

Neutrino Physics

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Part 3

Looking to the Future

The Open Questions

- What is the absolute scale of neutrino mass?
- Are neutrinos their own antiparticles?
- Are there *more* than 3 mass eigenstates?
 - Are there “sterile” neutrinos?
- What are the neutrino magnetic and electric dipole moments?

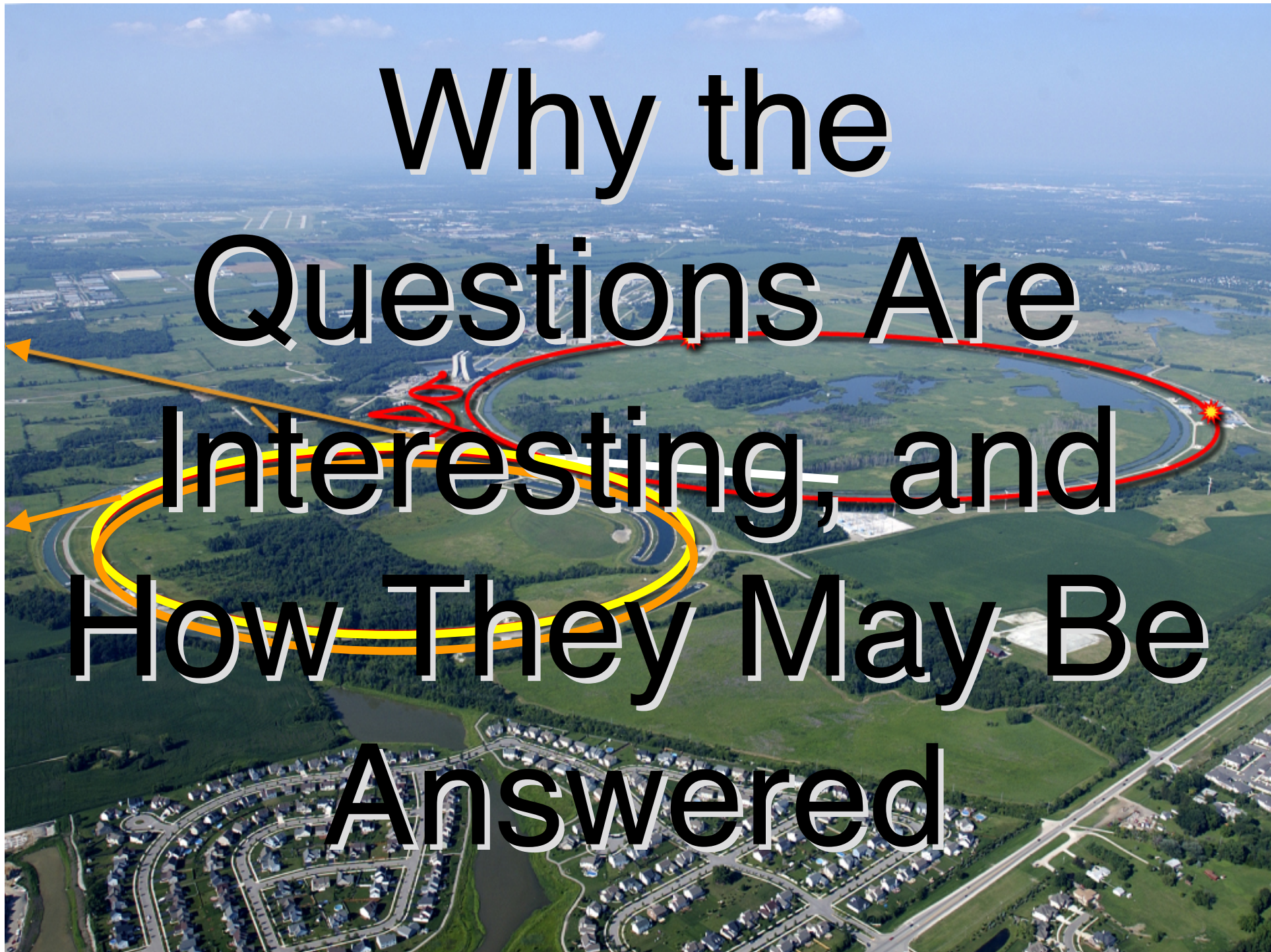
What is θ_{13} ?
How close to maximal is θ_{23} ?

• Is the spectrum like $\underline{\underline{=}}$ or $\underline{=}$?

• Do neutrino interactions
violate CP?

Is $P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta) \neq P(\nu_\alpha \rightarrow \nu_\beta)$?

- What can neutrinos and the universe tell us about one another?
- Is CP violation involving neutrinos the key to understanding the matter – antimatter asymmetry of the universe?
- What physics is behind neutrino mass?
- What *surprises* are in store?



**Why the
Questions Are
Interesting, and
How They May Be
Answered**

Does $\bar{v} = v$?

What Is the Question?

For each *mass eigenstate* ν_i , and *given helicity* h , does —

- $\bar{\nu}_i(h) = \nu_i(h)$ (Majorana neutrinos)

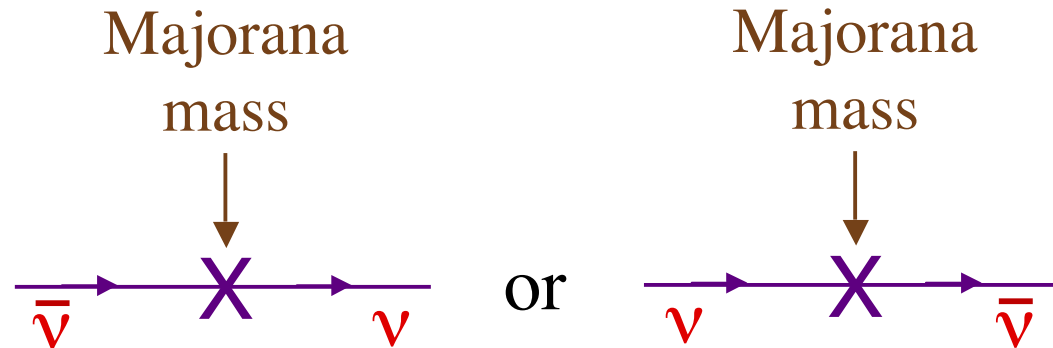
or

- $\bar{\nu}_i(h) \neq \nu_i(h)$ (Dirac neutrinos) ?

Equivalently, do neutrinos have *Majorana masses*? If they do, then the mass eigenstates are *Majorana neutrinos*.

Majorana Masses

Their effect:



Majorana masses mix ν and $\bar{\nu}$, so they do not conserve the **Lepton Number L** that distinguishes leptons from antileptons:

$$L(\nu) = L(\ell^-) = -L(\bar{\nu}) = -L(\ell^+) = 1$$

A Majorana mass for any fermion f causes $f \leftrightarrow \bar{f}$.

