\langle v_{rms} \rangle = 160 \text{ km/s } (1 \text{ eV} / m_\nu)$$

Dicke, Peebles, Roll, Wilkinson (1965)

Cosmic neutrino background

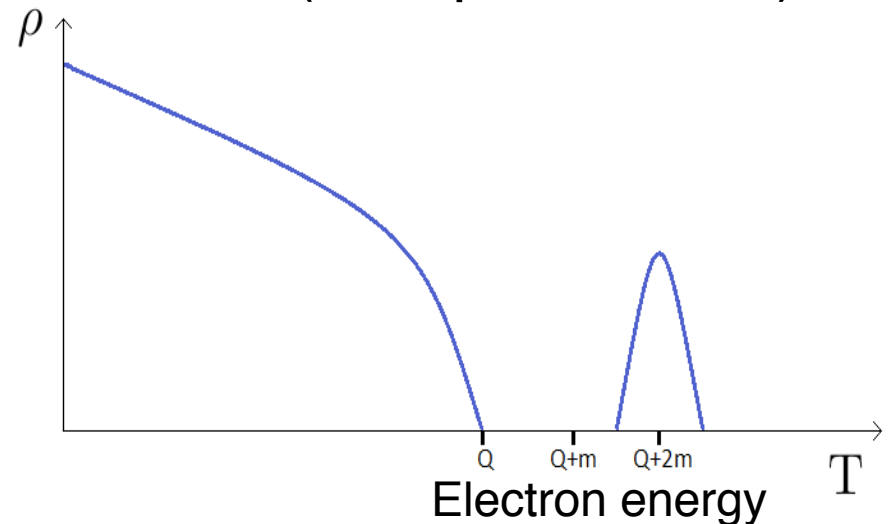


Neutrino capture



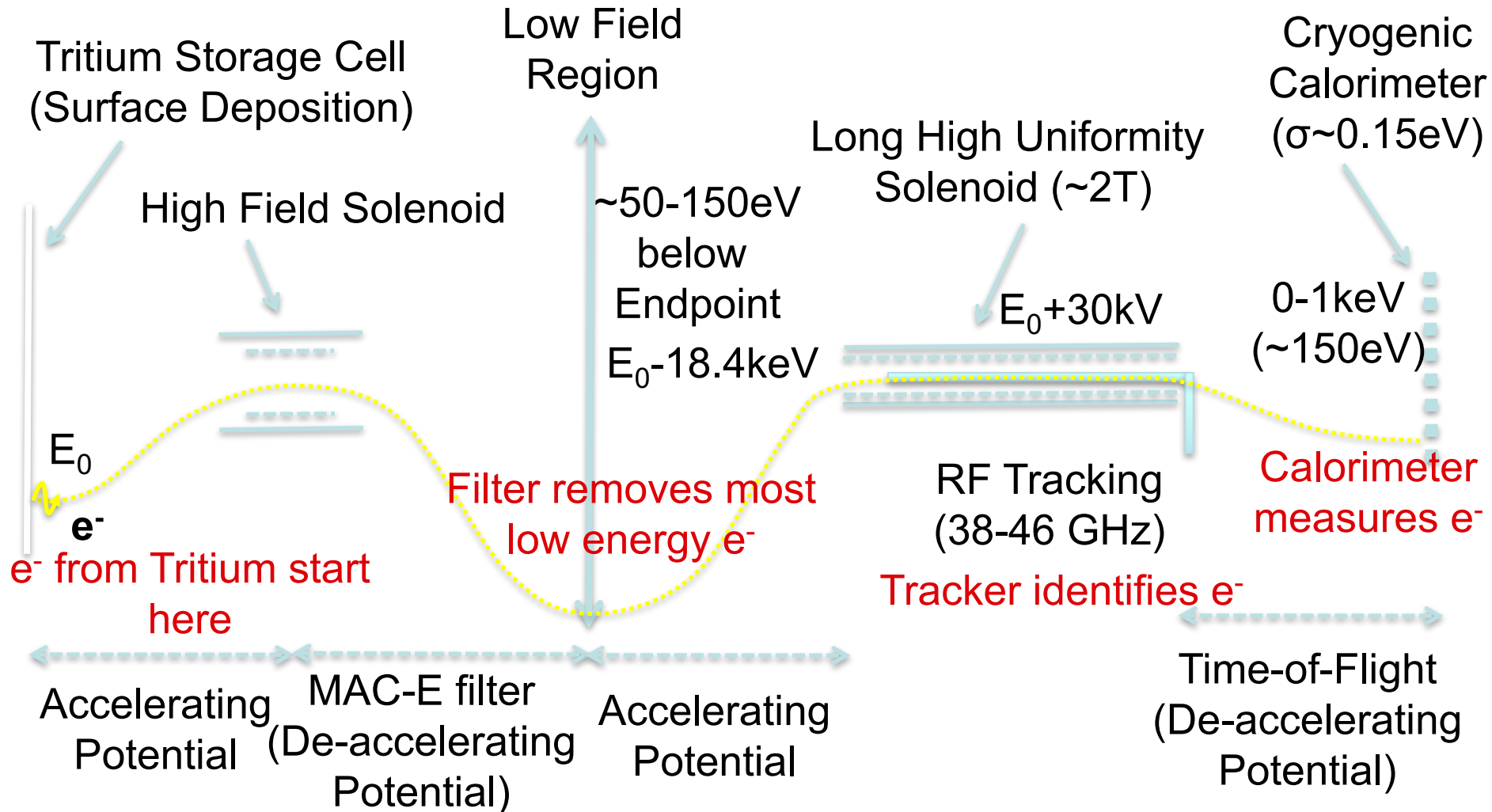
Capture cross section $\times (v/c) \sim 10^{-44} \text{ cm}^2$ (flat up to 10 keV)

