

Nuclear structure III (theory) Witek Nazarewicz (UTK/ORNL) National Nuclear Physics Summer School 2014 William & Mary, VA

- Nuclear Force
- General principles
- Examples: quantitative nuclear theory; predictive capability
- Realities of highperformance





The Force





- Nucleon r.m.s. radius ~0.86 fm
- Comparable with interaction range
- Half-density overlap at max. attarction

• $V_{\rm NN}$ not fundamental (more like inter-molecular van der Waals interaction)

are expected.

Since nucleons are composite objects, three-and higher-body forces

Nuclear force

A realistic nuclear force force: schematic view



Nucleon-Nucleon interaction (qualitative analysis)



There are infinitely many equivalent nuclear potentials! $\hat{H}\Psi = E\Psi$ $(\hat{U}\hat{H}\hat{U}^{-1})\hat{U}\Psi = E\hat{U}\Psi$

Reid93 is from V.G.J.Stoks et al., PRC**49**, 2950 (1994).

AV16 is from R.B.Wiringa et al., PRC**51**, 38 (1995).

nucleon-nucleon interactions



N3LO: Entem et al., PRC68, 041001 (2003)

Epelbaum, Meissner, et al.

Renormalization group (RG) evolved nuclear potentials



Bogner, Kuo, Schwenk, Phys. Rep. 386, 1 (2003)

three-nucleon interactions

ridal bulges from moon and Sun Three-body forces between protons and neutrons are Earth analogous to tidal forces: the gravitational force on the Moo Earth is not just the sum of Earth-Moon and Earth-Sun Water forces (if one employs point masses for Earth, Moon, Orbital Paths of Earth and Sun) Moon The computational cost of nuclear 3-body forces can © 2003 Stuart J. Robbins be greatly reduced by decoupling low-energy parts from high-energy parts, which can then be discarded. k'^{2} (fm⁻²) k'^{2} (fm⁻²) k'² (fm⁻²) k'² (fm⁻²) k'^{2} (fm⁻²) 4 8 12 0 4 8 12 0 4 8 12 0 4 8 12 0 4 8 12 0 4 8 12 0.5 $k^{2} \; (fm^{-2})$ 0 (fm) 12 $\lambda = 2.0 \text{ fm}^{-1}$ $\lambda = 1.5 \text{ fm}^{-1}$ $\lambda = 3.0 \text{ fm}^{-1}$ λ =4.0 fm⁻' -0.5

Recently the first consistent Similarity Renormalization Group softening of three-body forces was achieved, with rapid convergence in helium. With this faster convergence, calculations of larger nuclei are possible!

The challenge and the prospect: NN force





Optimizing the nuclear force

input matters: garbage in, garbage out

- The derivative-free minimizer POUNDERS was used to systematically optimize NNLO chiral potentials
- The optimization of the new interaction NNLO_{opt yields a χ} 2/datum ≈ 1 for laboratory NN scattering energies below 125 MeV. The new interaction yields very good agreement with binding energies and radii for A=3,4 nuclei and oxygen isotopes
- Ongoing: Optimization of NN + 3NF



http://science.energy.gov/np/highlights/2014/np-2014-05-e/

- Used a coarse-grained representation of the short-distance interactions with 30 parameters
- The optimization of a chiral interaction in NNLO 2/datum ≈ 1 for a mutually consistent set of 6713 NN scattering data
- Covariance matrix yields correlation between LECCs and predictions with error bars.



Deuteron, Light Nuclei

Deuteron

Binding energy	2.225 MeV
Spin, parity	1+
Isospin	0
Magnetic moment	μ=0.857 μ _N
Electric quadrupole moment	Q=0.282 e fm ²

 $\mu_p + \mu_n = 2.792\mu_N - 1.913\mu_N = 0.879\mu_N$

$$|\psi_d\rangle = 0.98 |{}^3S_1\rangle + 0.20 |{}^3D_1\rangle$$

produced by tensor force!



Nucleon-Nucleon Interaction NN, NNN, NNNN,..., forces

GFMC calculations tell us that:

 $\langle V_{\pi} \rangle / \langle V \rangle \sim 70 - 80\%$ $\langle V_{\pi} \rangle \sim -15 \text{ MeV/pair}$ $\langle V^{R} \rangle \sim -5 \text{ MeV/pair}$ $\langle V^{3} \rangle \sim -1 \text{ MeV/three}$ $\langle T \rangle \sim 15 \text{ MeV/nucleon}$ $\langle V_{C} \rangle \sim 0.66 \text{ MeV/pair of protons}$

Few-nucleon systems (theoretical struggle)



A=2: many years ago...