Quark and *charged-lepton* Majorana masses are forbidden by electric charge conservation.

Neutrino Majorana masses would make the neutrinos *very* distinctive.

Majorana neutrino masses have a different origin than the quark and charged-lepton masses.

Why Majorana Masses \longrightarrow Majorana Neutrinos

As a result of $K^0 \longleftrightarrow \bar{K}^0$ mixing, the neutral K mass eigenstates are —

$$K_{S,L} \cong (K^0 \pm \bar{K}^0)/\sqrt{2} . \quad \overline{K_{S,L}} = K_{S,L} .$$

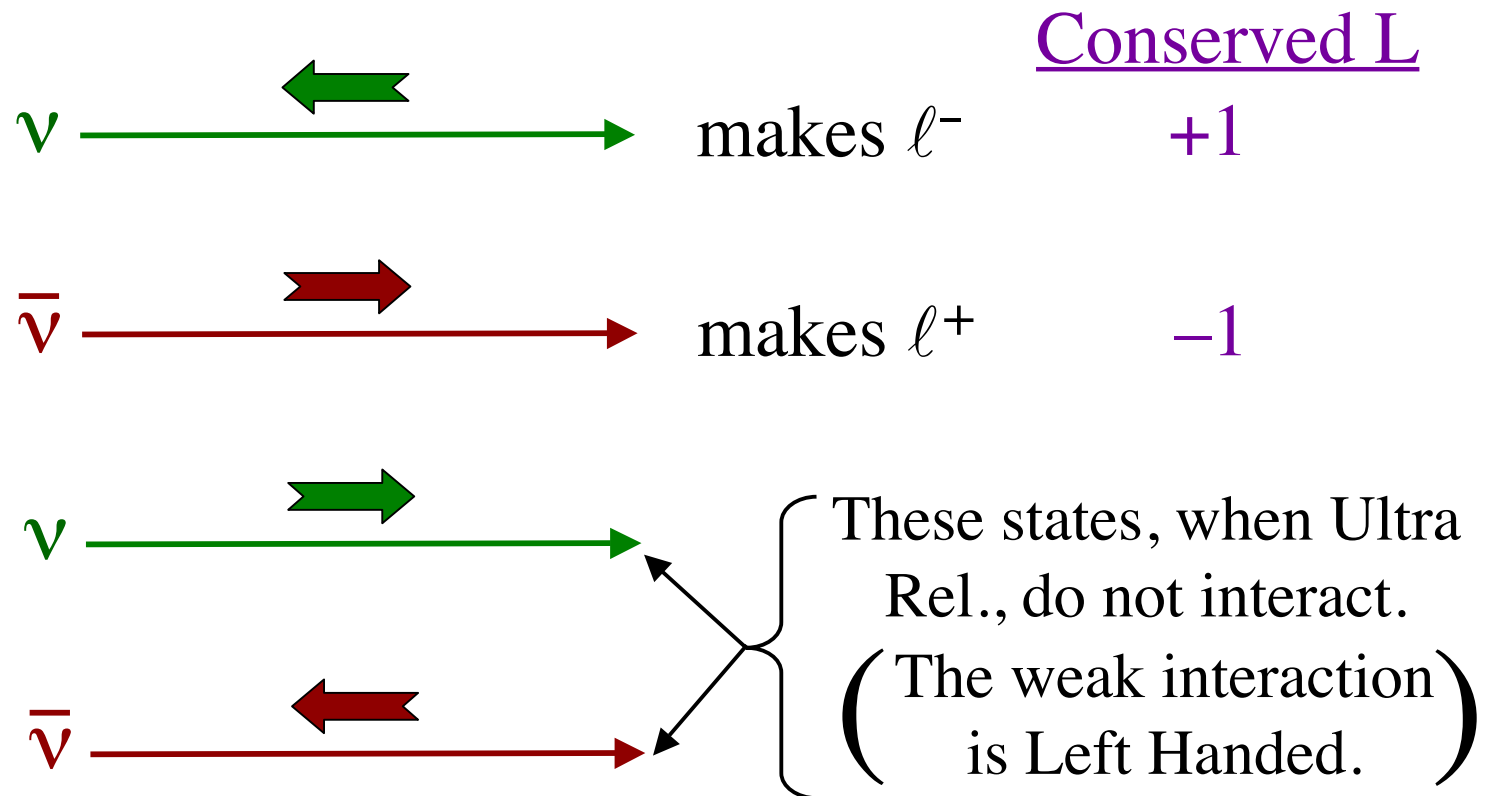
Majorana masses induce $\nu \longleftrightarrow \bar{\nu}$ mixing.

As a result of $\nu \longleftrightarrow \bar{\nu}$ mixing, the neutrino mass eigenstate is —

$$\nu_i = \nu + \bar{\nu} . \quad \bar{\nu}_i = \nu_i .$$

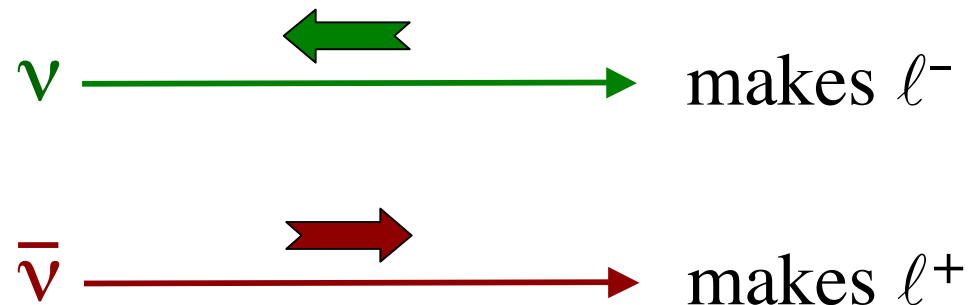
SM Interactions Of A Dirac Neutrino

We have 4 mass-degenerate states:



SM Interactions Of A Majorana Neutrino

We have only 2 mass-degenerate states:



The weak interactions violate *parity*.

