

The Electron Ion Collider (EIC)

Abhay Deshpande

Lecture 2:

What do we need polarized beams?

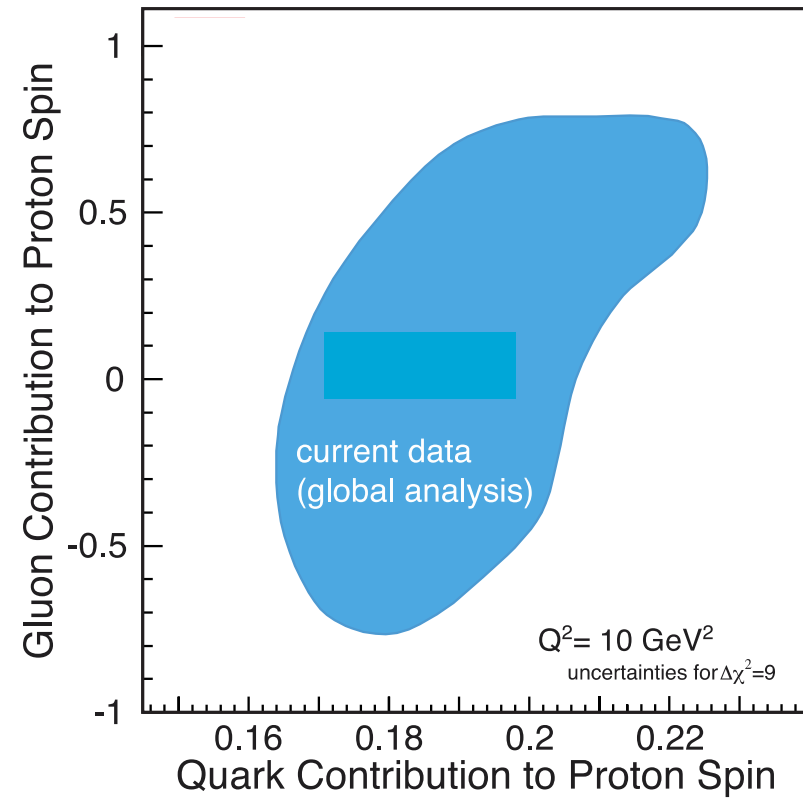
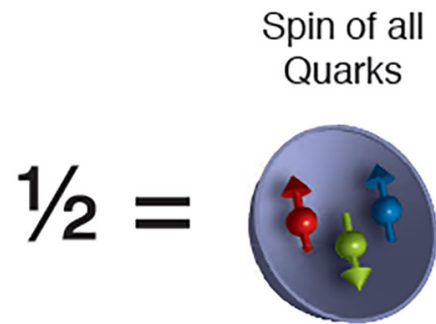
Nucleon Spin Structure: What's the problem?



Stony Brook
University

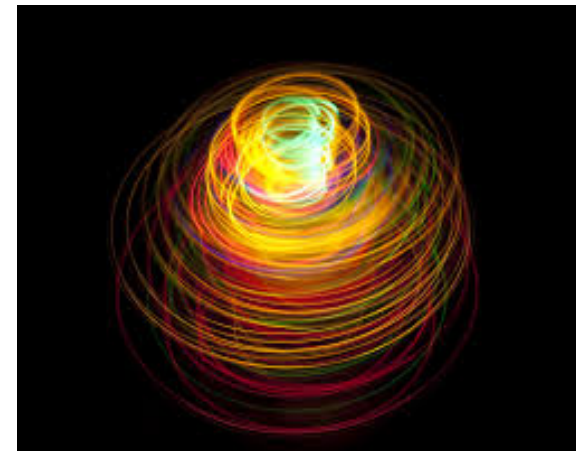
The nucleon spin puzzle....

Since 1988..



$$\frac{1}{2} = \frac{1}{2} \Delta \Sigma + \Delta G + L_Q + L_G$$

? ?



A brief history of nucleon spin: “Crisis” → “Puzzle”

Some equations...

Assume only γ^ exchange*

- Lepton Nucleon Cross Section

$$\frac{d^3\sigma}{dx dy d\phi} = \frac{\alpha^2 y}{2Q^4} L_{\mu\nu}(k, q, s,) W^{\mu\nu}(P, q, S)$$

- Lepton tensor $L_{\mu\nu}$ affects the kinematics (QED)
- Hadronic tensor $W^{\mu\nu}$ has information about the hadron structure

$$W^{\mu\nu}(P, q, S) = -\left(g^{\mu\nu} - \frac{q^\mu q^\nu}{q^2}\right) \underline{F_1(x, Q^2)} + \left(p^\mu - \frac{P \cdot q}{q^2} q^\mu\right) \left(p^\nu - \frac{P \cdot q}{q^2} q^\nu\right) \frac{1}{P \cdot q} \underline{F_2(x, Q^2)}$$

$$-i\epsilon^{\mu\nu\lambda\sigma} q_\lambda \left[\frac{M S_\sigma}{P \cdot q} \left(g_1(x, Q^2) + g_2(x, Q^2) \right) - \frac{M(S \cdot q) P_\sigma}{P \cdot q} g_2(x, Q^2) \right]$$

Lepton-nucleon Cross Section

$$\sigma = \bar{\sigma} - \frac{1}{2} h_l \delta\sigma.$$

$$\gamma = \frac{2Mx}{\sqrt{Q^2}} = \frac{\sqrt{Q^2}}{\nu}.$$

$$\frac{d^2\bar{\sigma}}{dx dQ^2} = \frac{4\pi\alpha^2}{Q^4 x} \left[xy^2 \left(1 - \frac{2m_l^2}{Q^2} \right) F_1(x, Q^2) \right. \\ \left. + \left(1 - y - \frac{\gamma^2 y^2}{4} \right) F_2(x, Q^2) \right],$$

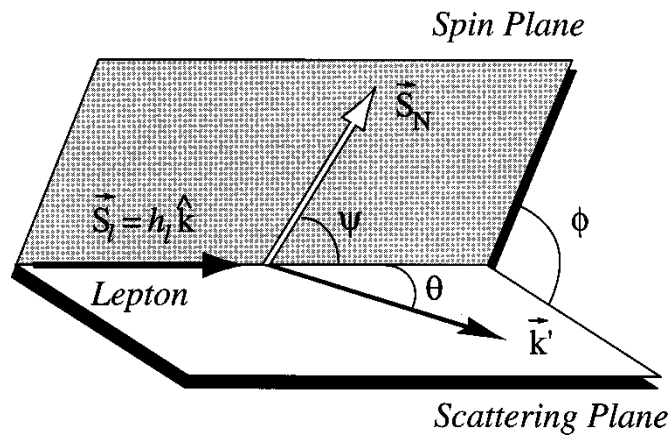
lepton helicity $h_l = \pm 1$

unpolarized structure functions $F_{1,2}(x, Q^2)$

scaling variable $x = Q^2 / 2M\nu$

exchanged virtual photon energy $= \nu$

Polarized lepton-nucleon cross section...



$$\Delta\sigma = \cos\psi \Delta\sigma_{\parallel} + \sin\psi \cos\phi \Delta\sigma_{\perp}$$

$$\gamma = \frac{2Mx}{\sqrt{Q^2}} = \frac{\sqrt{Q^2}}{\nu}$$

For high energy γ is small

$$\frac{d^2\Delta\sigma_{\parallel}}{dx dQ^2} = \frac{16\pi\alpha^2 y}{Q^4} \left[\left(1 - \frac{y}{2} - \frac{\gamma^2 y^2}{4} \right) g_1 - \frac{\gamma^2 y}{2} g_2 \right]$$

$$\frac{d^3\Delta\sigma_T}{dx dQ^2 d\phi} = -\cos\phi \frac{8\alpha^2 y}{Q^4} \gamma \sqrt{1 - y - \frac{\gamma^2 y^2}{4}} \left(\frac{y}{2} g_1 + g_2 \right)$$

Relation to spin structure function g_1

$$g_1(x) = \frac{1}{2} \sum_{i=1}^{n_f} e_i^2 \Delta q_i(x)$$

$$\Delta q_i(x) = q_i^+(x) - q_i^-(x) + \bar{q}_i^+(x) - \bar{q}_i^-(x)$$

$$q_i^+ \quad (\bar{q}_i^+) \quad \text{and} \quad q_i^- \quad (\bar{q}_i^-)$$

Quark and anti-quark with spin orientation along and against the proton spin.

- In QCD quarks interact with each other through gluons, which gives rise to a weak Q^2 dependence of structure functions
- At any given Q^2 the spin structure function is related to polarized quark & gluon distributions by coefficients C_q and C_g

Composition & Q^2 or t dependence of Structure Functions

$$g_1(x,t) = \frac{1}{2} \sum_{k=1}^{n_f} \frac{e_k^2}{n_f} \int_x^1 \frac{dy}{y} \left[C_q^S\left(\frac{x}{y}, \alpha_s(t)\right) \Delta\Sigma(y,t) \right. \\ \left. + 2n_f C_g\left(\frac{x}{y}, \alpha_s(t)\right) \Delta g(y,t) \right. \\ \left. + C_q^{NS}\left(\frac{x}{y}, \alpha_s(t)\right) \Delta q^{NS}(y,t) \right].$$

In this equation:

$$t = \ln(Q^2/\Lambda^2)$$

α_s = strong interaction constant
S & NS stand for flavor singlet & flavor non-singlet

$$\Delta\Sigma(x,t) = \sum_{i=1}^{n_f} \Delta q_i(x,t),$$

$$\Delta q^{NS}(x,t) = \left[\sum_{i=1}^{n_f} \left(e_i^2 - \frac{1}{n_f} \sum_{k=1}^{n_f} e_k^2 \right) / \frac{1}{n_f} \sum_{k=1}^{n_f} e_k^2 \right] \Delta q_i(x,t).$$

To study the Q^2 evolution experimentally: you need a wide Q^2 range in measurements

Spin Crisis → Puzzle (1960-2000)

Limitations of the Quark Parton Model (QPM) & lessons learnt

Early (fixed target) spin experiments and their limitations

@1960's # of hadrons > # of chemical elements

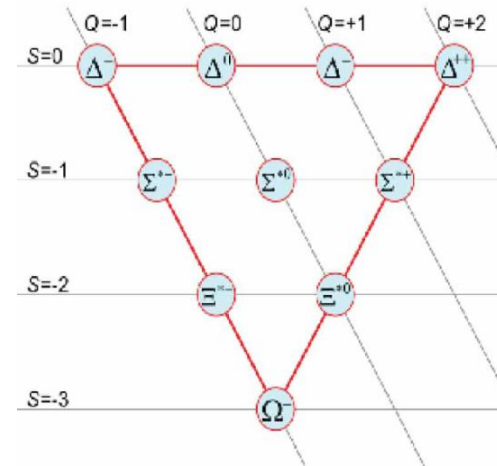
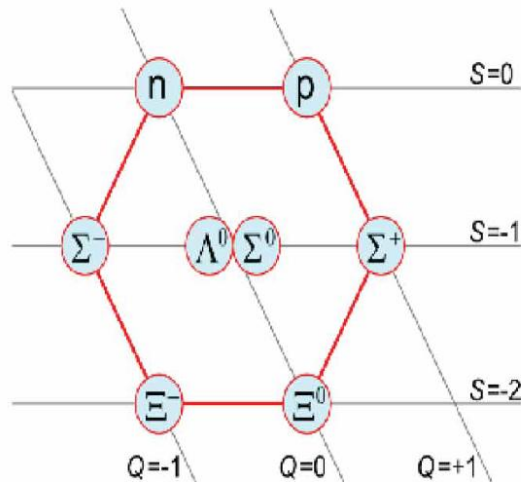
*"If I had known this earlier,
I would have done Biology"*
--W. Pauli

- baryons ($J = n/2$)_{n=1,2,...}
- mesons: ($J = n$)_{n=0,1,...}

1961 'The Eightfold way' : all baryons and mesons grouped in multiplets defined by SU(3) symmetry

e.g., baryons:

$J^P = 1/2^+$
(same spin and parity)



$J^P = 3/2^+$

isolines of same charge and strangeness

Understanding the proton structure:

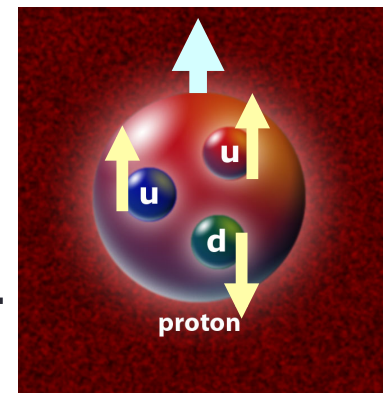
Friedman, Kendall, Taylor: 1960's SLAC Experiment

1990 Nobel Prize: *"for their pioneering investigations concerning **deep inelastic scattering of electrons on protons and bound neutrons**, which have been of essential importance for the development of the **quark model** in particle physics".*

Obvious next Question:

Could we understand other properties of proton, e.g. SPIN, in the quark-parton model?

Proton Spin = $\frac{1}{2}$, each quark is a spin $\frac{1}{2}$ particle...



Structure Functions & PDFs

- The F_1 and F_2 are unpolarized structure functions or momentum distributions
- The g_1 and g_2 are polarized structure functions or spin distributions
- In QPM
 - $F_2(x) = 2xF_1$ (Callan Gross relation)
 - $g_2 = 0$ (Twist 3 quark gluon correlations)

$$F_1(x) = \frac{1}{2} \sum_f e_f^2 \{q_f^+(x) + q_f^-(x)\} = \frac{1}{2} \sum_f e_f^2 q_f(x)$$

$$g_1(x) = \frac{1}{2} \sum_f e_f^2 \{q_f^+(x) - q_f^-(x)\} = \frac{1}{2} \sum_f e_f^2 \Delta q_f(x)$$



Remember
this

Nucleon spin & Quark Probabilities

- Define

$$\Delta q = q^+ - q^-$$

- With q^+ and q^- probabilities of quark & anti-quark with spin parallel and anti-parallel to the nucleon spin
- Total quark contribution then can be written as:

$$\Delta\Sigma = \Delta u + \Delta d + \Delta s$$

- The nucleon spin composition

$$\frac{1}{2} = \frac{1}{2}\Delta\Sigma$$

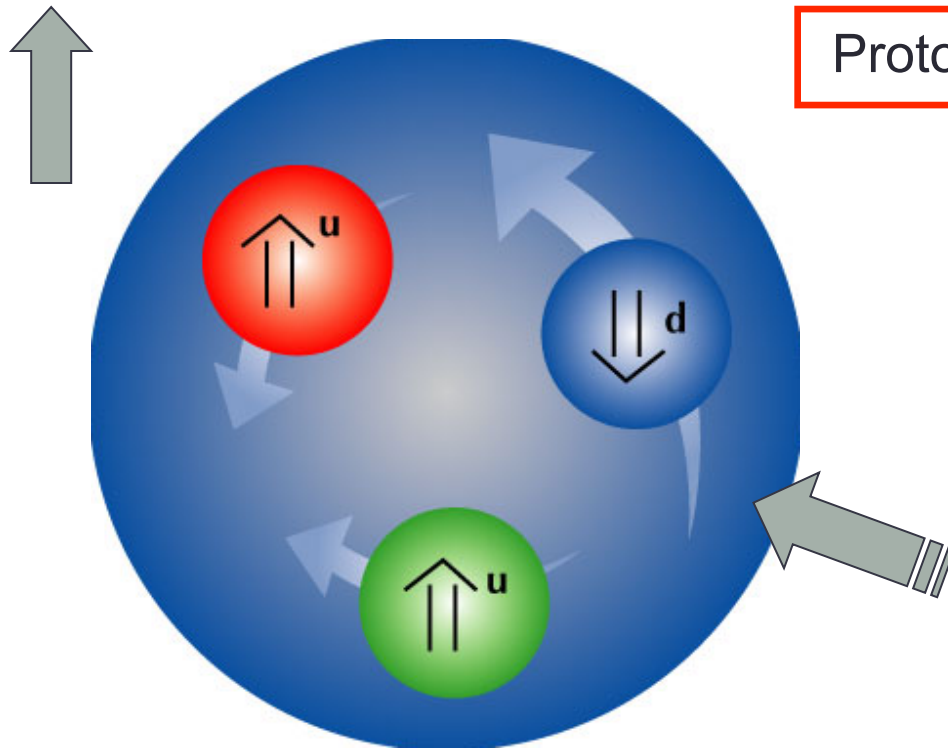
Nucleon's Spin: Naïve Quark Parton Model

- Protons and Neutrons are spin 1/2 particles
- Quarks that constitute them are also spin 1/2 particles