3H: 1984 (1% accuracy)

- Faddeev
- Schroedinger



3He: 1987

4He: 1987

5He: 1994 (n-α resonance)



A=6,7,..12: 1995-2014

Happy the man who has been able to discern the cause of things

Virgil, Georgica

Theories Models

- A first rate theory predicts
- A second rate theory forbids
- A third rate theory explains after the facts

Alexander I. Kitaigorodskii

Weinberg's Laws of Progress in Theoretical Physics From: "Asymptotic Realms of Physics" (ed. by Guth, Huang, Jaffe, MIT Press, 1983)

First Law: "The conservation of Information" (*You will get nowhere* by churning equations)

Second Law: "Do not trust arguments based on the lowest order of perturbation theory"

Third Law: "You may use any degrees of freedom you like to describe a physical system, but if you use the wrong ones, you'll be sorry!"



Patient: Doctor, doctor, it hurts when I do this! Doctor: Then don't do that.





Modeling the Atomic Nucleus



The Nuclear Many-Body Problem

 $2^A \times \frac{A!}{N!Z!} \begin{array}{c} \text{coupled integro-differential} \\ \text{equations in 3A dimensions} \end{array}$

How to explain the nuclear landscape from the bottom up? Theory roadmap



Theory of nuclei is demanding

- rooted in QCD
- insights from EFT
- many-body interactions
- in-medium renormalization
- microscopic functionals
- low-energy coupling constants optimized to data
- crucial insights from exotic nuclei

- many-body techniques
 - direct ab initio schemes
 - symmetry breaking and restoration
- high-performance computing
- interdisciplinary connections



- nuclear structure impacted by couplings to reaction and decay channels
- clustering, alpha decay, and fission still remain major challenges for theory
- unified picture of structure and reactions

Illustrative physics examples

Ab initio theory for light nuclei and nuclear matter



(nuclei, neutron droplets, nuclear matter)



Green's Function Monte Carlo (imaginary-time method)

$$\begin{split} \left|\psi_{0}\right\rangle &= \lim_{\tau \to \infty} e^{-(\hat{H} - E_{0})\tau} \left|\psi_{V}\right\rangle \\ \hline \mathbf{Trial \ wave \ function} \\ \left|\psi(\tau)\right\rangle &= e^{-(\hat{H} - E_{0})\tau} \left|\psi_{V}\right\rangle \\ \left|\psi(0)\right\rangle &= \left|\psi_{V}\right\rangle, \quad \left|\psi(\infty)\right\rangle &= \left|\psi_{0}\right\rangle \\ \tau &= n\Delta\tau \quad \Rightarrow \quad \left|\psi(\tau)\right\rangle &= \left[e^{-(\hat{H} - E_{0})\Delta\tau}\right]^{n} \left|\psi_{V}\right\rangle \\ \end{split}$$

Other methods:

- Faddeev-Yakubovsky method
- Hyperspherical harmonics method
- Coupled-cluster expansion method, exp(S)
- Cluster approaches (resonating group method)
- No-core shell model
- Lattice methods



GFMC: S. Pieper, ANL

1-2% calculations of A = 6-12 nuclear energies are possible excited states with the same quantum numbers computed

12C: ground state and Hoyle state

state-of-the-art computing

 10^{-1}

10-2

 10^{-1}

 2^{+}

-82.6(1)

0⁺ -84.51

2⁺ -87.72

0+

-92.16

Exp

 $f_{pt}^{}(k)$

Wiringa et al. Phys. Rev. C 89,

024305 (2014); A. Lovato et al.,

10

10

10

10

e x

• ρ_{1b+2b}

[[[]]

Phys. Rev. Lett. 112, 182502 (2014)

0.08

ŕ (fm

-82

-84

-86

-88

-90

-92

E [MeV]

q (fm ^{.1})



Epelbaum et al., Phys. Rev. Lett. 109, 252501 (2012). Lattice EFT Lahde et al., Phys. Lett. B 732, 110 (2014).