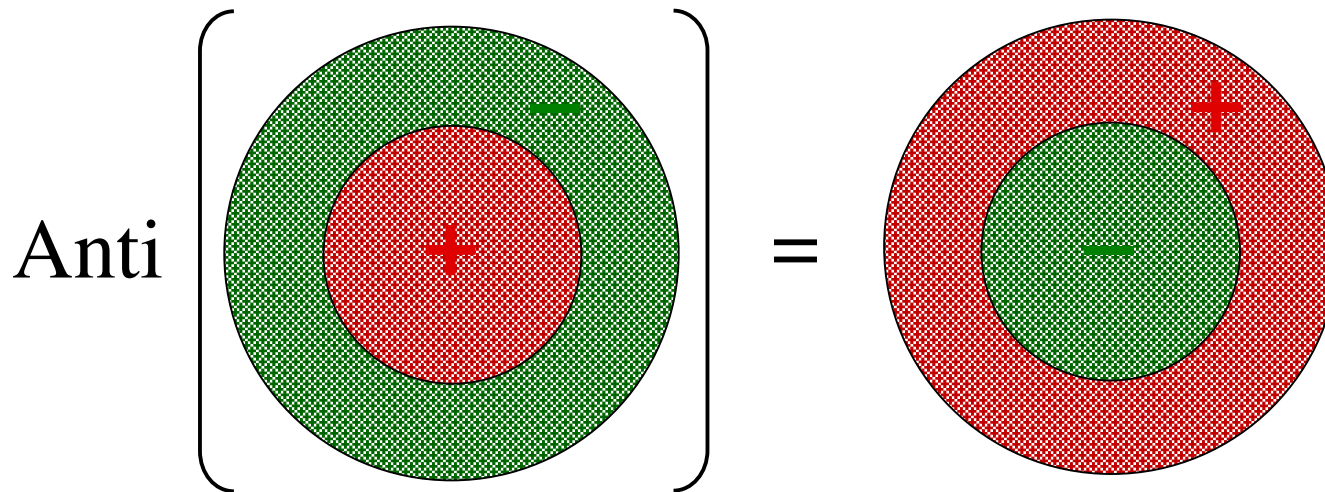
(They can tell *Left* from *Right*.)

An incoming left-handed neutral lepton makes ℓ^- .

An incoming right-handed neutral lepton makes ℓ^+ .

Can a Majorana Neutrino Have an Electric Charge *Distribution*?

No!

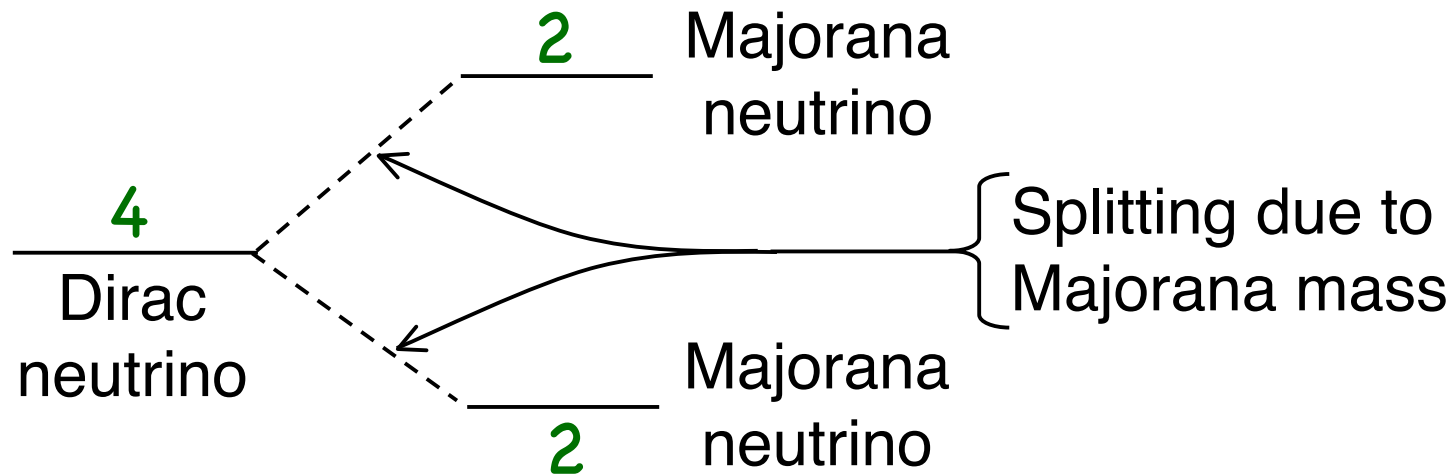


But for a Majorana neutrino —

$$\text{Anti}(\nu) = \nu$$

Majorana Masses Split Dirac Neutrinos

A Majorana mass term splits a Dirac neutrino into two Majorana neutrinos.



Why Most Theorists Expect Majorana Masses

The Standard Model (SM) is defined by the fields it contains, its *symmetries* (notably weak isospin invariance), and its renormalizability.

Leaving neutrino masses aside, anything allowed by the SM symmetries occurs in nature.

Majorana mass terms
are allowed by the SM symmetries.

Then quite likely *Majorana masses*
occur in nature too.

To Determine
Whether
Majorana Masses
Occur in Nature

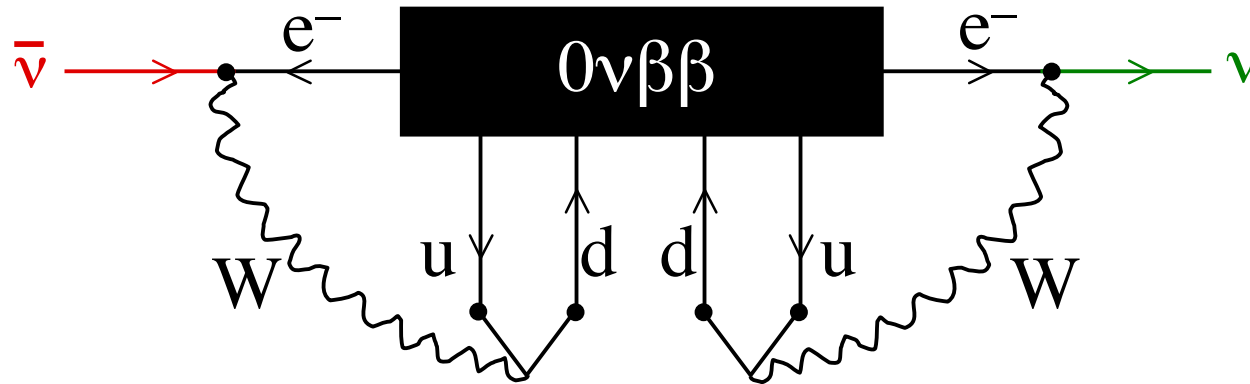
The Promising Approach — Seek Neutrinoless Double Beta Decay [$0\nu\beta\beta$]



We are looking for a *small* Majorana neutrino mass. Thus, we will need *a lot* of parent nuclei (say, one ton of them).