• And there are three of them in the

Proton: u u d

Neutron: u d d



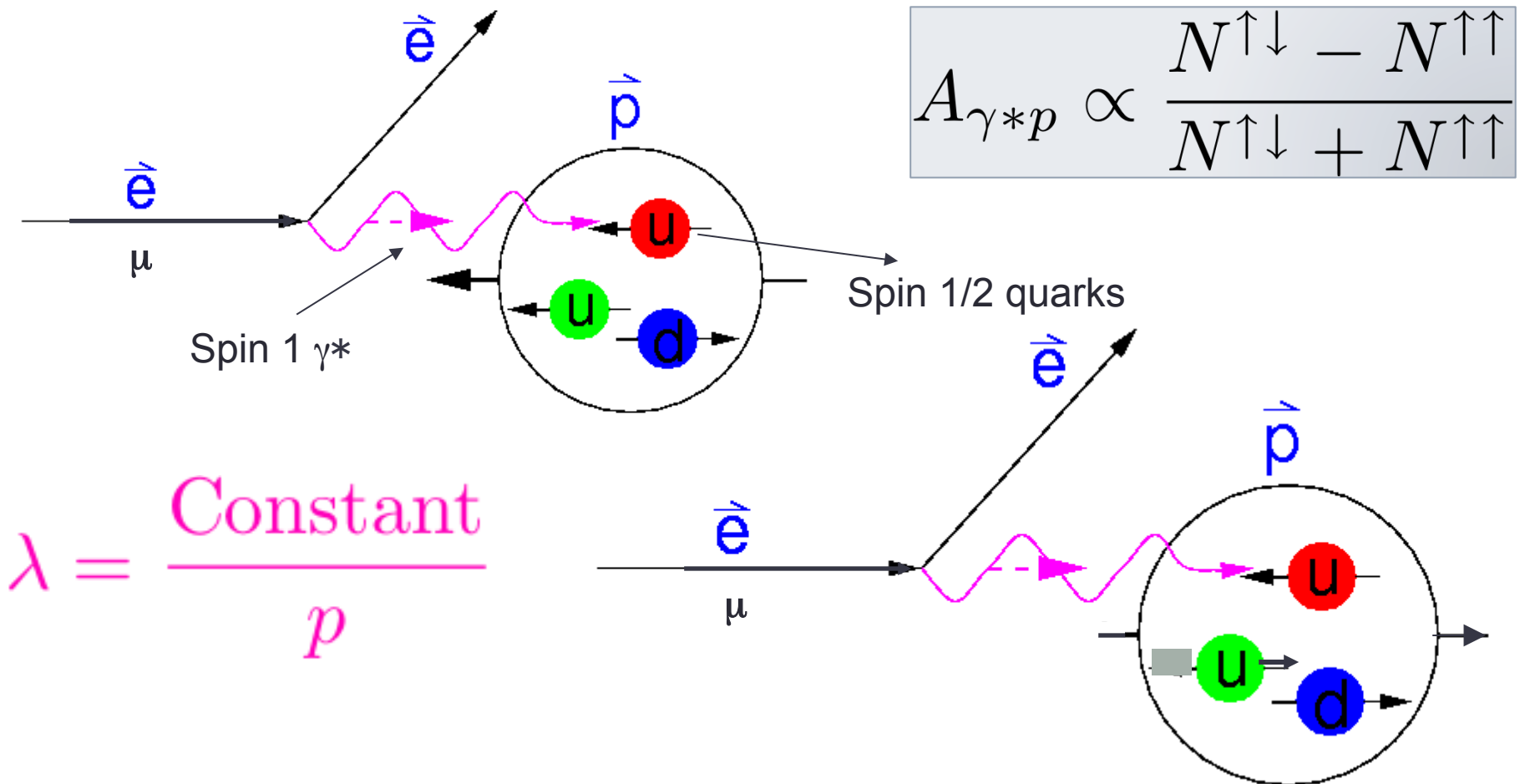
$S_{\text{proton}} = \text{Sum of all quark spins!}$

$$1/2 \quad ? = 1/2 + 1/2 + 1/2$$

$$1/2 = 1/2 - 1/2 + 1/2$$

How was the Quark Spin measured?

- Deep Inelastic polarized electron or muon scattering



Measurements of spin structure functions:
What issues we need to worry about?

- 1) Design of experiments, operational issues
- 2) Calculations of spin structure functions

Experimental Needs in DIS

Polarized target, polarized beam

- Polarized targets: hydrogen (p), deuteron (pn), helium (^3He : 2p+n)
- Polarized beams: electron, muon used in DIS experiments

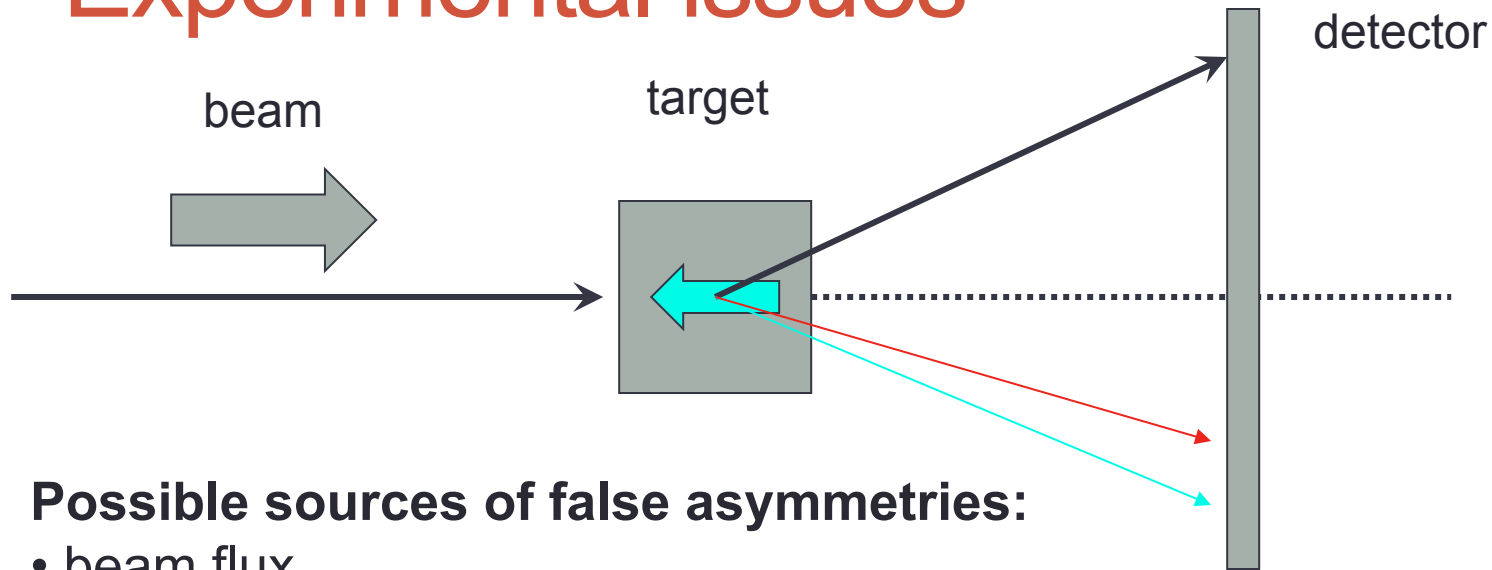
Determine the kinematics: measure with high accuracy:

- Energy of **incoming lepton**
- Energy, direction of **scattered lepton**: energy, direction
- Good identification of **scattered lepton**

Control of false asymmetries:

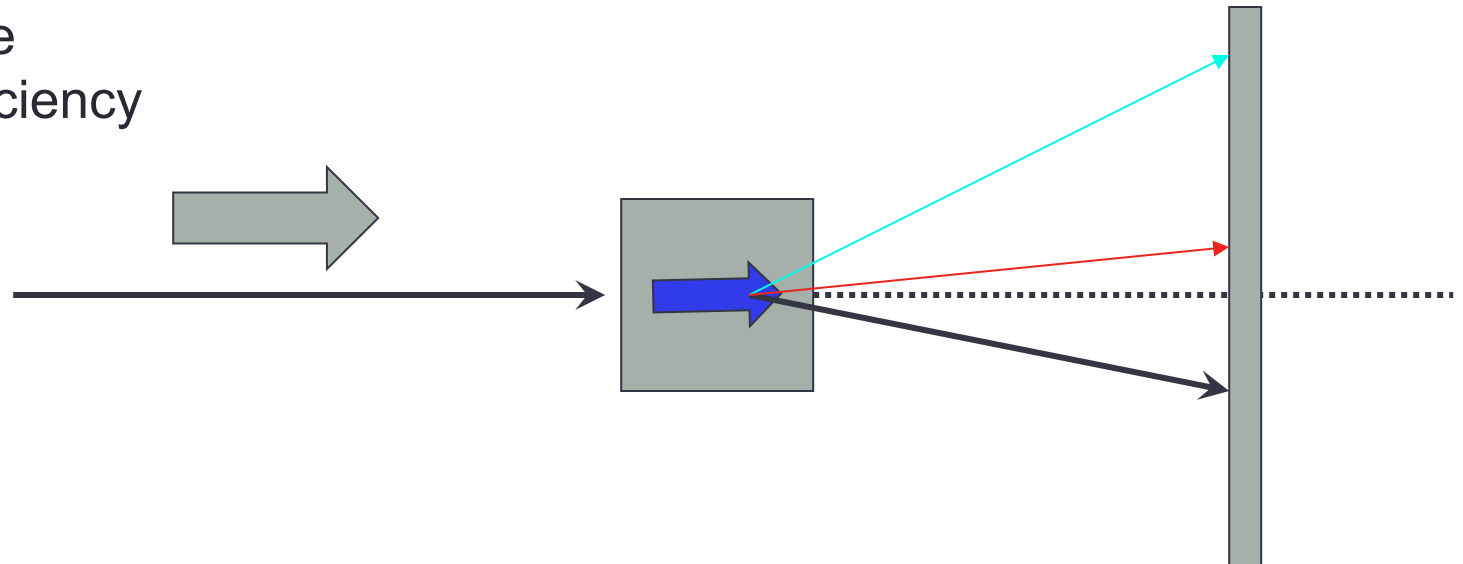
- **Need excellent understanding and control of false asymmetries (time variation of the detector efficiency etc.)**

Experimental issues



Possible sources of false asymmetries:

- beam flux
- target size
- detector size
- detector efficiency



An Ideal Situation

$$A_{measured} = \frac{N^{\rightarrow\leftarrow} - N^{\rightarrow\rightarrow}}{N^{\rightarrow\leftarrow} + N^{\rightarrow\rightarrow}}$$

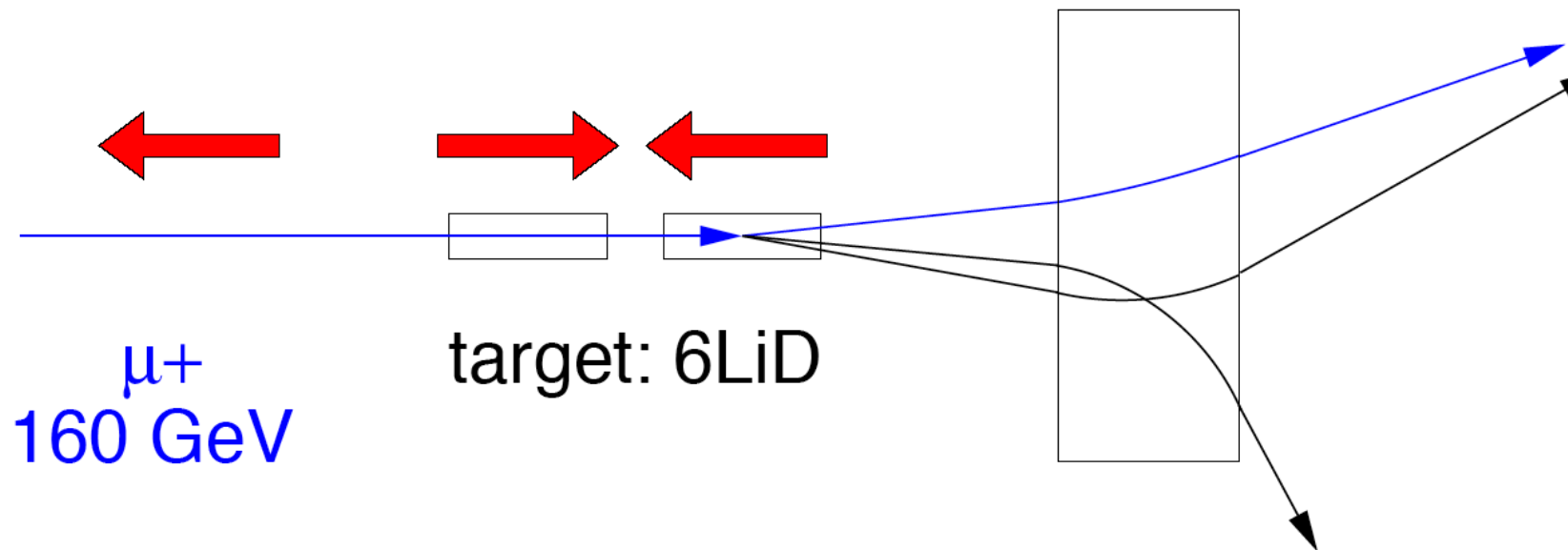
$$N^{\leftarrow\rightarrow} = N_b \cdot N_t \cdot \sigma^{\leftarrow\rightarrow} \cdot D_{acc} \cdot D_{eff}$$

$$N^{\rightarrow\rightarrow} = N_b \cdot N_t \cdot \sigma^{\rightarrow\rightarrow} \cdot D_{acc} \cdot D_{eff}$$

If all other things are equal, they cancel in the ratio and....

$$A_{measured} = \frac{\sigma^{\rightarrow\leftarrow} - \sigma^{\rightarrow\rightarrow}}{\sigma^{\rightarrow\leftarrow} + \sigma^{\rightarrow\rightarrow}}$$

A Typical Setup



- Target polarization direction reversed every 6-8 hrs
- Typically experiments try to limit false asymmetries to be about 10 times smaller than the physics asymmetry of interest

Asymmetry Measurement

$$\frac{N^{\uparrow\downarrow} - N^{\uparrow\uparrow}}{N^{\uparrow\downarrow} + N^{\uparrow\uparrow}} = A_{measured} = P_{beam} \cdot P_{target} \cdot f \cdot A_{\parallel}$$

- f = dilution factor proportional to the polarizable nucleons of interest in the target “material” used, for example for NH_3 , $f=3/17$

$$g_1 \approx \frac{A_{\parallel}}{D} \cdot F_1 \approx \frac{A_{\parallel}}{D} \frac{F_2}{2 \cdot x} \quad \int_0^1 g_1^p(x, Q_0^2) dx = \Gamma_1^p(Q_0^2)$$

- D is the depolarization factor, kinematics, polarization transfer from polarized lepton to photon, $D \sim y^2$

First Moments of SPIN SFs

- With $\Delta q = \int \Delta q(x) dx$

$$g_1(x) = \frac{1}{2} \sum_f e_f^2 \{q_f^+(x) - q_f^-(x)\} = \frac{1}{2} \sum_f e_f^2 \Delta q_f(x)$$

$$\Gamma_1^p = \frac{1}{2} \left[\frac{4}{9} \Delta u + \frac{1}{9} \Delta d + \frac{1}{9} \Delta s \right]$$

$$= \frac{1}{12} \underbrace{(\Delta u - \Delta d)}_{a_3 = g_a} + \frac{1}{36} \underbrace{(\Delta u + \Delta d - 2\Delta s)}_{a_8} + \frac{1}{9} \underbrace{(\Delta u + \Delta d + \Delta s)}_{a_0}$$

Neutron decay
(3F-D)/3
Hyperon Decay

$\Delta\Sigma$

$$\Gamma_1^{p,n} = \frac{1}{12} \left[\pm a_3 + \frac{1}{\sqrt{3}} a_8 \right] + \frac{1}{9} a_0$$

First moment of $g_1^p(x)$: Ellis-Jaffe SR

$$\Gamma_1^{p,n} = \frac{1}{12} \left[\pm a_3 + \frac{1}{\sqrt{3}} a_8 \right] + \frac{1}{9} a_0$$

$$a_3 = \frac{g_A}{g_V} = F + D = 1.2601 \pm 0.0025$$

$$a_8 = 3F - D \implies F/D = 0.575 \pm 0.016$$

Assuming $SU(3)_f$ & $\Delta s = 0$, Ellis & Jaffe: $\Gamma_1^p = 0.170 \pm 0.004$

Measurements were done at SLAC (E80, E130) Experiments:

Low 8-20 GeV electron beam on fixed target

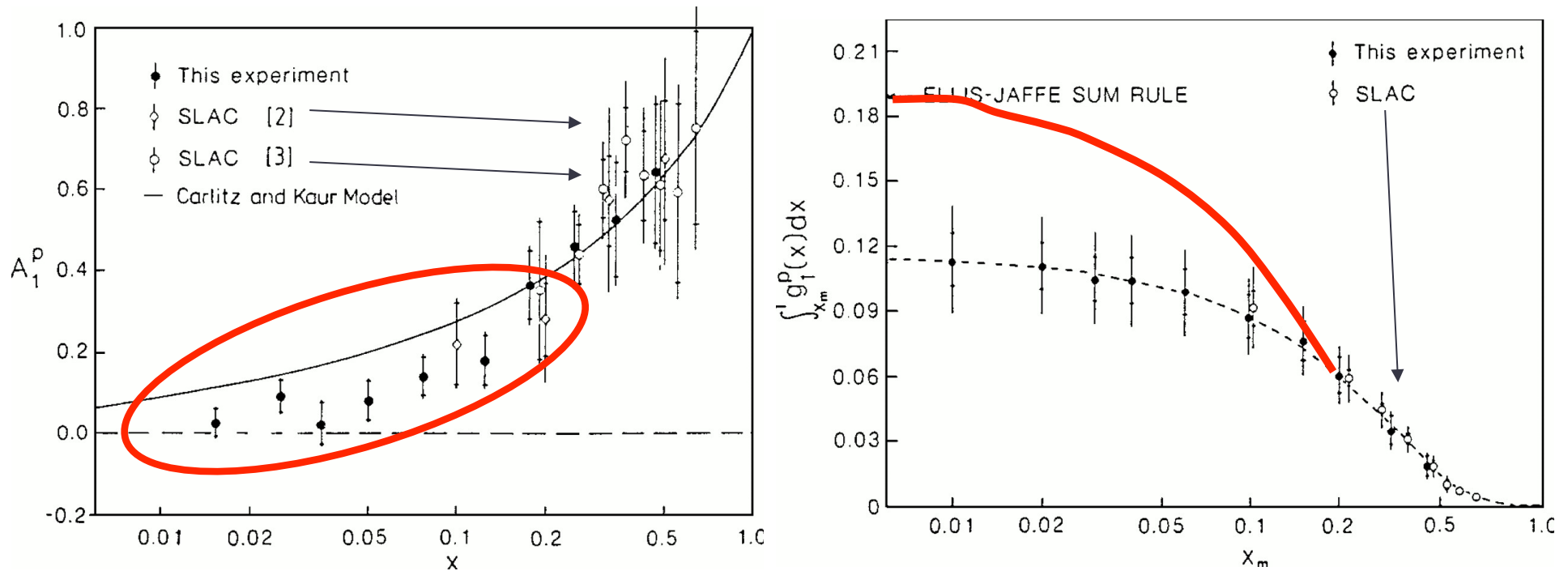
Did not reach low enough $x \rightarrow x_{\min} \sim 10^{-2}$

Found consistency of data and E-J sum rule above

European Muon Collaboration at CERN

- 160 GeV muon beam (lower intensity), but significantly higher energy
- Significantly LOWER X reach $\rightarrow x_{\min} \sim 10^{-3}$
- Polarized target
- Repeated experiment for A_1 and measured g_1 of the proton!