0⁺ -92(3)

The frontier: neutron-rich calcium isotopes probing nuclear forces and shell structure in a neutron-rich medium



Ab Initio Path to Heavy Nuclei Binder et al., arXiv:1312.5685



- The first accurate ab initio cupled cluster calculations for heavy nuclei using *SRG-evolved chiral interactions*. A number of technical hurdles eliminated
- Many-body calculations up to ^{132Sn are now possible with controlled uncertainties on the order of 2%}
- A first direct validation of chiral Hamiltonians in the regime of heavy nuclei using ab initio techniques.
- Future studies will have to involve consistent chiral Hamiltonians at N3LO considering initial and SRG-induced 4N interactions and provide an exploration of other observables.

Configuration interaction techniques

- light and heavy nuclei
- detailed spectroscopy
- quantum correlations (lab-system description)







One-body Hamiltonia

- Construct basis states with good $(J_{z, Tz) \text{ or } (J,T)}$
- Compute the Hamiltonian matrix
- Diagonalize Hamiltonian matrix for lowest eigenstates
- Number of states increases dramatically with particle number

Full *fp* shell for ${}^{60}Zn : \approx 2 \times 10^9 J_z$ states 5,053,594 J = 0, T = 0 states 81,804,784 J = 6, T = 1 states



Residual

interactioni

= X + X + X →



• Can we get around this problem? Effective interactions in truncated spaces (*P*-included, finite; *Q*-excluded, infinite)

- Residual interaction (G-matrix) depends on the configuration space. Effe 🗔 = 🔀 + 💢 🕻
- Breaks down around particle drip lines

G-matrix, obtained from the Bethe-Goldstone equation (scattering within a nuclear medium)

Microscopic valence-space Shell Model Hamiltonian

Coupled Cluster Effective Interaction (valence cluster expansion)



In-medium SRG Effective Interaction

G.R. Jansen et al., arXiv:1402.2563



Anomalous Long Lifetime of ^{14C}

Determine the microscopic origin of the suppressed β -decay rate: 3N force







Dimension of matrix solved for 8 lowest states ~ 10⁹ Solution took ~ 6 hours on 215,000 cores on Cray XT5 Jaguar at ORNL

Fusion of Light Nuclei

Computational nuclear physics enables us to reach into regimes where <u>experiments and</u> <u>analytic theory are not possible</u>, such as the cores of fission reactors or hot and dense evolving environments such as those found in inertial confinement fusion environment.



Ab initio theory reduces uncertainty due to conflicting

- The n-3H elastic cross section for 14 MeV neutrons, important for NIF, was not known precisely enough.
 - Delivered evaluated data with required 5% uncertainty and successfully compared to measurements using an Inertial Confinement Facility
- "First measurements of the differential cross sections for the elastic n-2H and n-3H scattering at 14.1 MeV using an Inertial Confinement Facility", by J.A. Frenje *et al.*, Phys. Rev. Lett. **107**, 122502 (2011)

http://physics.aps.org/synopsis-for/10.1103/PhysRevLett.107.122502

Nuclear Density Functional Theory and Extensions

Mean-Field Theory ⇒ Density Functional Theory Degrees of freedom: nucleonic densities

Nuclear DFT

- two fermi liquids
- self-bound
- superfluid
- mean-field ⇒ one-body densities
- zero-range \Rightarrow local densities
- finite-range \Rightarrow gradient terms
- particle-hole and pairing channels
- Has been extremely successful. A broken-symmetry generalized product state does surprisingly good job for nuclei.

Examples: Nuclear Density Functional Theory

Traditional (limited) functionals provide quantitative description

BE differences

Cwiok et al., Nature, 433, 705 (2005)

Example: Large Scale Mass Table Calculations HFB+LN mass table, HFBTHO

- 5,000 even-even nuclei, 250,000 HFB runs, 9,060 processors about 2 CPU hours
- Full mass table: 20,000 nuclei, 12M configurations full JAGUAR

Description of observables and model-based extrapolation

- Systematic errors (due to incorrect assumptions/poor modeling)
- Statistical errors (optimization and numerical errors)

Erler et al., Nature (2012)

Quantified Nuclear Landscape

How many protons and neutrons can be bound in a nucleus? Erler et al. Nature 486, 509 (2012) Skyrme-DFT: 6,900±500_{syst}

From nuclei to neutron stars (a multiscale problem)

Gandolfi et al. PRC85, 032801 (2012)

J. Erler et al., PRC 87, 044320 (2013)

The covariance ellipsoid for the neutron skin R_{skin in}^{208Pb and the radius of a}

The mean values are: R(1.4M \odot)=12 km and Rskin= 0.17 fm.

Major uncertainty: density dependence of the symmetry energy. Depends on T=3/2three-nucleon forces