Whatever diagrams cause $0\nu\beta\beta$, its observation would imply the existence of a Majorana mass term:

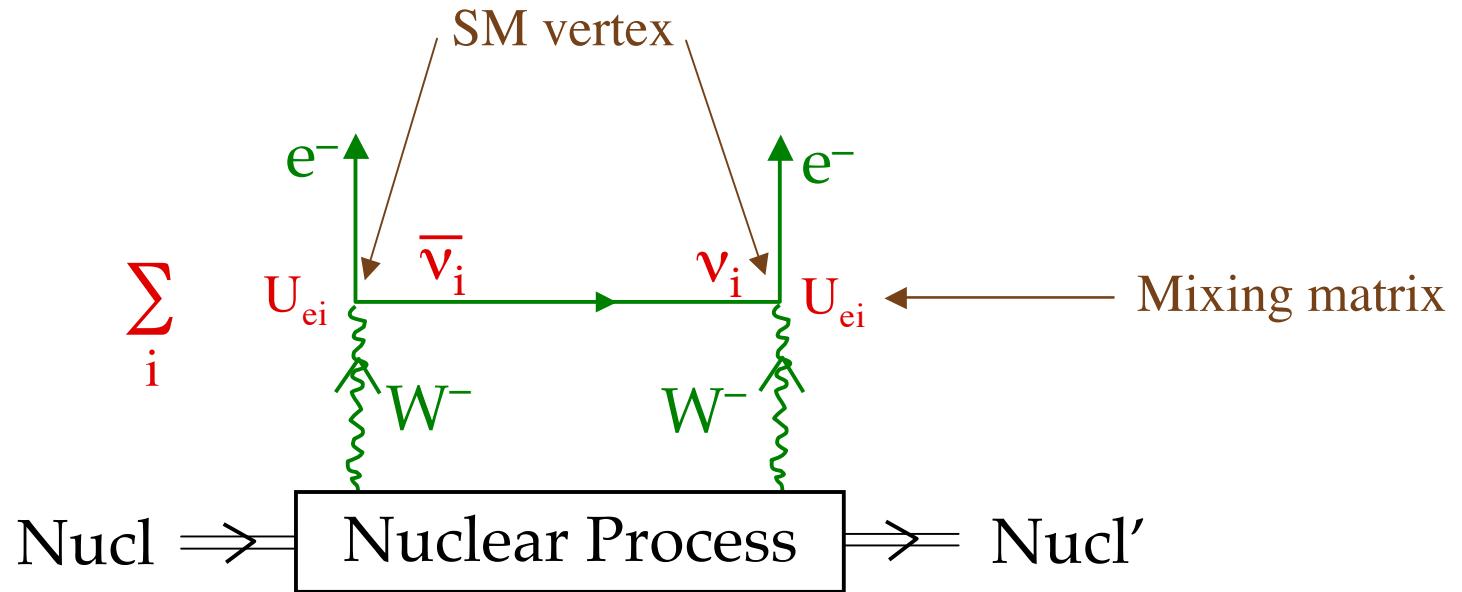
(Schechter and Valle)



$\bar{\nu} \rightarrow \nu$: A (tiny) Majorana mass term

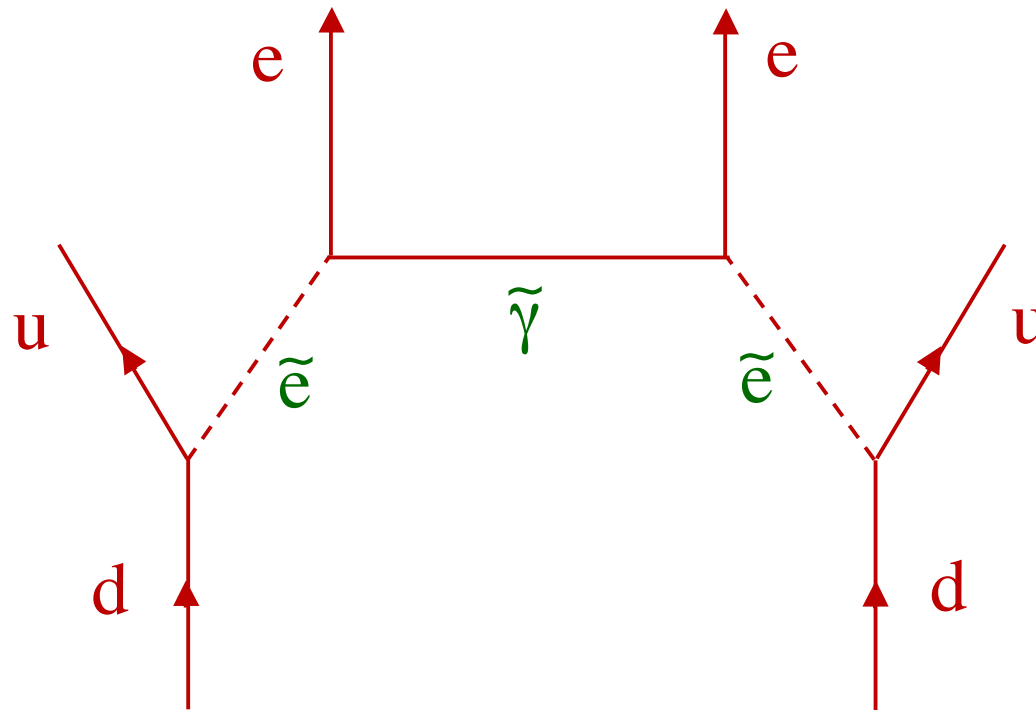
$$\therefore 0\nu\beta\beta \longrightarrow \bar{\nu}_i = \nu_i$$

We anticipate that $0\nu\beta\beta$ is dominated by a diagram with Standard Model vertices:

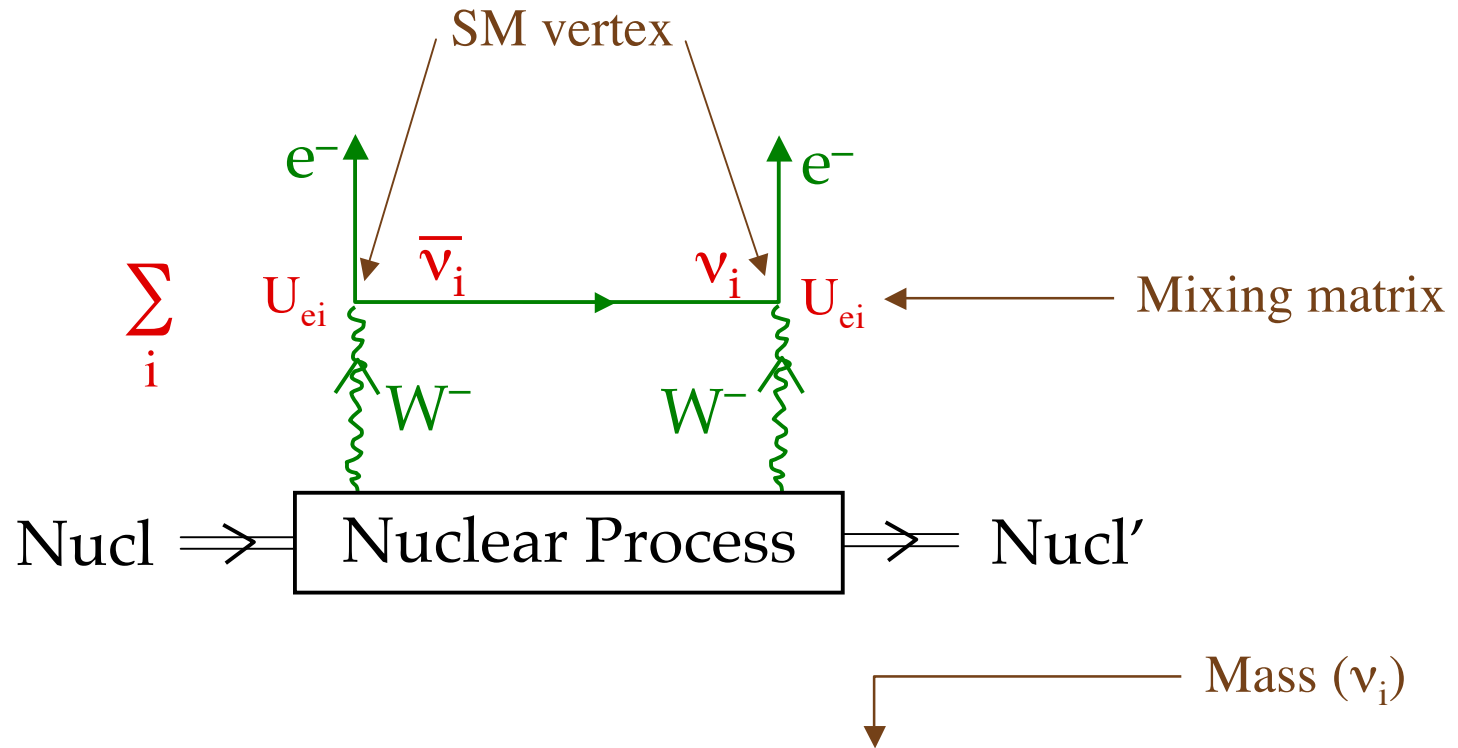


But there could be other contributions to $0\nu\beta\beta$,
which at the quark level is the process
 $dd \rightarrow uuee$.

An example from Supersymmetry:



Assume the dominant mechanism is —



The $\bar{\nu}_i$ is emitted [RH + O{ m_i/E }LH].

Thus, Amp [ν_i contribution] $\propto m_i$

$$\text{Amp}[0\nu\beta\beta] \propto \left| \sum m_i U_{ei}^2 \right| \equiv m_{\beta\beta}$$

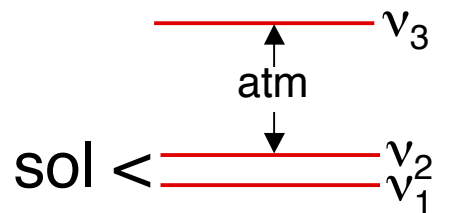
How Large is $m_{\beta\beta}$?

How sensitive need an experiment be?

Note: $\Gamma = m_{\beta\beta}^2 |\text{Nuclear M.E.}|^2 \text{Phase Space}$

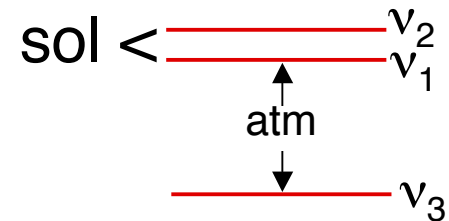
Suppose there are only 3 neutrino mass eigenstates.
(More might help.)