Proton Spin Crisis (1989)!

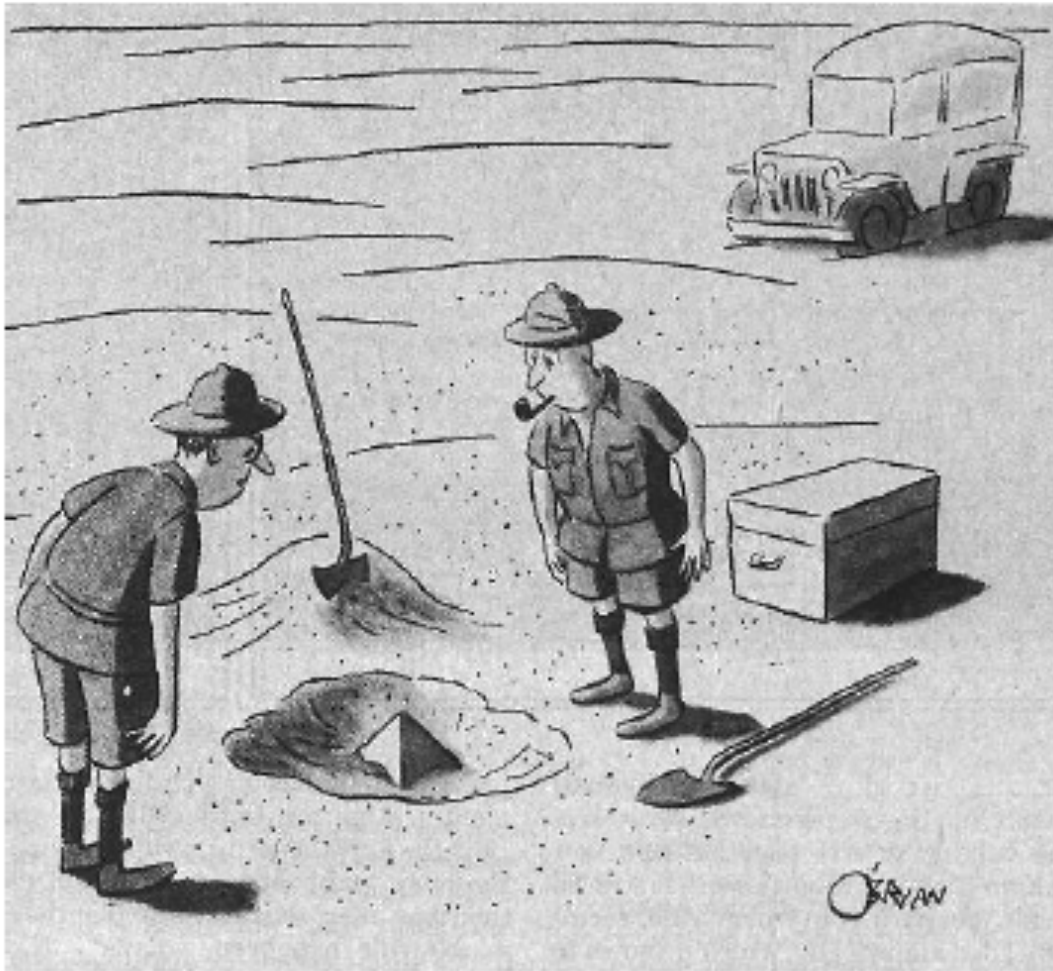


$$\Delta\Sigma = (0.12) \pm (0.17) \text{ (EMC, 1989)}$$

Ashman et al., EMC Collaboration, NPB 328 (1989) I

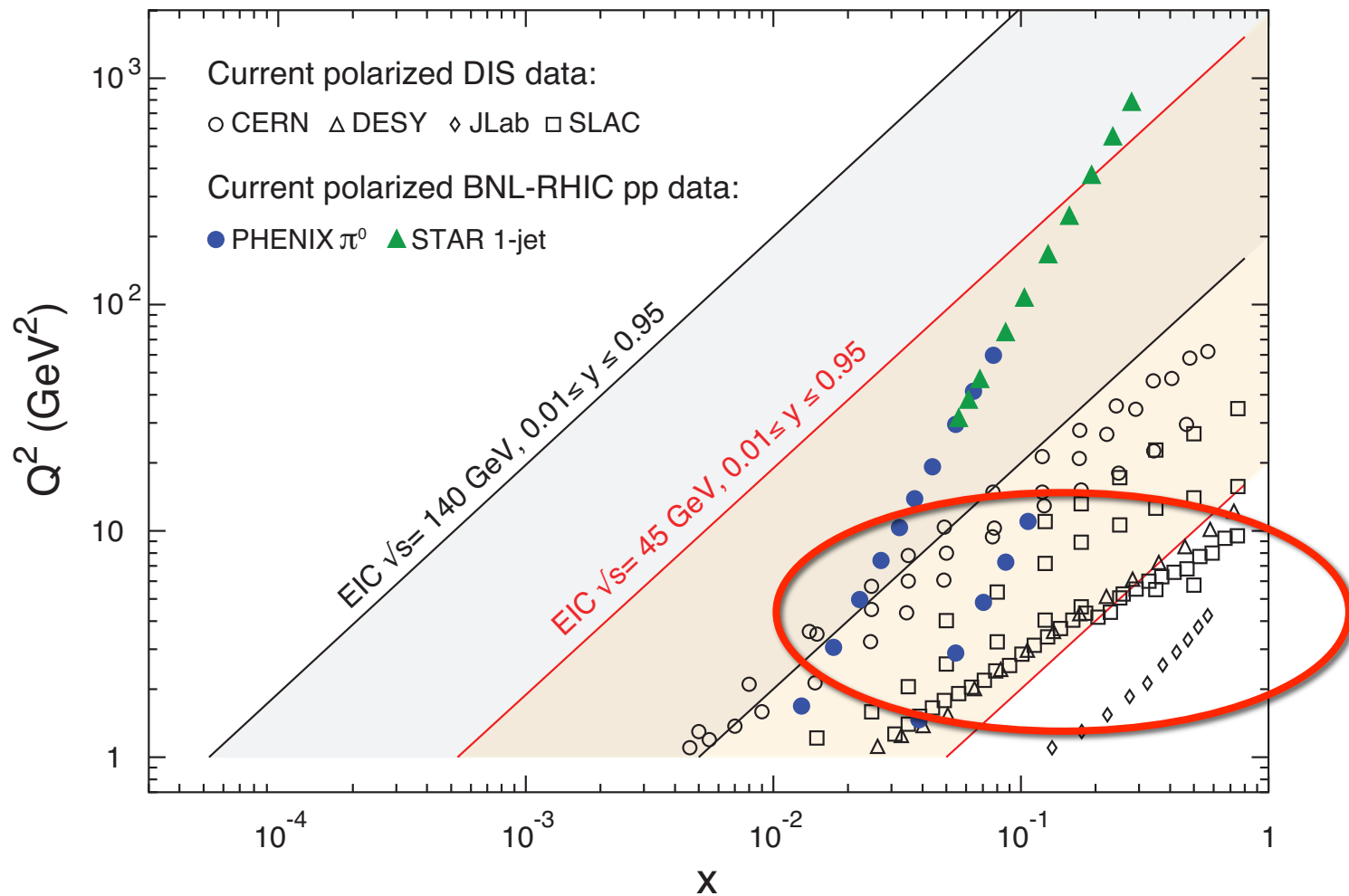
$\Delta\Sigma = 0.58$ expected from E-J sum rule....

How significant is this?



“It could be the discovery of the century. Depending, of course on how far below it goes...”

Fixed target experiments:



Extrapolations!

The most simplistic but intuitive theoretical predictions for the polarized deep inelastic scattering are the **sum rules** for the nucleon structure function g_1 .

$$\Gamma_1(Q^2) = \int_0^1 g_1(x, Q^2) dx$$

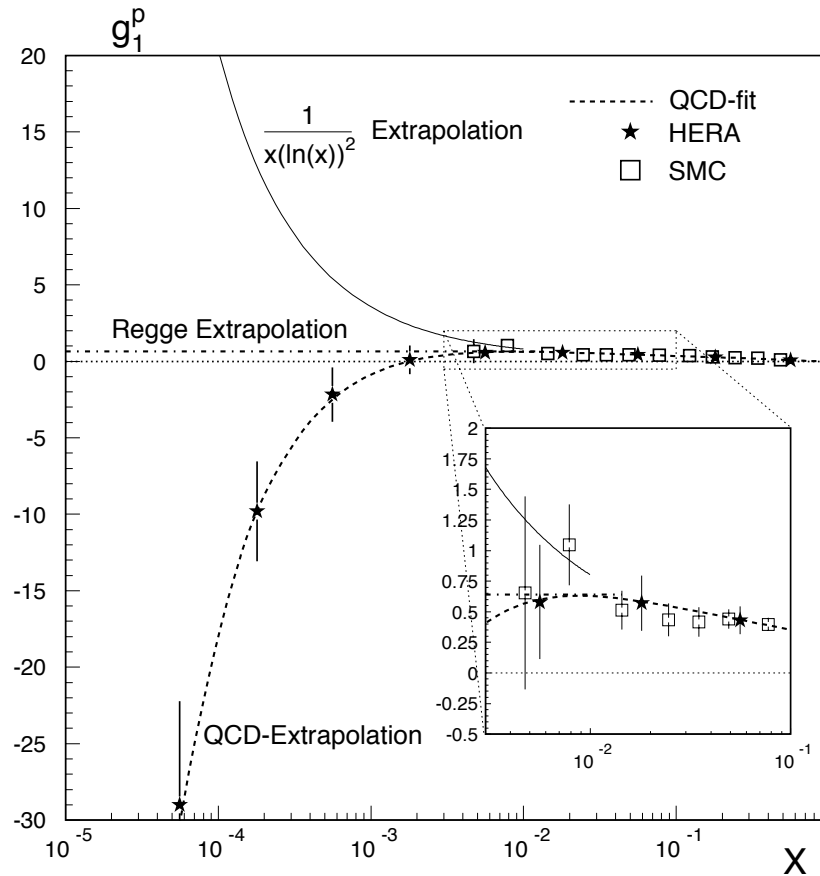
Due to experimental limitations, accessibility of x range is limited, and extrapolations to $x=0$ and $x=1$ are **unavoidable**.

Extrapolations to $x=1$, are *somewhat* less problematic: $|A_1| \leq 1$
Small contribution to the integral

Future precision studies at JLab 12GeV of great interest

Low x behavior of $g_1(x)$ is theoretically not well established
hence of **significant debate and excitement** in the community

A collection of low x behaviors:



Deshpande, Hughes, Lichtenstadt, HERA low x WS (1999)
 Simulated data for polarized e-p scattering shown in the figure.

- Low x behavior all over the place
- No theoretical guidance for which one is correct
- Only logical path is through measurements.
 - Polarized HERA not easy
 - **Now being considered in view of the Electron Ion Collider**

A. De Roeck, A. Deshpande, V. Hughes, J. Lichtenstadt, G. Radcliff, EPJ, C6 (1999), 121

Aftermath of the EMC Spin “Crisis”

Naïve quark model yields: $\Delta u = 4/3$ and $\Delta d = -1/3 \implies \Delta\Sigma = 1$

Relativistic effects included quark model: $\Delta\Sigma = 0.6$

After much discussions, arguments an idea that became emergent, although not without controversy: “gluon anomaly”

- True quark spin is screened by large gluon spin: Altarelli, Ross
Carlitz, Collins
Mueller et al.

$$\Delta\Sigma(Q^2) = \Delta\Sigma' - N_f \frac{\alpha_S(Q^2)}{2\pi} \Delta g(Q^2)$$

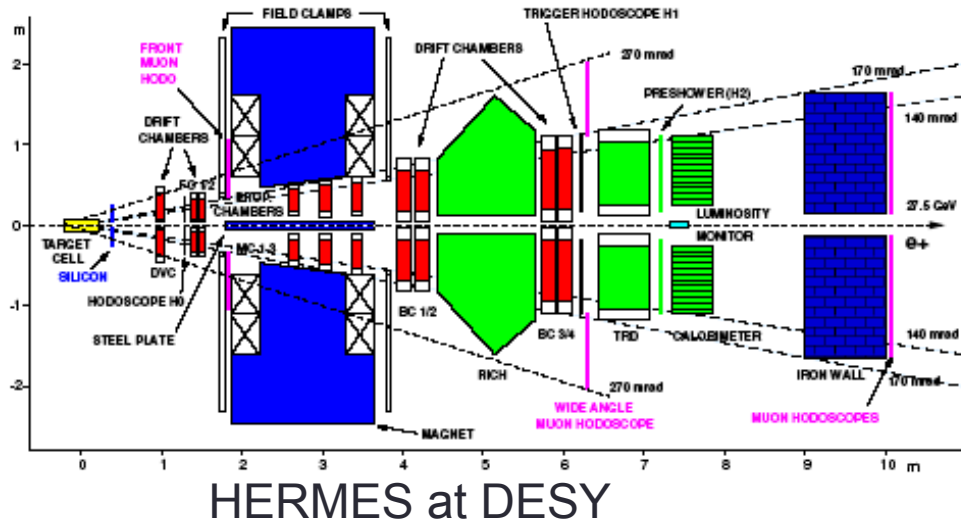
- But there were strong alternative scenarios proposed that blamed the remaining spin of the proton on:

- Gluon spin (same as above)
- Orbital motion of quarks and gluons (OAM)

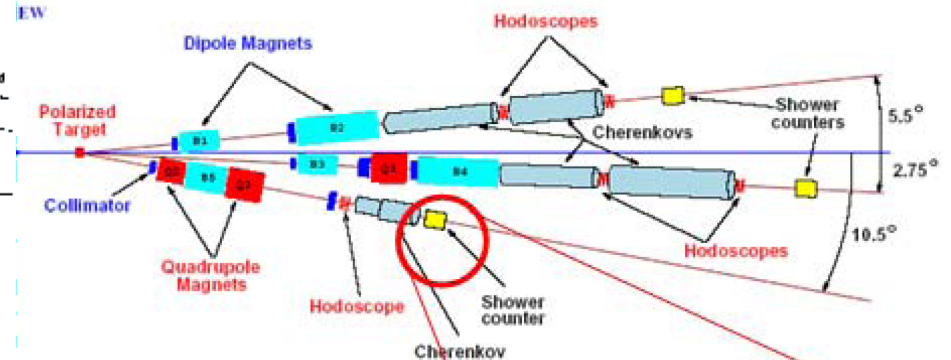
Jaffe, Manohar
Ji et al

It became clear that precision measurements of nucleon spin constitution was needed!