Original idea: Steven Weinberg in 1962,
Phys. Rev. 128:3, 1457
JCAP 0706 (2007)015, hep-ph/0703075,
Cocco, Mangano, Messina



A little bit of everything: PTOLEMY

Princeton Tritium Observatory for Light, Early-universe, Massive-neutrino Yield





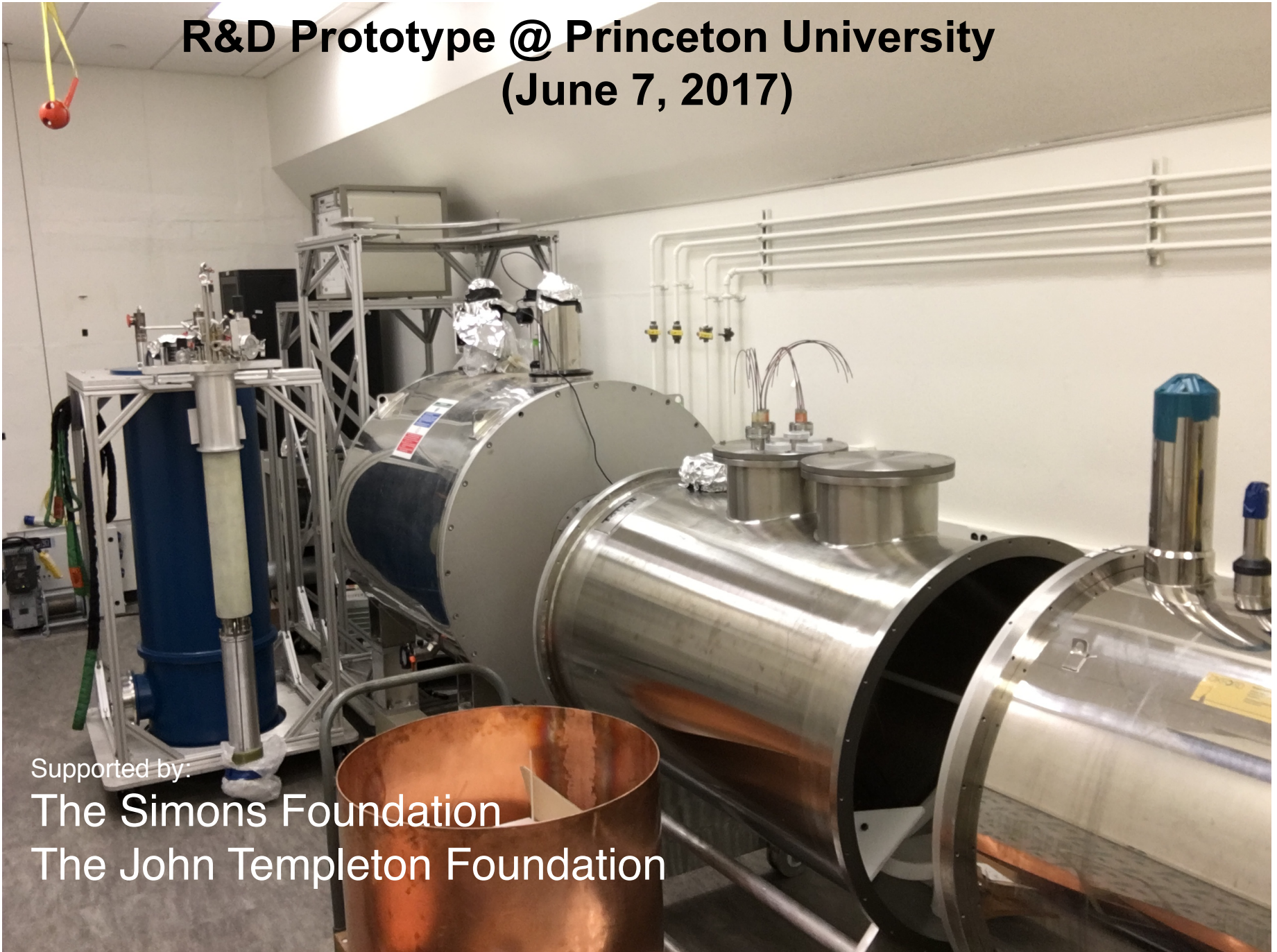
PTOLEMY

**R&D Prototype @ PPPL
(August 2, 2016)**

Supported by:
The Simons Foundation
The John Templeton Foundation

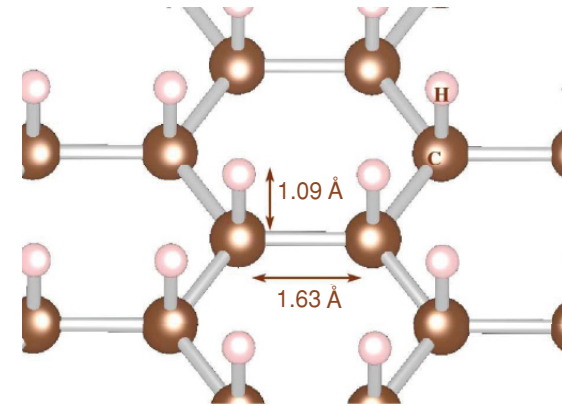
R&D Prototype @ Princeton University (June 7, 2017)

Supported by:
The Simons Foundation
The John Templeton Foundation



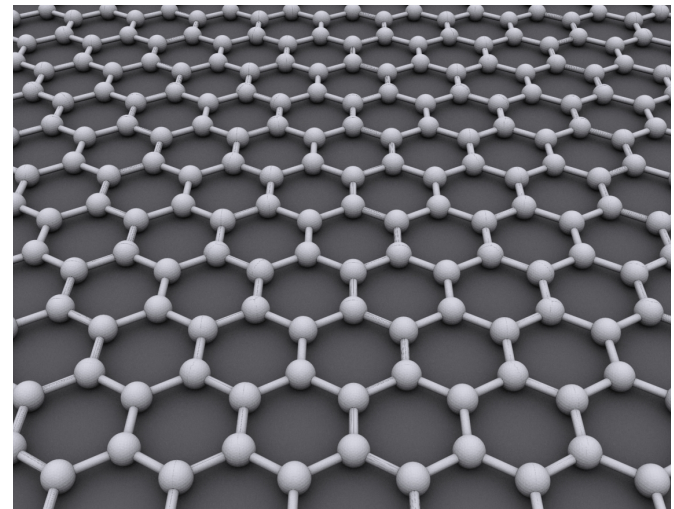
Major challenges

- Reduce molecular smearing
 - New source (Tritiated-Graphene or Cryogenic Au(111))
- Measure the energy spectrum directly with a resolution comparable to the neutrino mass
 - High-resolution electron microcalorimeters
- Compress a 70m spectrometer length – KATRIN's length – down to ~cm scale and replicate it at lower precision – final measurement from microcalorimeter
 - New filter concept (Newtonian vs. Galilean)
 - RF trigger system (Project 8 development)
 - G-FET as a potential trigger system



< 3eV binding energy

Graphene



Summary

- Beta decay measurements are the best tool we have for a precise measurement of small neutrino masses.
- KATRIN, following in the footsteps of MAINZ/TROITSK, will begin taking data soon, bringing limits on m_β to the 0.2 eV level.
- Bolometer experiments HOLMES and ECHo are also pursuing direct mass measurements in holmium.
- Radio frequency measurement developed by Project 8 is another promising technique with some recent technical success.
- These experiments can also potentially probe the CNB.
- The future of the field is dependent on nuclear theory!