Then the spectrum looks like —

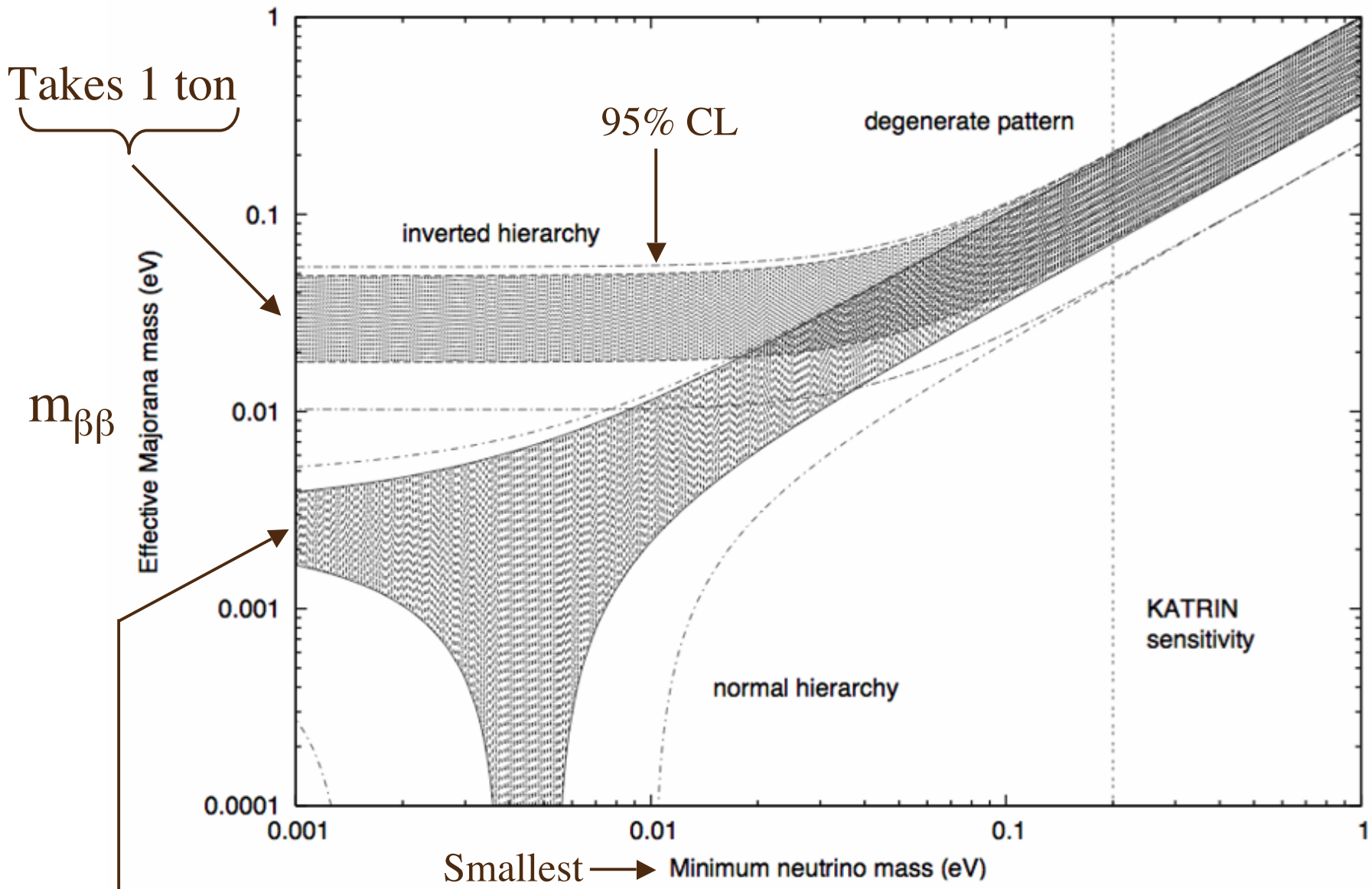


Normal hierarchy

or



Inverted hierarchy



$m_{\beta\beta}$ For Each Hierarchy

There is no clear theoretical preference
for either hierarchy.

If the hierarchy is **inverted**—

then $0\nu\beta\beta$ searches with sensitivity
to $m_{\beta\beta} = 0.01$ eV have
a very good chance to see a signal.

*Sensitivity in this range is the target
for the next generation of experiments.*

We Must Be
Alert
To *Surprises!*

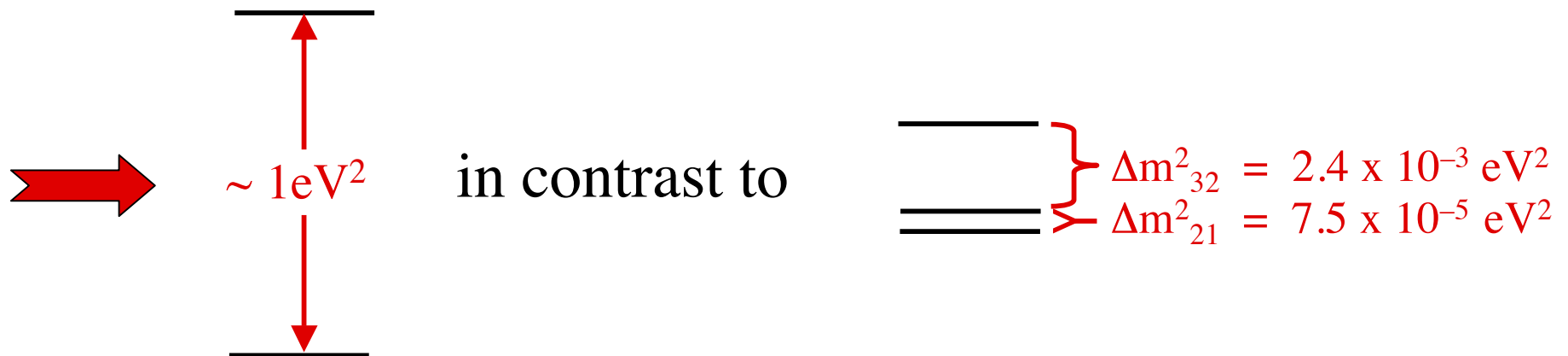
Are There
More Than 3
Mass Eigenstates?

Are There
Sterile Neutrinos?

The Hint From LSND

The **LSND** experiment at Los Alamos reported a *rapid* $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ oscillation at $L(\text{km})/E(\text{GeV}) \sim 1$.

$$P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e) = \sin^2 2\theta \sin^2 \left[1.27 \Delta m^2 (\text{eV}^2) \frac{L(\text{km})}{E(\text{GeV})} \right]$$



➡ At least **4** mass eigenstates

➡ (From the **Z** width) At least **1** sterile neutrino

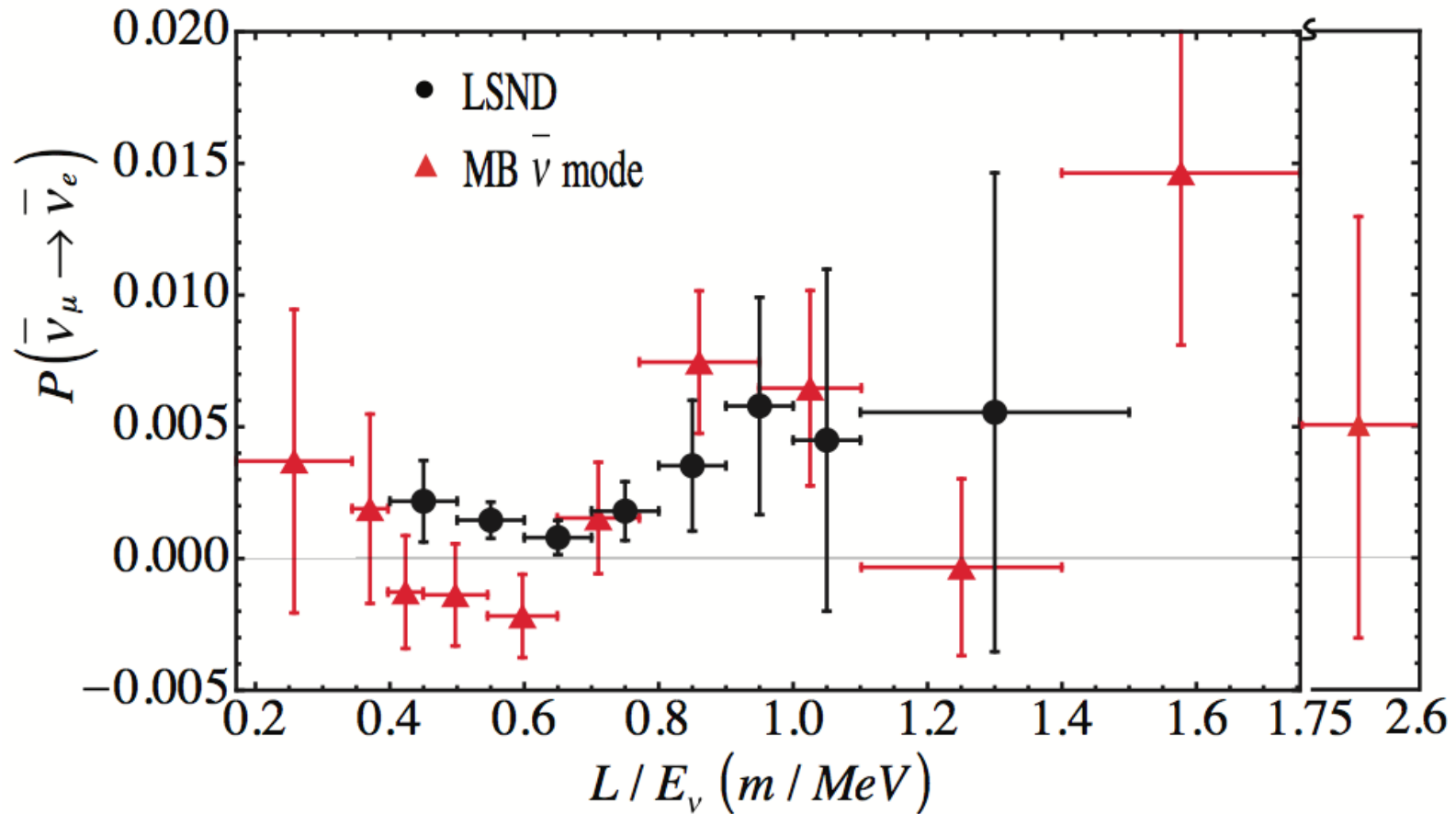
Is the LSND Signal Genuine Neutrino Oscillation?