Experiments



HERMES at DESY



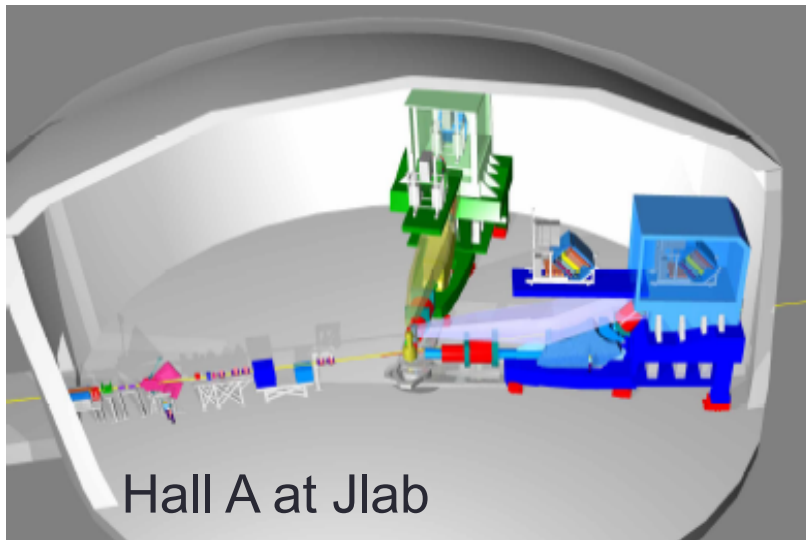
- high energy beams
- large angular acceptance
- broad kinematical range

- two stages spectrometer
- Large Angle Spectrometer (SM1)
- Small Angle Spectrometer (SM2)

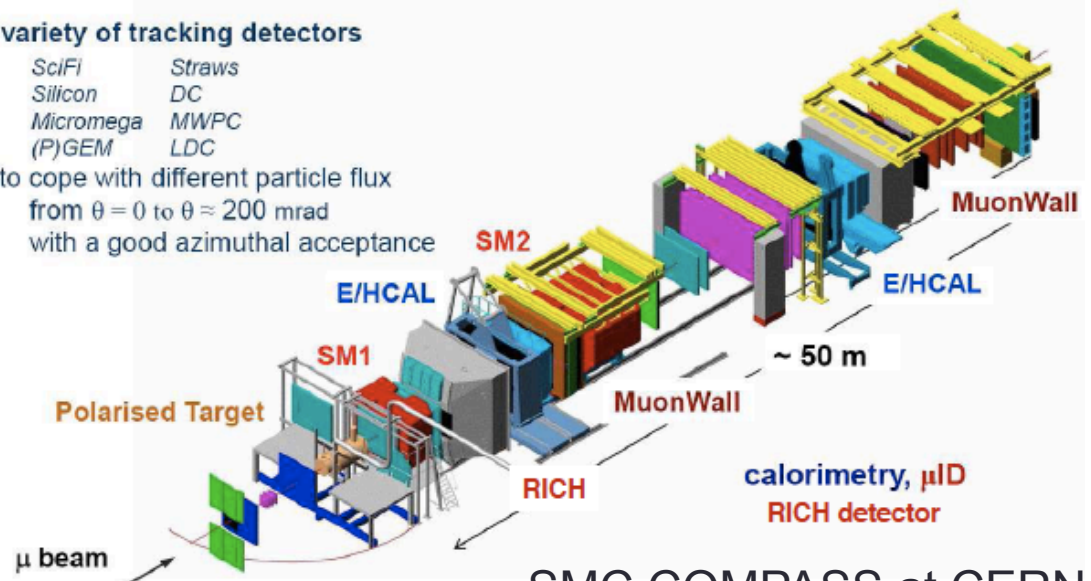
variety of tracking detectors

ScIFI	Straws
Silicon	DC
Micromegas	MWPC
(P)GEM	LDC

to cope with different particle flux
from $\theta = 0$ to $\theta \approx 200$ mrad
with a good azimuthal acceptance

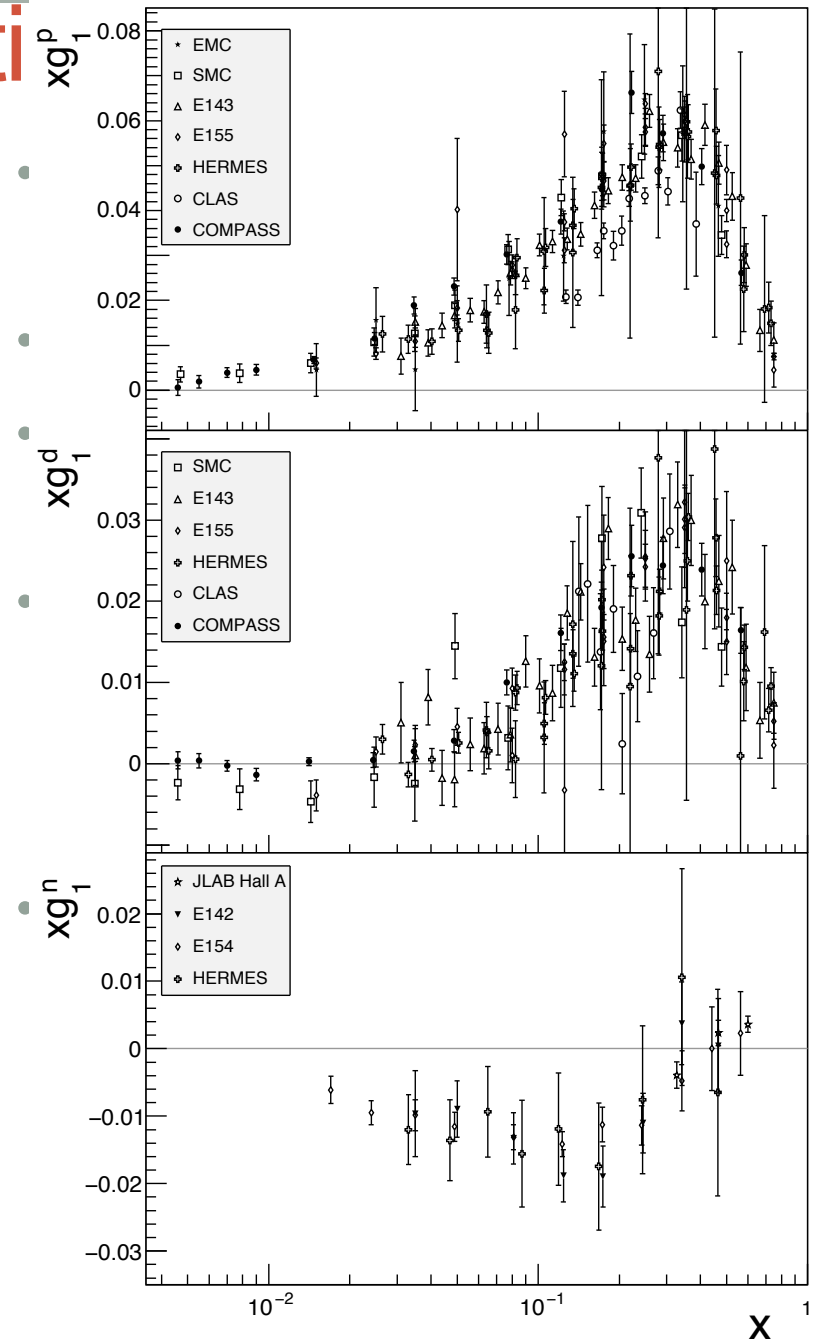
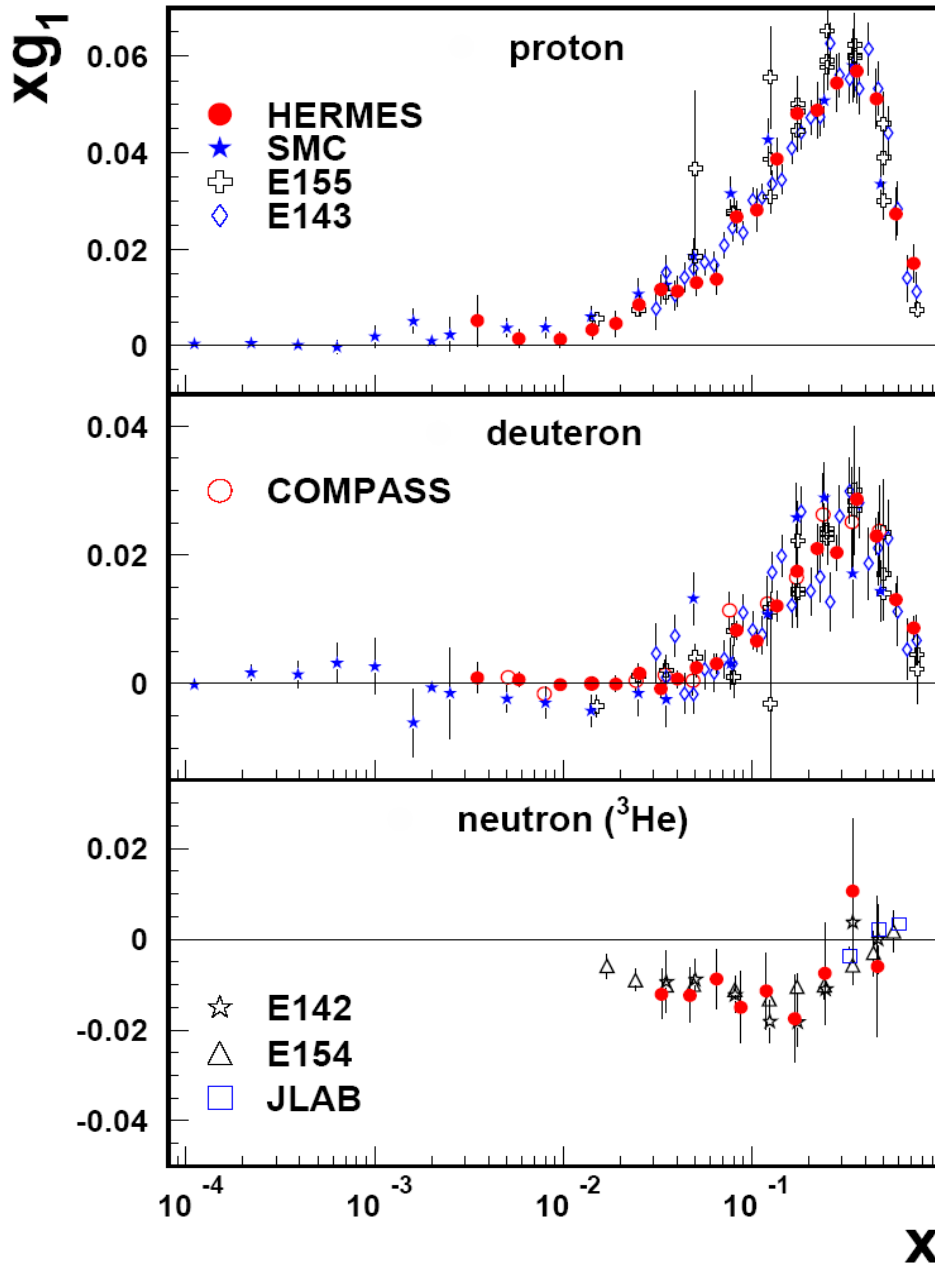


Hall A at Jlab



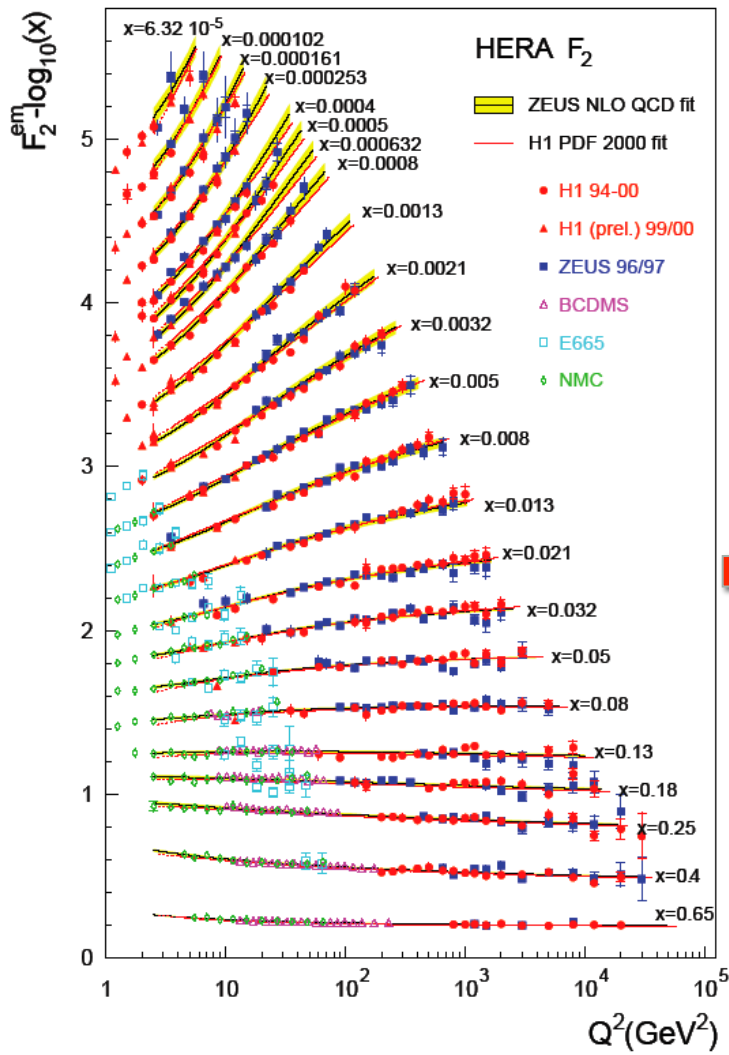
SMC, COMPASS at CERN

Spin structure Functi

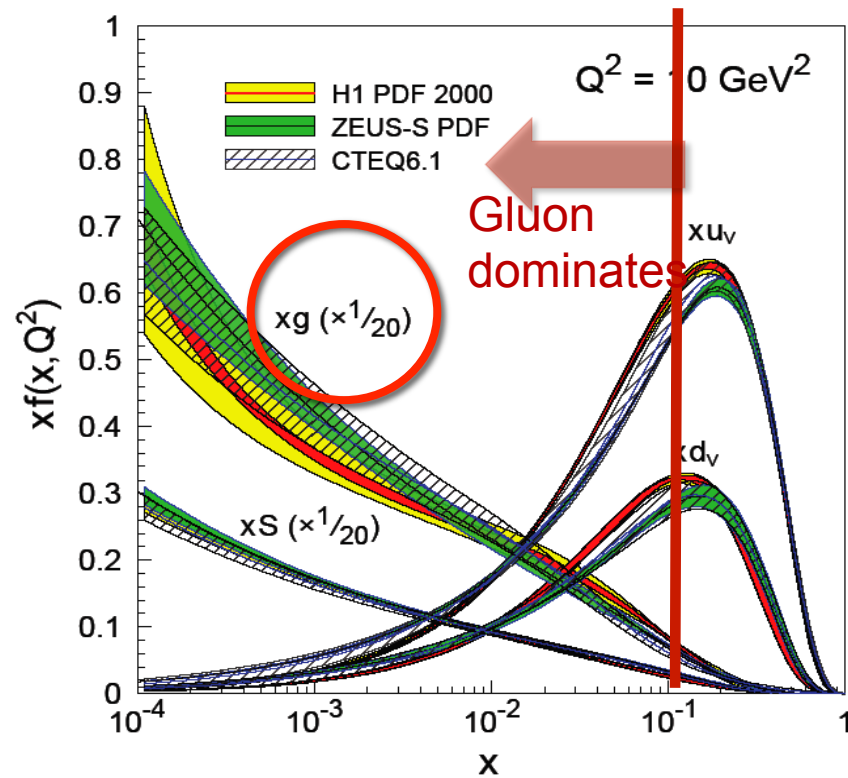


How to extract the polarized gluon distribution?

Similar to extraction of PDFs at HERA (RECALL)



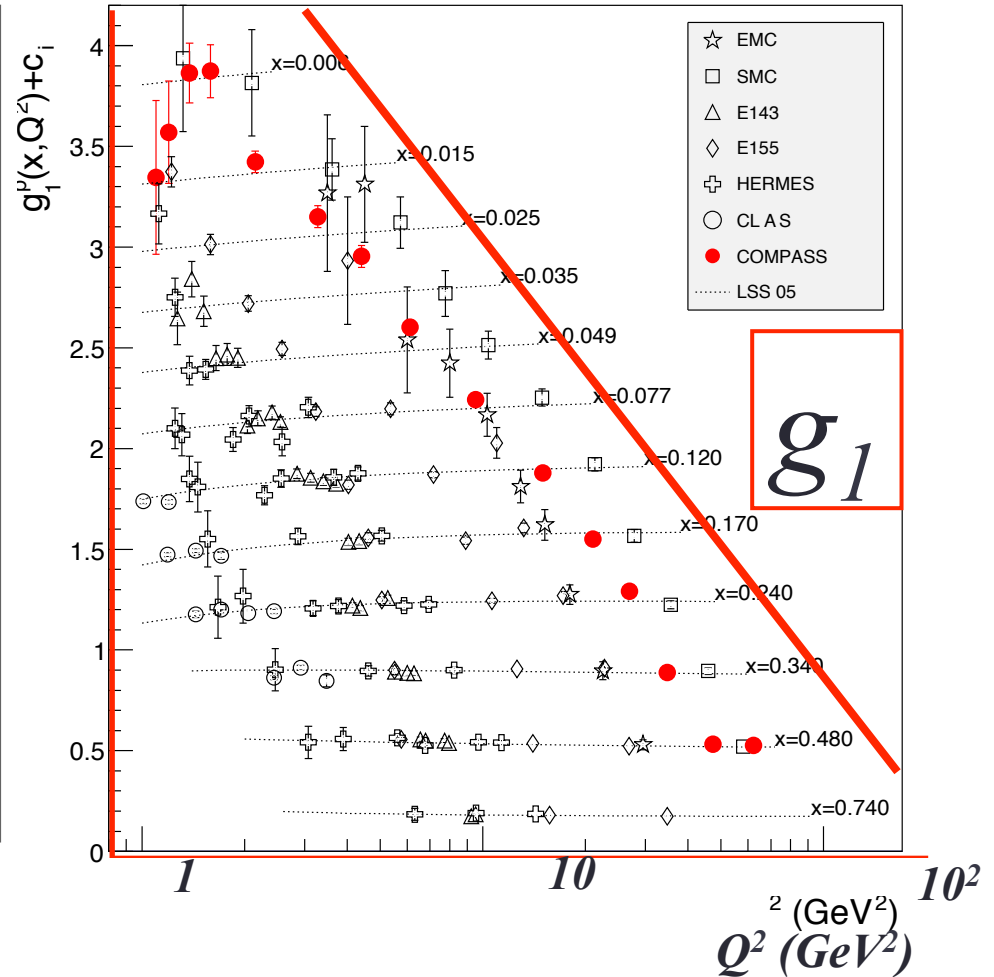
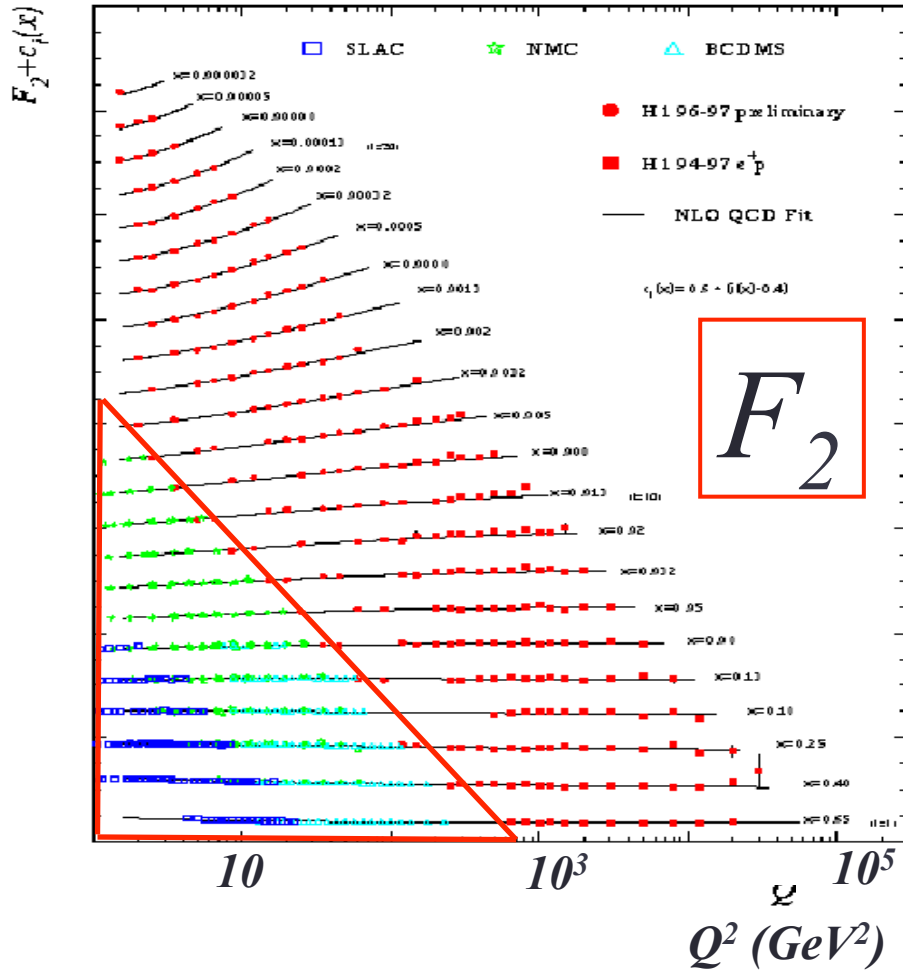
NLO pQCD analyses: fits with linear DGLAP* equations



*Dokshitzer, Gribov, Lipatov, Altarelli, Parisi

F_2 vs. g_1 structure function measurements

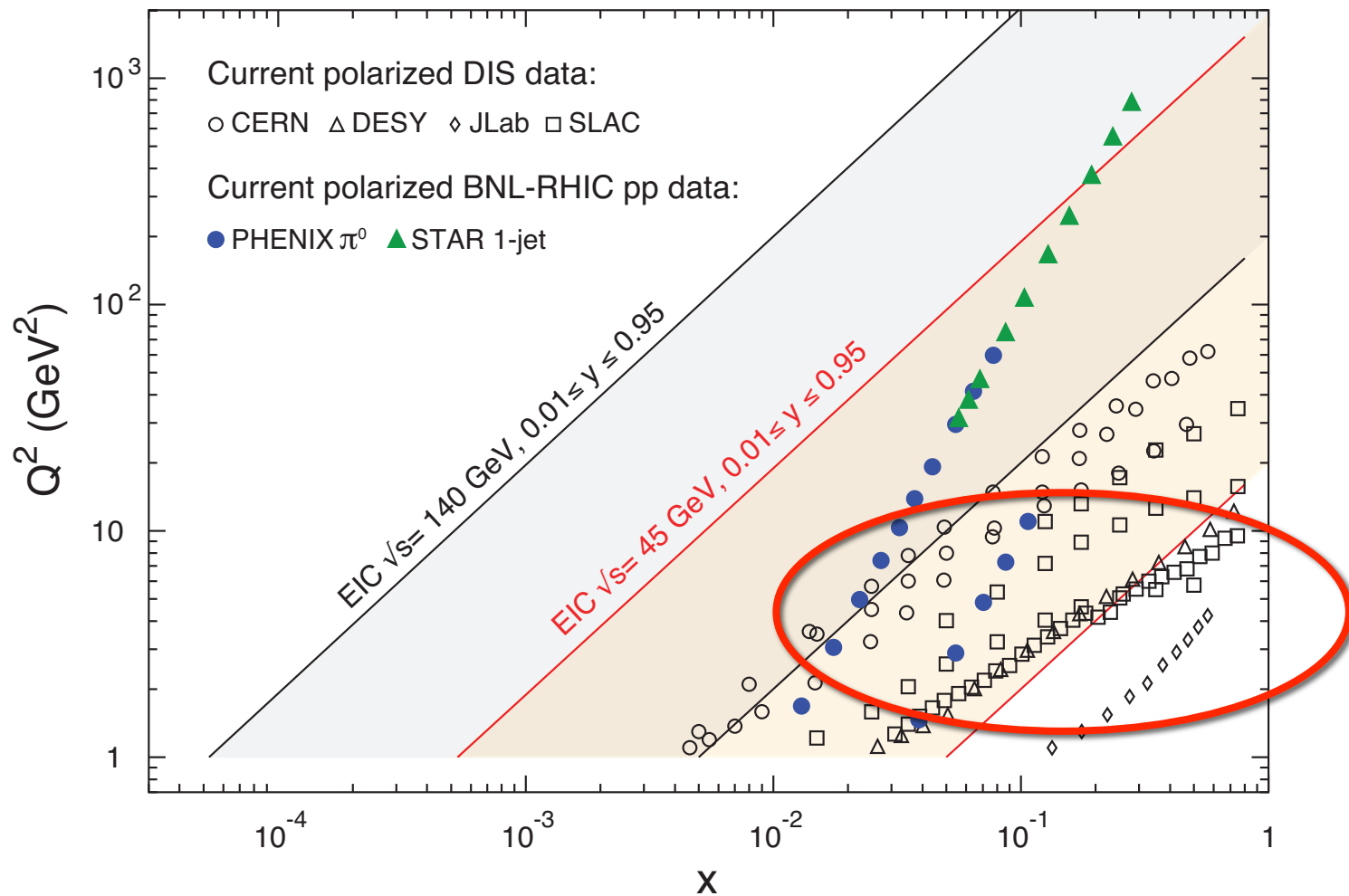
Aidala et al.1209.2803v2



Large amount of polarized data since 1998... but not in NEW kinematic region!

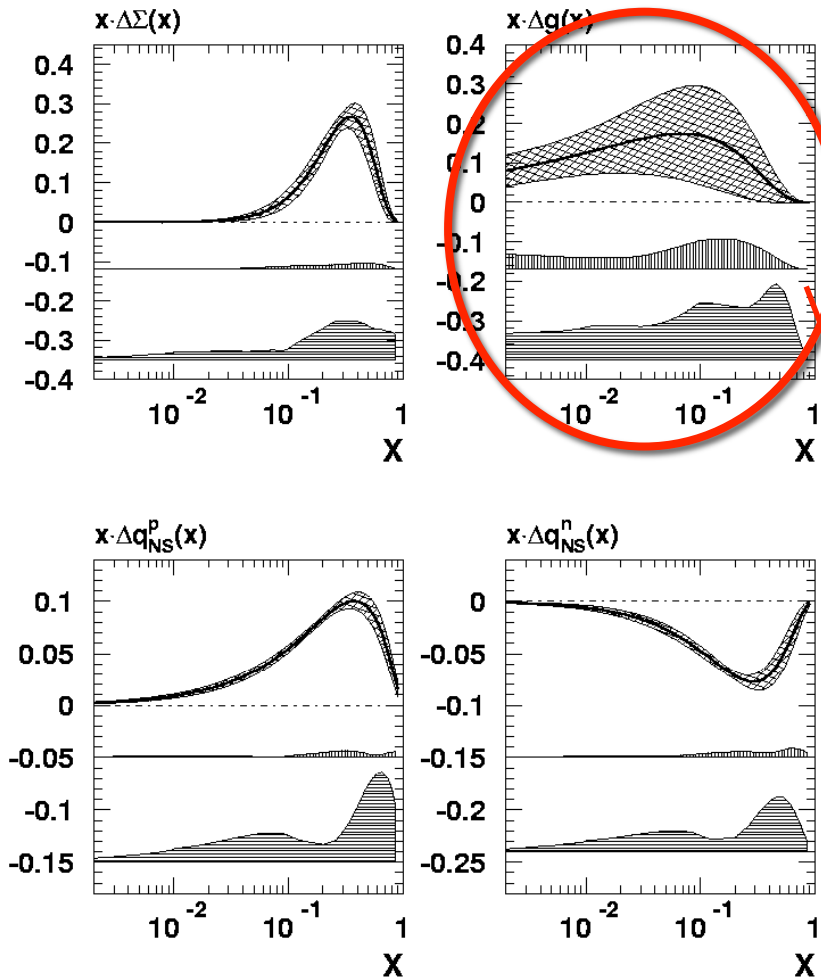
Large uncertainty in gluon polarization (+/-1.5) results from lack of wide Q^2 arm

Fixed target experiments:



Global analysis of Spin SF

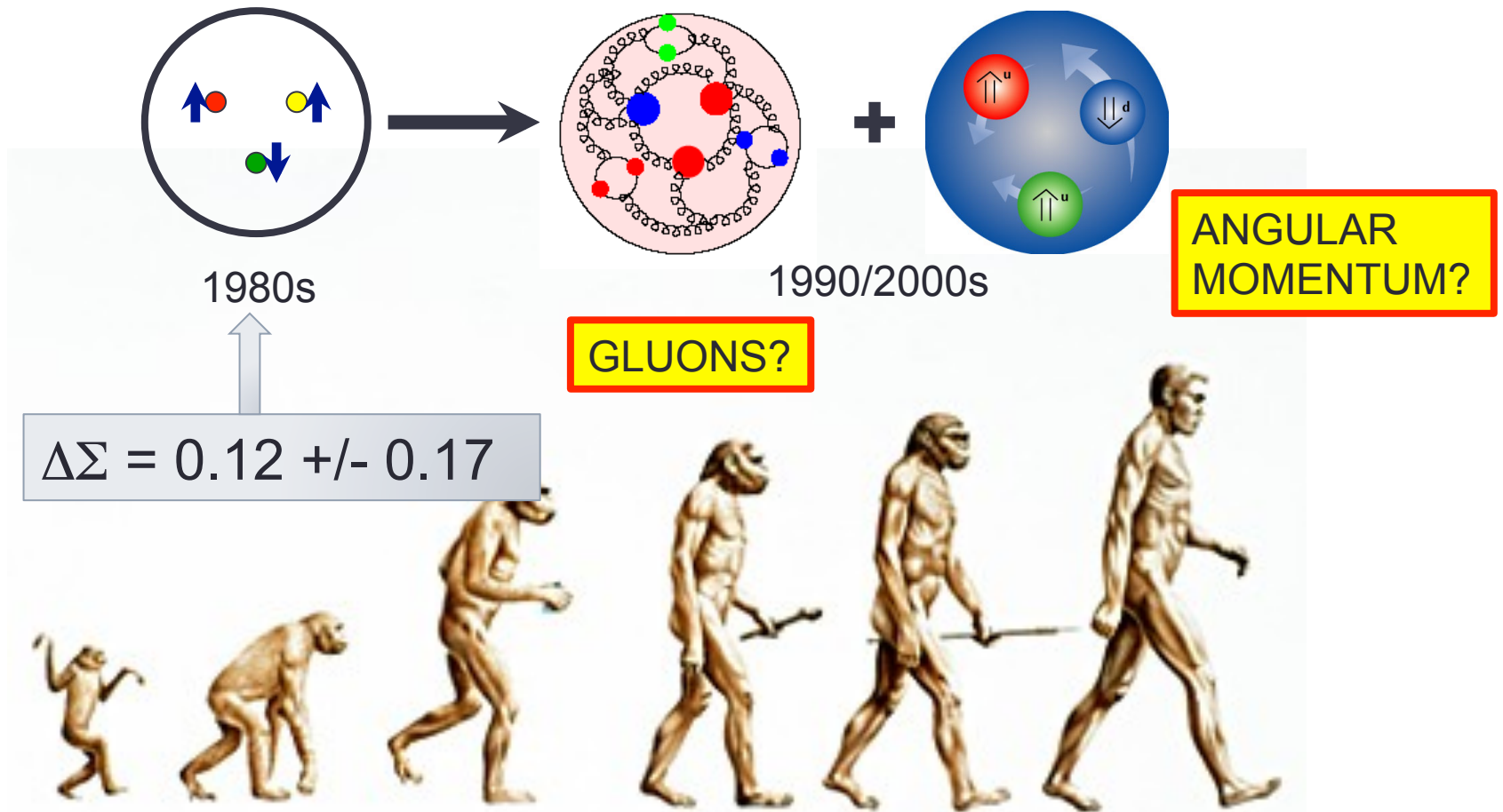
SMC PRD 58 112002 (1998)



- World's all available g_1 data
- Coefficient and splitting functions in QCD at NLO
- Evolution equations: DGLAP

$$f(x) = x^\alpha (1-x)^\beta (1 + ax + bx^2)$$
- Quark distributions fairly well determined, with small uncertainty
 - $\Delta\Sigma = 0.23 \pm 0.04$
- **Polarized Gluon distribution has largest uncertainties**
 - $\Delta G = 1 \pm 1.5$

Evolution: Our Understanding of Nucleon Spin



Spin Crisis for the prevailing models But their limitations were quickly appreciated and discussed....

Limitations of fixed target experiments:

- Kinematics of fixed target experiments: Does not allow exploration of low-x region
- Extraction of gluon polarization needed large Q^2 arm, and fixed target experiments did not allow that either....
- In 1990's ideas to achieve high energy polarized proton beams evolved... Siberian Snake Magnets
- High energy polarized proton beam polarimetry was developed as a future need...
- Ideally we needed a polarized e-p collider to overcome this, but non was under consideration! (although polarized HERA was proposed but failed)

NEED FOR A POLARIZED COLLIDER WAS UNDERSTOOD,
TO MEASURE THE EXPECTED **LARGE GLUON
POLARIZATION**

HOWEVER NO PLANS EXISTED FOR SUCH A COLLIDEER.