The **MiniBooNE** experiment is trying to confirm or refute **LSND**.

In **MiniBooNE**, both L and E are ~ 17 times larger than they were in **LSND**, and L/E is comparable.

MiniBooNE has recently reported its $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ results.

Direct MiniBooNE-LSND Comparison of $\bar{\nu}$ Data



(Phys.Rev.Lett.105:181801, 2010)

Do LSND and MiniBooNE see the same thing???

The Reactor $\bar{\nu}_e$ Flux Surprise

The prediction for the un-oscillated $\bar{\nu}_e$ flux from reactors has increased by about 3%.

(Mueller et al.)

Measurements of the $\bar{\nu}_e$ flux at (10 – 100)m from reactor cores now show a $\sim 6\%$ disappearance.

(Mention et al.)

Disappearance at $L(\text{m})/E(\text{MeV}) \sim 1$ suggests oscillation with $\Delta m^2 \sim 1 \text{ eV}^2$, like LSND and MiniBooNE.

Fits to all data with 2 extra neutrinos are improved.

(Kopp et al.)

Clearly, more information is needed.

While awaiting further news –

*We will assume there are
only 3 neutrino mass eigenstates,
and no sterile neutrinos.*



Do Neutrino Interactions
Violate CP?


*Are we descended
from heavy neutrinos?*


The Challenge – A Cosmic Broken Symmetry

The universe contains baryons,
but essentially no antibaryons.

$$\frac{n_B}{n_\gamma} = 6 \times 10^{-10} \quad ; \quad \frac{n_{\bar{B}}}{n_B} \sim 0 (< 10^{-6})$$

Standard cosmology: Any initial
baryon – antibaryon asymmetry
would have been erased.

How did $n_{\bar{B}} = n_B$  $n_{\bar{B}} \ll n_B$?

Sakharov: $n_{\bar{B}} = n_B$  $n_{\bar{B}} \ll n_B$ requires \cancel{CP} .

The \cancel{CP} in the quark mixing matrix, seen in B and K decays, leads to much too small a $B-\bar{B}$ asymmetry.

If *quark* \cancel{CP} cannot generate the observed $B-\bar{B}$ asymmetry, can some scenario involving *leptons* do it?

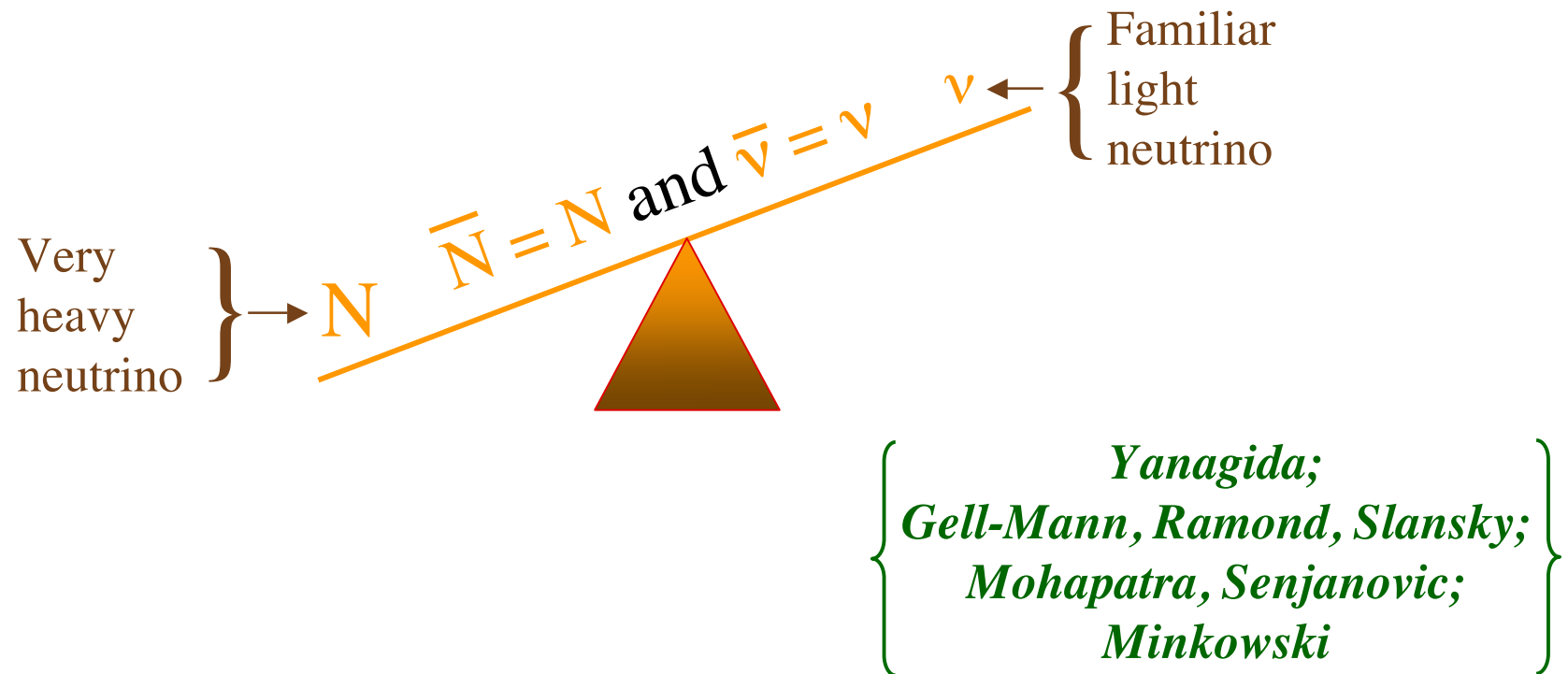
The candidate scenario: *Leptogenesis*.