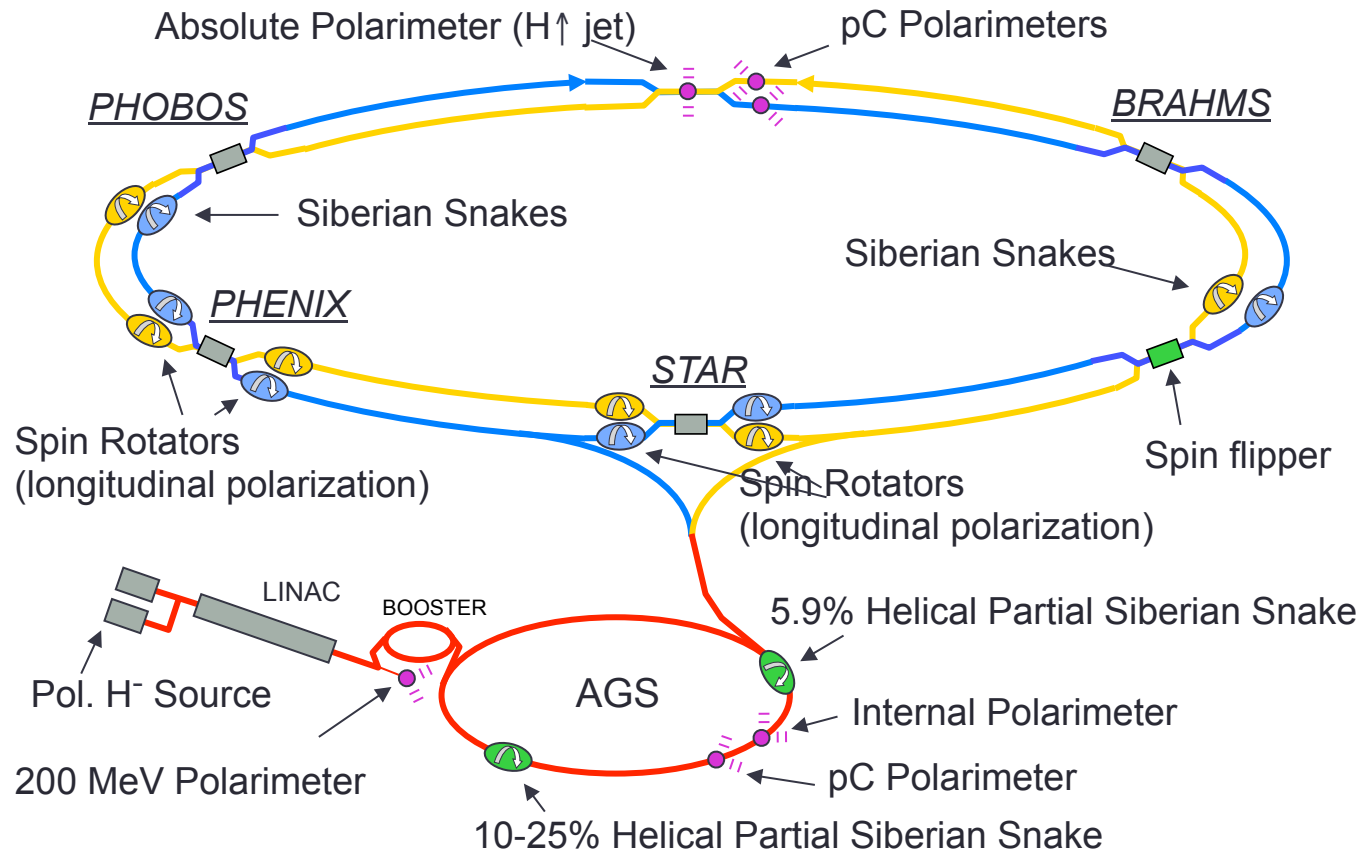
RHIC: RELATIVISTIC HEAVY ION COLLIDER WAS BEING
PLANNED FOR THE INVESTIGATIONS OF QUARK GLUON
PLASMA (AT BNL)

THE JAPANESE (RIKEN INSTITUTE) JUMPED IN WITH T.D.
LEE AND DIR. SAMIOS, TO INSTALL **SIBERIAN SNAKE
MAGNETS AND SPIN ROTATOR MAGNETS** TO ENABLE
POLARIZED PROTON COLLISIONS IN RHIC.

Motivation for RHIC Spin:

- If gluons really carry the bulk of nucleon's spin, why not use polarized proton (known by then to be predominantly made of gluons!)?
 - Technical know-how (Siberian Snakes, Spin Rotators, polarimetry ideas) to do this at high energy evolved around the time (mid/late-1990s)
- Why $\Delta\Sigma$ (quark + anti-quark's spin) small? **Are quark and anti-quark spins anti-aligned?** Polarized $p+p$ at high energy, through $W^{+/-}$ production could address this
- A severe need for investigations of the surprising transverse spin effects was naturally possible and needed with the proposed polarized $p+p$ collider...

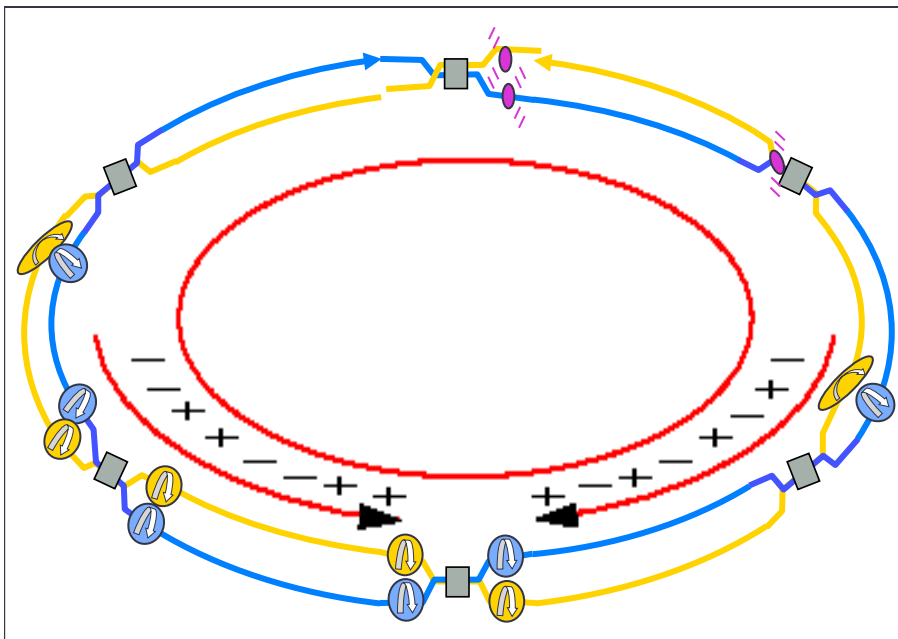
RHIC as a Polarized Proton Collider



Without Siberian snakes: $\nu_{sp} = G\gamma = 1.79 E/m \rightarrow \sim 1000$ depolarizing resonances
 With Siberian snakes (local 180° spin rotators): $\nu_{sp} = \frac{1}{2} \rightarrow$ no first order resonances
 Two partial Siberian snakes (11° and 27° spin rotators) in AGS

Measuring A_{LL}

$$A_{LL} = \frac{d\sigma_{++} - d\sigma_{+-}}{d\sigma_{++} + d\sigma_{+-}} = \frac{1}{|P_1 P_2|} \frac{N_{++} - RN_{+-}}{N_{++} + RN_{+-}}; \quad R = \frac{L_{++}}{L_{+-}}$$

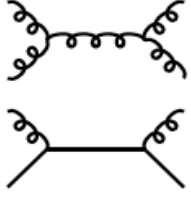
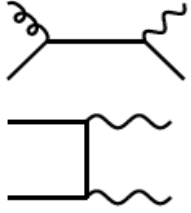
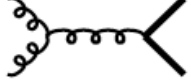
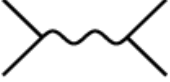
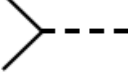


(N) Yield
(R) Relative Luminosity
(P) Polarization

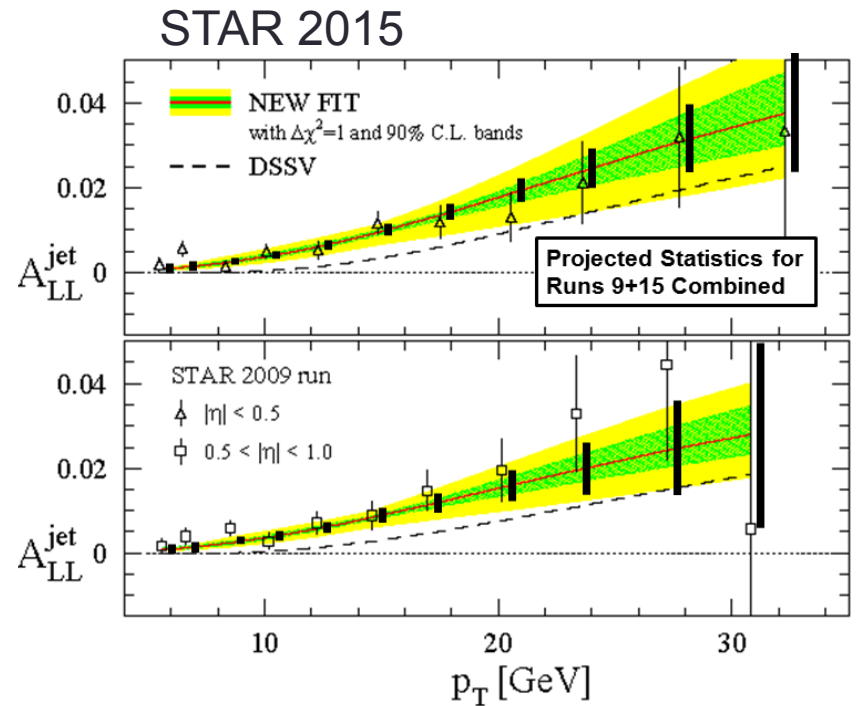
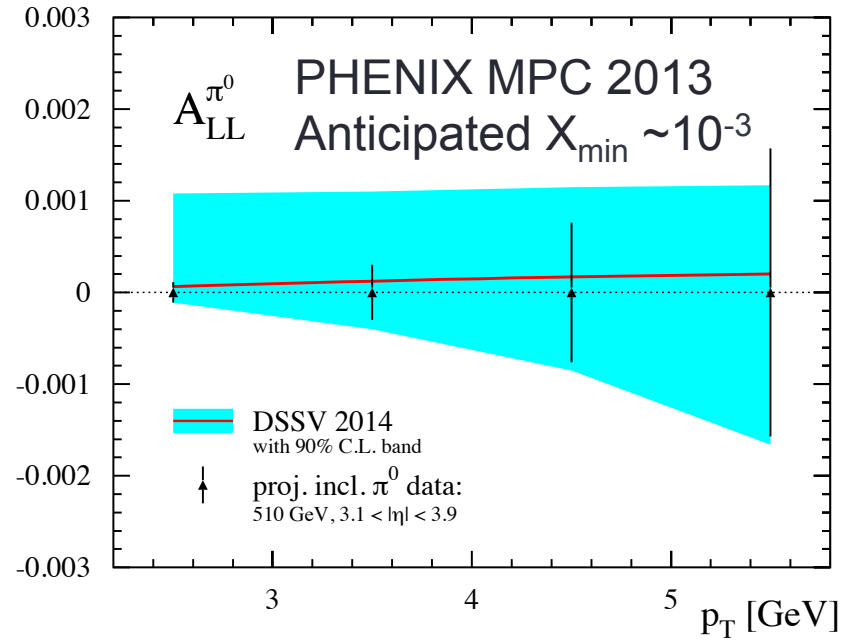
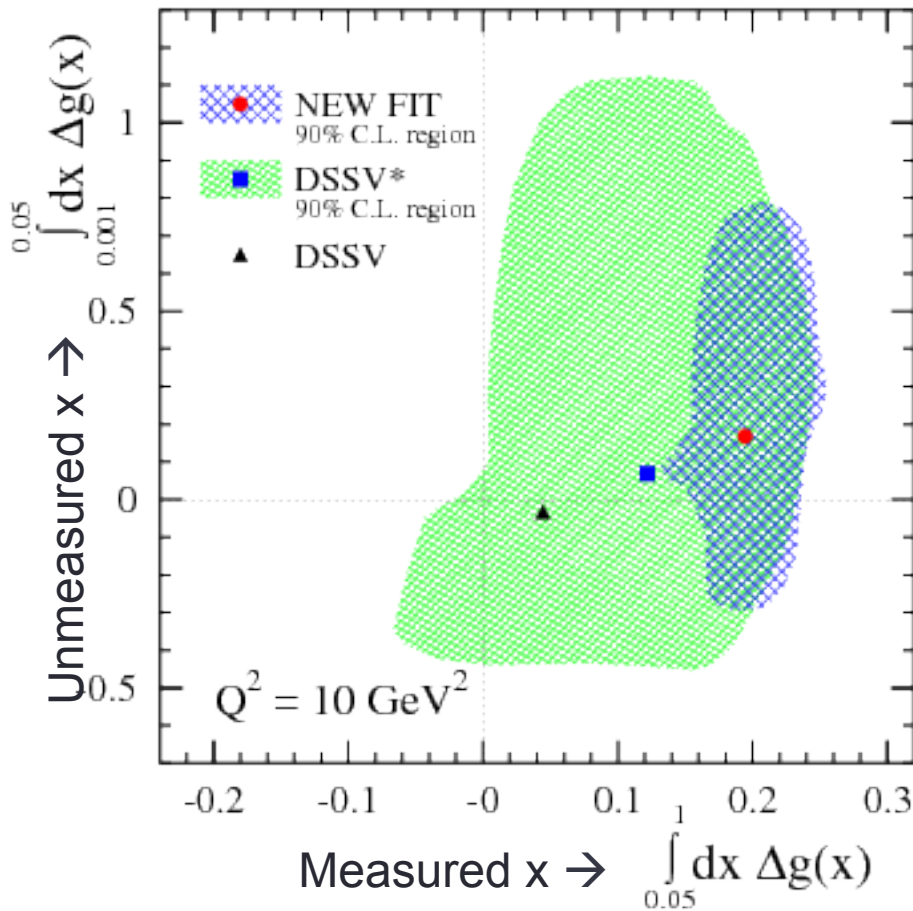
Exquisite control over false asymmetries due to ultra fast rotations of the target and probe spin.

- ✓ Bunch spin configuration alternates every 106 ns
- ✓ Data for all bunch spin configurations are collected at the same time
- ⇒ Possibility for false asymmetries are greatly reduced

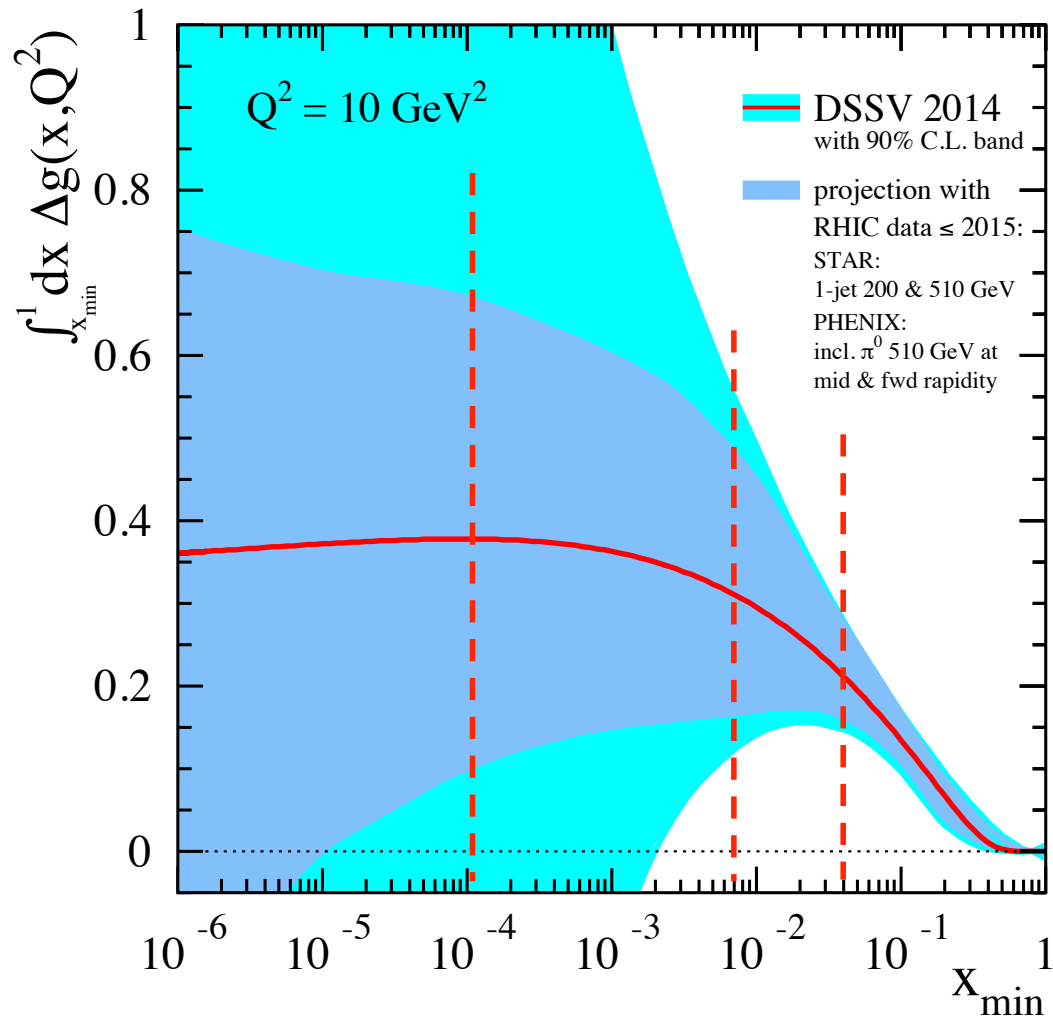
The probes and techniques at RHIC

Reaction	Dom. partonic process	probes	LO Feynman diagram
$\vec{p}\vec{p} \rightarrow \pi + X$	$\vec{g}\vec{g} \rightarrow gg$ $\vec{q}\vec{g} \rightarrow qg$	Δg	
$\vec{p}\vec{p} \rightarrow \text{jet}(s) + X$	$\vec{g}\vec{g} \rightarrow gg$ $\vec{q}\vec{g} \rightarrow qg$	Δg	(as above)
$\vec{p}\vec{p} \rightarrow \gamma + X$ $\vec{p}\vec{p} \rightarrow \gamma + \text{jet} + X$ $\vec{p}\vec{p} \rightarrow \gamma\gamma + X$	$\vec{q}\vec{g} \rightarrow \gamma q$ $\vec{q}\vec{g} \rightarrow \gamma q$ $\vec{q}\vec{q} \rightarrow \gamma\gamma$	Δg Δg $\Delta q, \Delta\bar{q}$	
$\vec{p}\vec{p} \rightarrow DX, BX$	$\vec{g}\vec{g} \rightarrow c\bar{c}, b\bar{b}$	Δg	
$\vec{p}\vec{p} \rightarrow \mu^+\mu^- X$ (Drell-Yan)	$\vec{q}\vec{q} \rightarrow \gamma^* \rightarrow \mu^+\mu^-$	$\Delta q, \Delta\bar{q}$	
$\vec{p}\vec{p} \rightarrow (Z^0, W^\pm)X$ $p\vec{p} \rightarrow (Z^0, W^\pm)X$	$\vec{q}\vec{q} \rightarrow Z^0, \vec{q}'\vec{q} \rightarrow W^\pm$ $\vec{q}'\vec{q} \rightarrow W^\pm, q'\vec{q} \rightarrow W^\pm$	$\Delta q, \Delta\bar{q}$	

Δg @ RHIC Current status: $0.05 < x < 0.2$

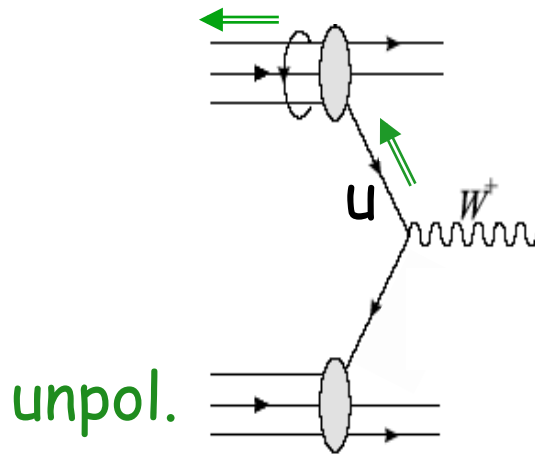


RHIC's limit... no/limited handle on low-x



RHIC Spin
White Paper (2015)
For NSAC LRP

Anti-Quark Polarization measurement via W production and decay



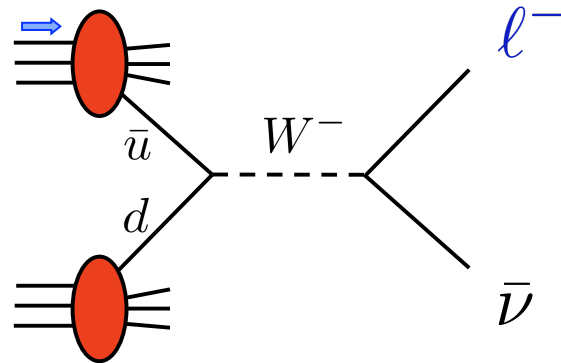
$$\sqrt{s} = 500 \text{ GeV}$$

- Large parity violating effect anticipated

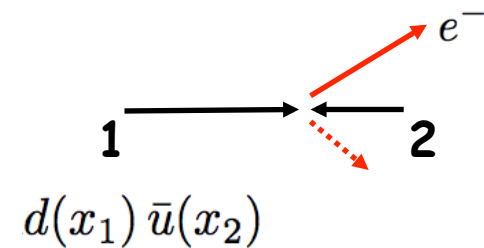
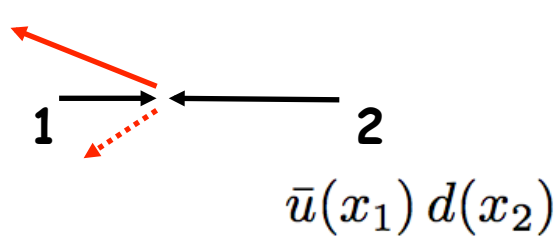
$$A_L = \frac{\sigma^+ - \sigma^-}{\sigma^+ + \sigma^-} \neq 0$$

- Measurement complimentary to SIDIS, but devoid of fragmentation function makes it cleaner!
- NLO analyses about now available

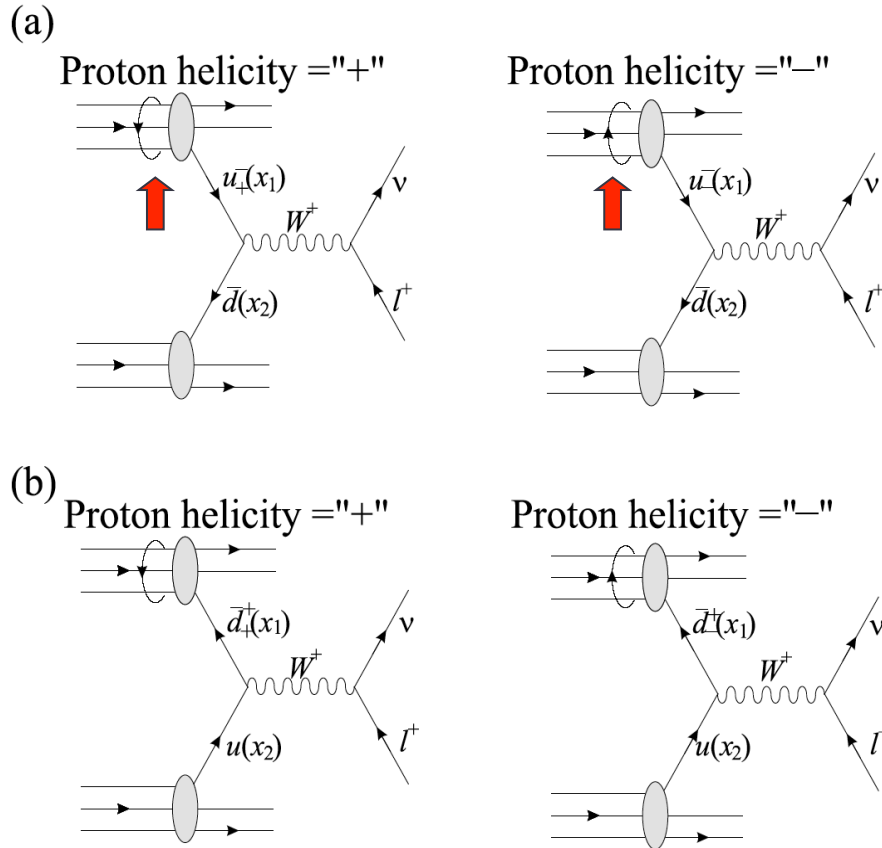
Some insight in to what goes on....



$$\sigma^{W^-} \propto \bar{u}(x_1) d(x_2) + d(x_1) \bar{u}(x_2)$$



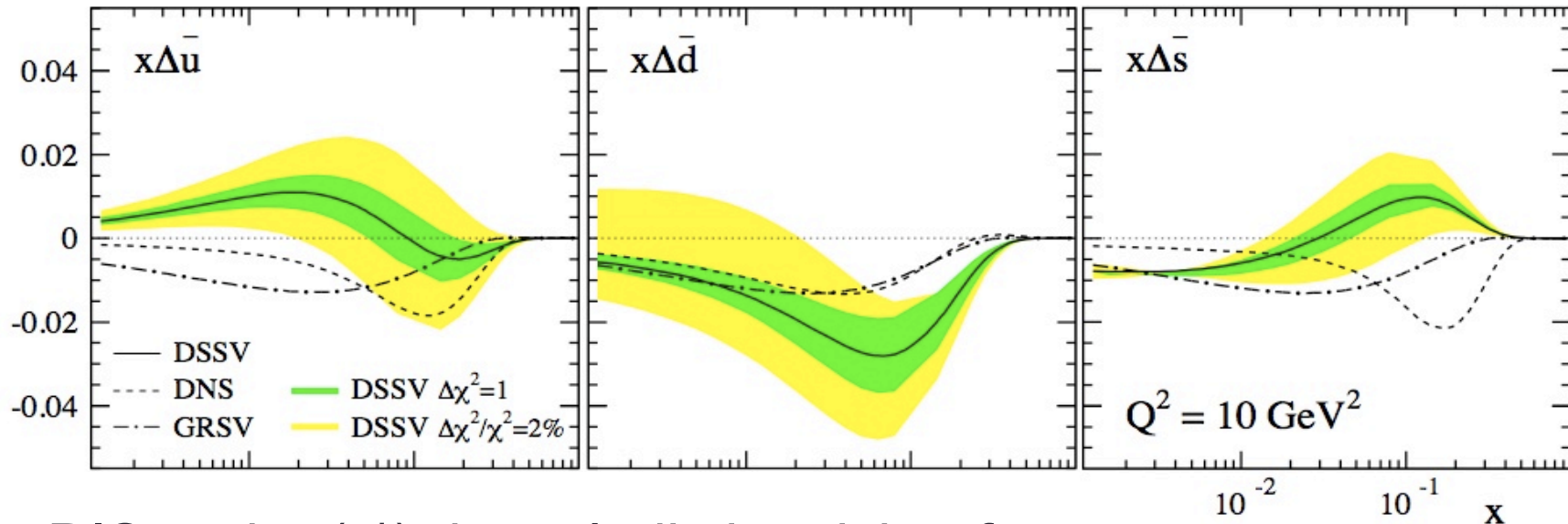
W production @ RHIC



$$A_L^{W^+} = \frac{u_-(x_1)d(x_2) - u_+(x_1)d(x_2)}{u_-(x_1)\bar{d}(x_2) + u_+(x_1)\bar{d}(x_2)} = \frac{\Delta u(x_1)}{u(x_1)}$$

$$A_L^{W^-} = \frac{\bar{d}_-(x_1)u(x_2) - \bar{d}_+(x_1)u(x_2)}{\bar{d}_-(x_1)u(x_2) + \bar{d}_+(x_1)u(x_2)} = -\frac{\Delta \bar{d}(x_1)}{\bar{d}(x_1)}$$

What about the anti-quark polarization?



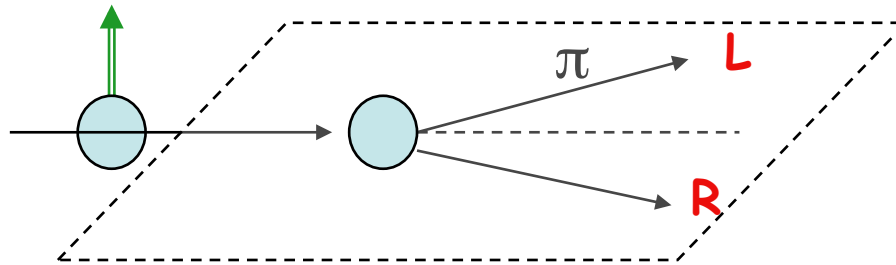
- DIS probe (γ^*) doesn't distinguish q from $q\bar{q}$
 - Has to take measure semi-inclusive (π , K production)
 - Uncertainties in fragmentation functions
- High energy p - p collisions enable probing $q, q\bar{q}$ through $W^{+/-}$ production \rightarrow Plan at RHIC

In parallel:

The Transverse Spin Puzzle

Had been observed but *ignored* for almost 3 decades...

Transverse spin introduction



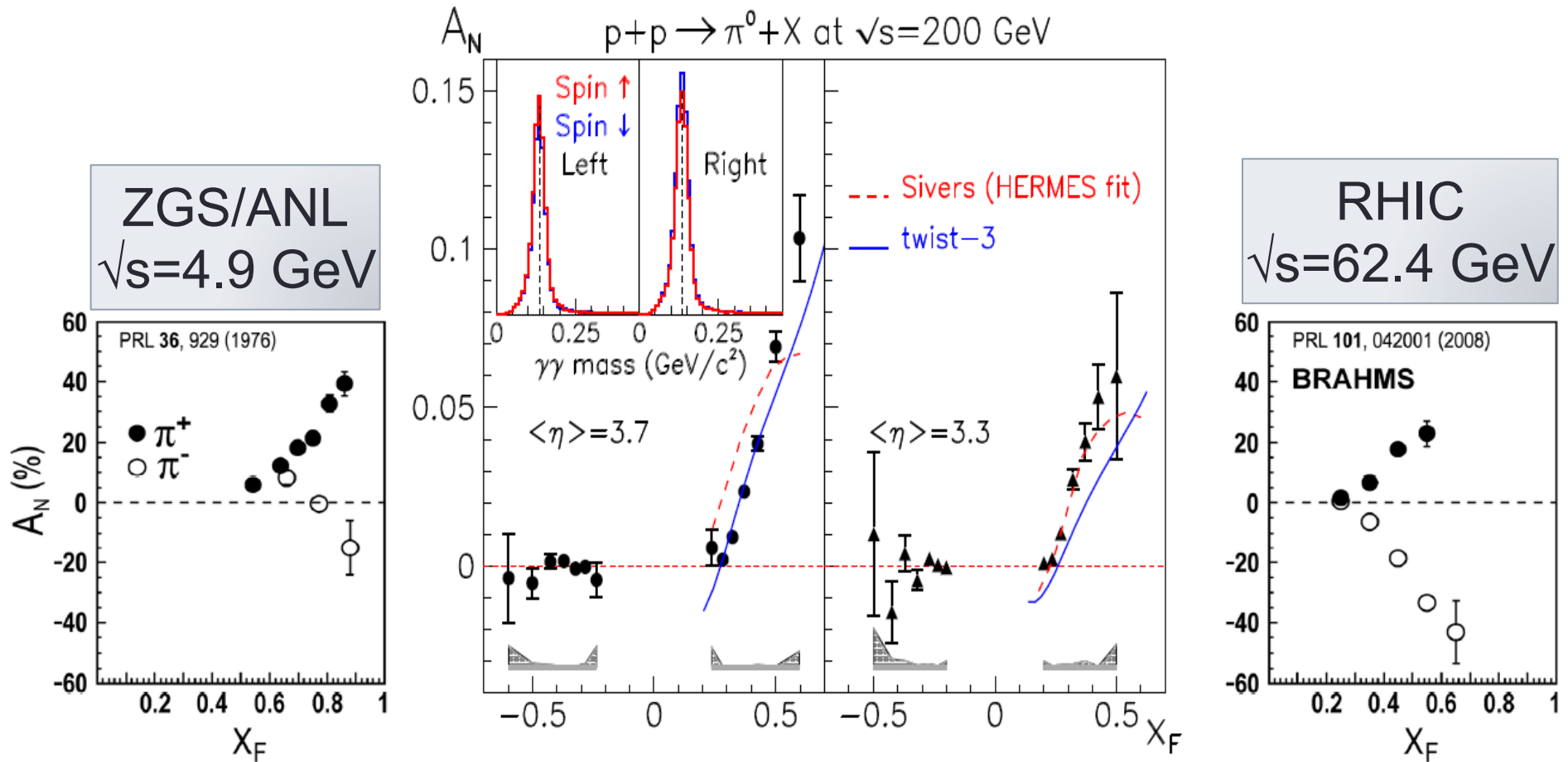
$$A_N = \frac{N_L - N_R}{N_L + N_R}$$

$$A_N \sim \frac{m_q}{p_T} \alpha_S$$

Kane, Pumplin, Repko 1978
PRL 41 1689 (1978)

- Since people started to measure effects at high p_T to interpret them in pQCD frameworks, this was “neglected” as it was expected to be small..... However....
- Pion production in single transverse spin collisions showed us something different....

Pion asymmetries: at most CM energies!