(Fukugita, Yanagida)

Leptogenesis – A Two-Step Process

Leptogenesis is an outgrowth of the most popular theory of why neutrinos are so light —

The See-Saw Mechanism



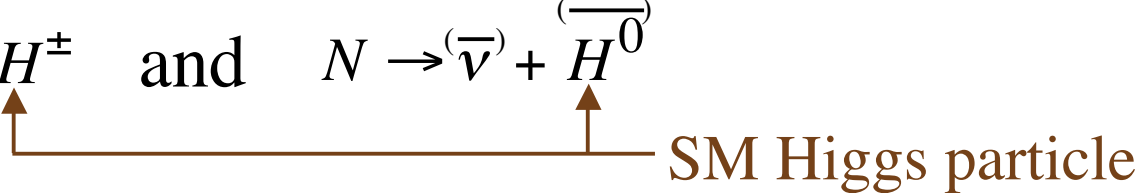
In standard leptogenesis, to account for the observed cosmic baryon – antibaryon asymmetry, *and* to explain the tiny light neutrino masses, we must have —

$$m_N \sim 10^{(9-10)} \text{ GeV} .$$

This puts the heavy neutrinos **N** far beyond LHC range.

But these heavy neutrinos would have been made in the *hot* Big Bang.

In the see-saw picture —

$$N \rightarrow \ell^{\bar{\tau}} + H^{\pm} \quad \text{and} \quad N \rightarrow (\bar{\nu}) + \overline{H^0}$$


SM Higgs particle

Assume 3 heavy neutrinos N_i to match the number (3) of light lepton (ℓ_α, ν_α) families.

By SM weak-isospin symmetry —

$$\Gamma(N_i \rightarrow \ell_\alpha^- + H^+) = \Gamma(N_i \rightarrow \nu_\alpha + H^0)$$

There are $3 \times 3 = 9$ independent “coupling constants” (lowest order decay amplitudes) $y_{\alpha i}$, forming a matrix y .

~~CP~~ phases in the matrix y will lead to —

$$\Gamma(N \rightarrow \ell^- + H^+) \neq \Gamma(N \rightarrow \ell^+ + H^-)$$

and

$$\Gamma(N \rightarrow \nu + H^0) \neq \Gamma(N \rightarrow \bar{\nu} + \overline{H^0})$$

This produces a universe with unequal numbers of **leptons** (ℓ^- and ν) and **antileptons** (ℓ^+ and $\bar{\nu}$).

In this universe the lepton number L , defined by $L(\ell^-) = L(\nu) = -L(\ell^+) = -L(\bar{\nu}) = 1$, is not zero.

This is Leptogenesis — Step 1

Leptogenesis — Step 2

The Standard-Model *Sphaleron* process, which does not conserve Baryon Number B , or Lepton Number L , but does conserve $B - L$, acts.



Initial state
from N decays

Final state

There is now a nonzero Baryon Number.

There are baryons, but ~ no antibaryons.

Reasonable parameters give the observed n_B/n_γ .

Leptogenesis and ~~CP~~ In Light ν Oscillation

In a convenient basis, the coupling matrix *y* is the only source of CP violation among the leptons.

The see-saw relation is —

$$M_\nu = -v^2 U^T \left(y^* M_N^{-1} y^\dagger \right) U$$

Heavy N mass eigenvalues

Light ν mass eigenvalues

Real number

Leptonic mixing matrix

Through U , the phases in \mathbf{y} lead to \cancel{CP} in light neutrino oscillation.

$$\begin{aligned}
 P(\overset{(-)}{\nu}_\alpha \rightarrow \overset{(-)}{\nu}_\beta) &= \\
 \text{e, } \mu, \text{ or } \tau & \begin{array}{c} \uparrow \quad \uparrow \\ \hline \end{array} \\
 &= \delta_{\alpha\beta} - 4 \sum_{i>j} \Re(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*) \sin^2\left(\Delta m_{ij}^2 \frac{L}{4E}\right) \\
 & \quad \overset{\text{Distance}}{\downarrow} \quad \uparrow \\
 & \quad \overset{\text{Neutrino (Mass)}^2 \text{ splitting}}{\uparrow} \quad \overset{\text{Energy}}{\uparrow} \\
 & \quad \overset{(+)}{\underset{(-)}{2}} \sum_{i>j} \Im(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*) \sin\left(\Delta m_{ij}^2 \frac{L}{2E}\right)
 \end{aligned}$$

*The observation of CP violation in neutrino oscillation would make it more plausible that **leptogenesis** occurred in the early universe.*

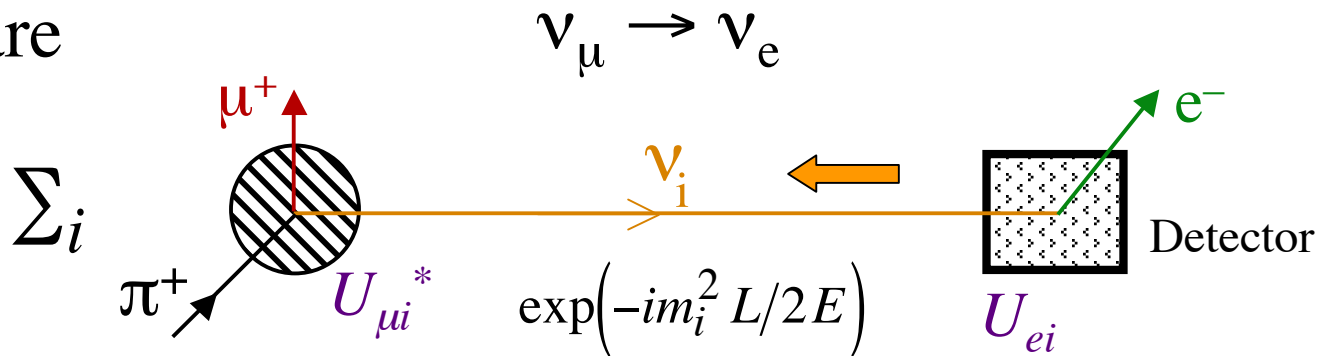
Seeking CP violation in neutrino oscillation is now a worldwide goal.

The search will use long-baseline accelerator neutrino beams to study $\nu_{\mu} \rightarrow \nu_e$ and $\bar{\nu}_{\mu} \rightarrow \bar{\nu}_e$, or their inverses.

Q : Can CP violation still lead to $P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e) \neq P(\nu_\mu \rightarrow \nu_e)$ when $\bar{\nu} = \nu$?

A : Certainly!

Compare



with