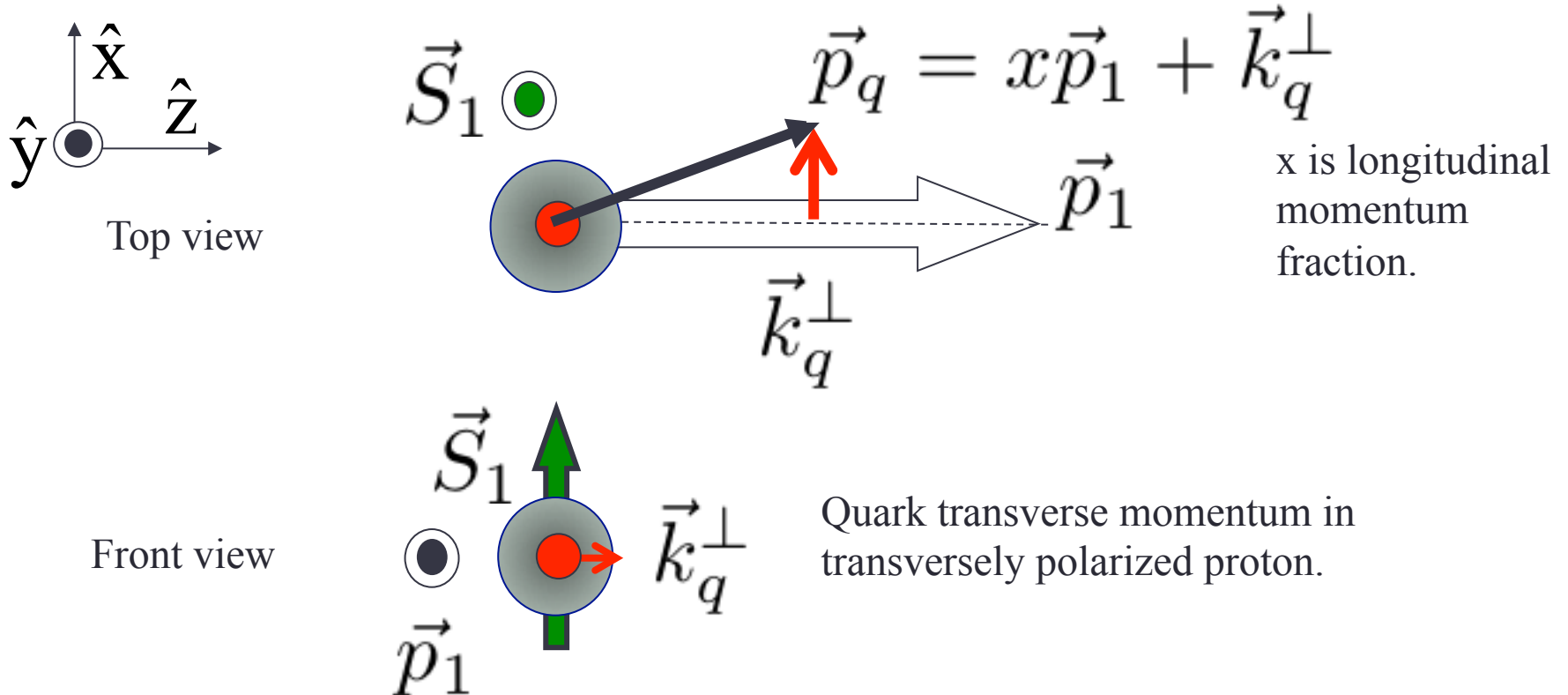


Suspect soft QCD effects at low scales, but they seem to remain relevant to perturbative regimes as well

Sivers effect: due to transverse motion of quarks in the nucleon: initial state effect

Phys Rev D41 (1990) 83; Phys Rev D43 (1991) 261

$$SSA_{Sivers} \propto \vec{S}_1 \cdot (\vec{p}_1 \times \vec{k}_q)$$



INITIAL STATE EFFECT: Orbital angular momentum?

What does “Sivers effect” probe?

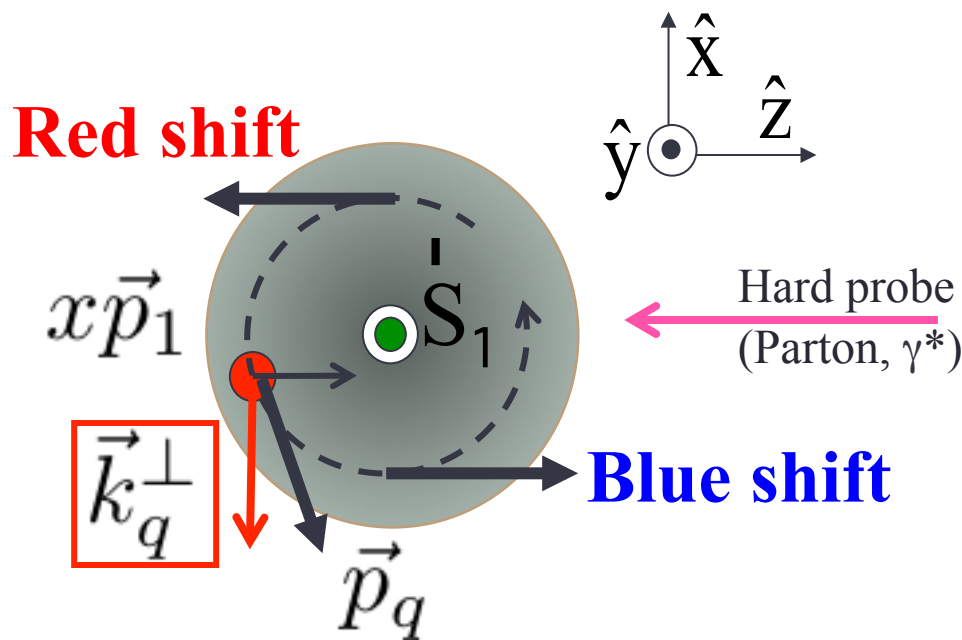
Top view, Breit frame

Quarks orbital motion adds/ subtracts longitudinal momentum for negative/positive \hat{x} .

PRD66 (2002) 114005

Parton Distribution Functions rapidly fall in longitudinal momentum fraction x .

Final State Interaction between outgoing quark and target spectator.



Sivers function
 $f_{1T}^\perp(x, \vec{k}_q^\perp)$

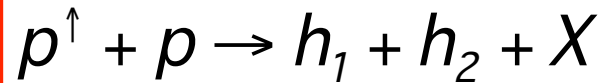
hep-ph/
0703176

Quark Orbital angular momentum

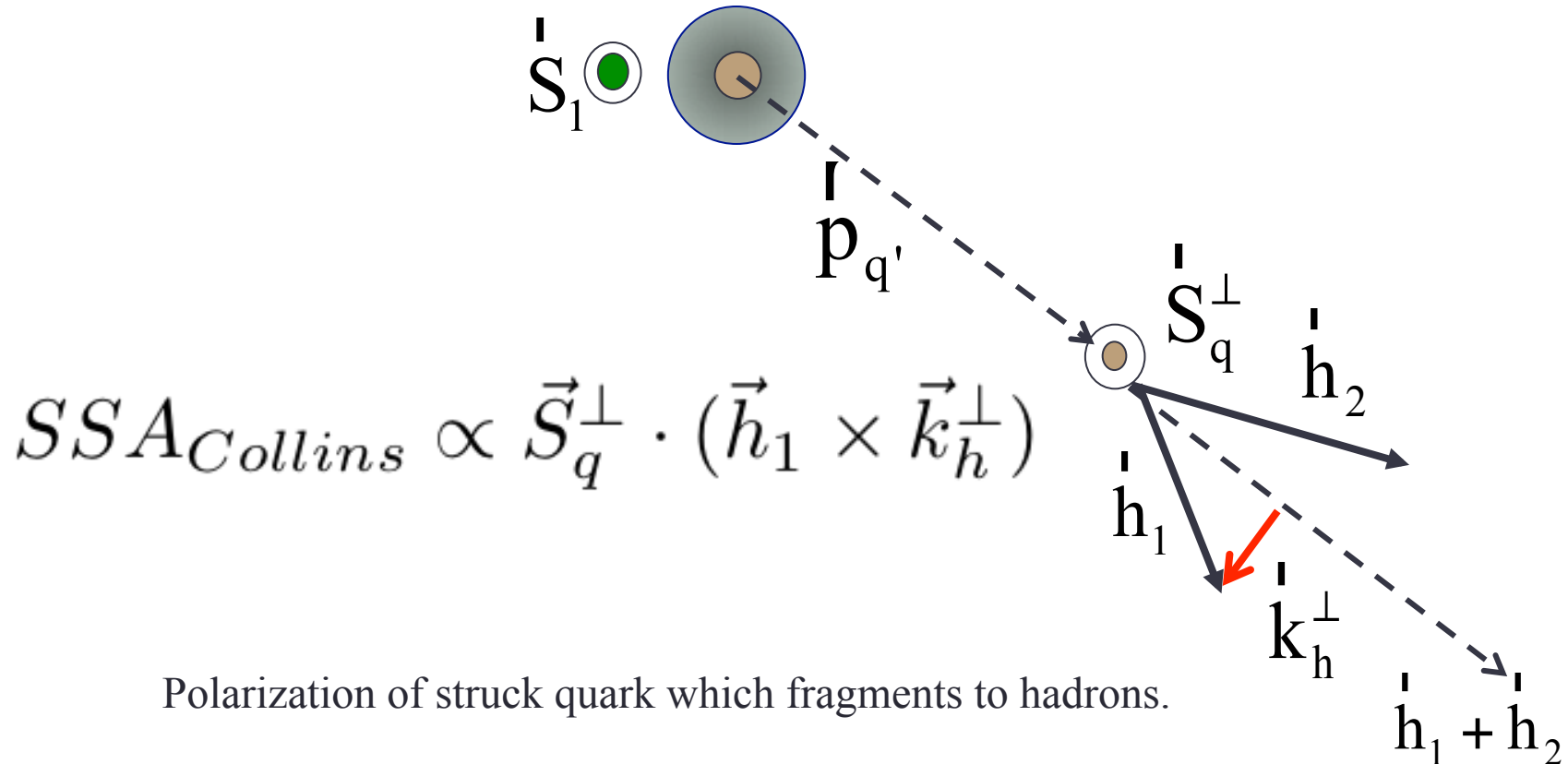
Generalized Parton Distribution Functions
 PRD59 (1999) 014013

Collins (Heppelmann) effect: Asymmetry in the fragmentation hadrons

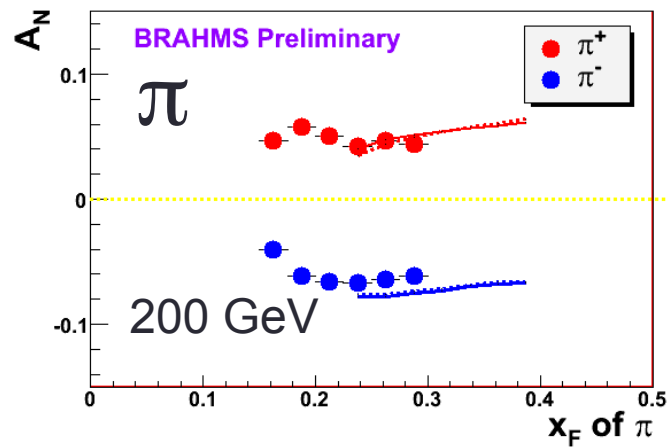
Example:



Nucl Phys B396 (1993) 161,
Nucl Phys B420 (1994) 565

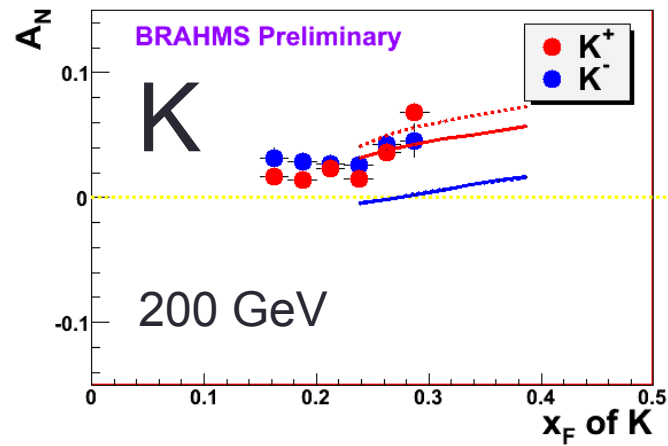


Polarization of struck quark which fragments to hadrons.

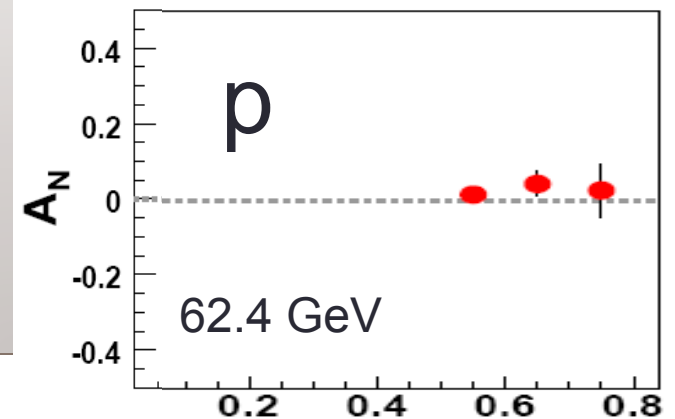
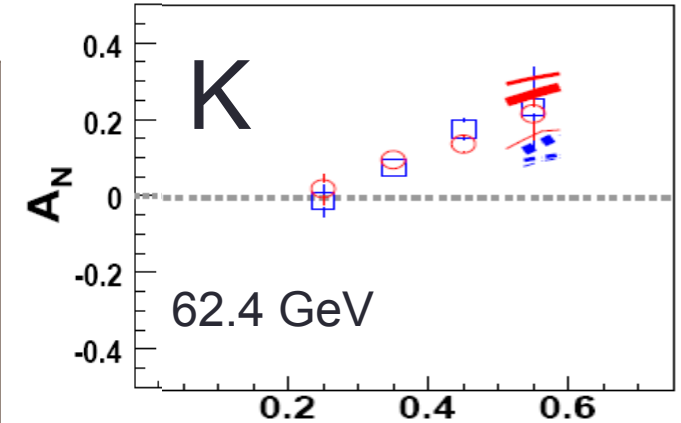
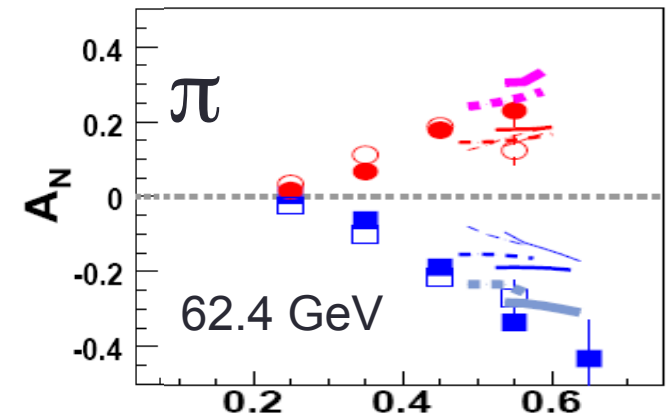
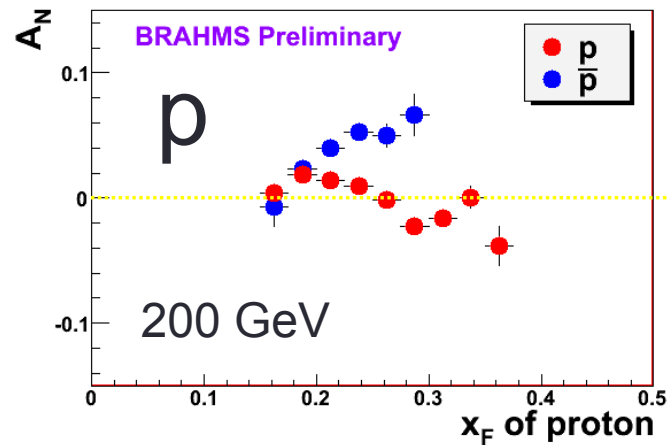


π, K, p
200 & 62.4
GeV

BRAHMS



- Scales on plots different
- Kaon asymmetries not predicted
- Unfortunately no anti-proton measurement



Although not expected, at any observable level, 400+ times the expected values of asymmetries have been routinely seen experimentally: both in ep and pp systems.

- Transverse motion/momentum of partons (indirect evidence for orbital angular momentum of the quarks & gluons?)
- Asymmetry in fragmentation process (final state) or
- Both May be responsible.

Summary

- RHIC Spin program made great strides in furthering our knowledge of nucleon, but we need something more to address some of the still open issues:

- **What remains?**

- ΔG at low x ? Spin structure functions and its behavior at low x ?
- If orbital angular motion plays a role, what is the orbital contribution from Gluons?

Precision measurements at a future facility are essential...

THE ELECTRON ION COLLIDER.... Can we do better?

Study of internal structure of a watermelon:



A-A (RHIC)
1) Violent collision of melons



2) Cutting the watermelon with a knife
Violent DIS e-A (EIC)



3) MRI of a watermelon
Non-Violent e-A (EIC)

