The Exact and Approximate Symmetries of Electroweak Interactions

> Krishna Kumar, UMass Amherst Joint NNPS/TSI 2010 TRIUMF, June 27 - July 2 2010

Unique Low Energy Tests exploiting the special properties of Leptons, Nucleons and Nuclei

Outline of Lectures

- Standard Model of Electroweak Interactions
- Searches for Violations of Discrete Symmetries
- Charged Lepton Flavor Violation and Precision Weak Neutral Current Experiments

 Parity-Violating Electron Scattering Experiments, Electroweak Probes of Hadron Structure & Precision Weak Charged Current Experiments

Review and Perspective

- In lecture 1, we introduced the electroweak interaction and its verification using colliders
- In lecture 2, we learned about symmetries and discussed EDM searches to find T-violation
- In lecture 3, we discussed lepton flavor violation and weak neutral current experiments
- Coupled with Vincenzo's lecture yesterday, we are now quite familiar with the language of BSM searches and the role of symmetries at low energies

Outline of Lecture #4

Parity-violating electron scattering as a probe of new flavor diagonal amplitudes at the TeV scale

Electroweak probes of hadron structure
Precision charged current experiments
Muon g-2

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Parity-Violating Electron Scattering

Weak Neutral Current (WNC) Interactions at $Q^2 << M_z^2$

Longitudinally Polarized Electron Scattering off Unpolarized Targets

$$\sigma \alpha | A_{\gamma} + A_{weak} |^2$$



$$-A_{LR} = A_{PV} = \frac{\sigma_{\uparrow} - \sigma_{\downarrow}}{\sigma_{\downarrow} + \sigma_{\downarrow}} \sim \frac{A_{weak}}{A_{\gamma}} \sim \frac{G_F Q^2}{4 \pi \alpha} (g_A^{\ e}g_V^{\ T} + \beta g_V^{\ e}g_A^{\ T})$$

and g_A are function of $sin^2\theta_W$ $\qquad A_{PV} \sim 10^{-5} \cdot Q^2$ to $10^{-4} \cdot Q^2$

Specific choices of kinematics and target nuclei probes different physics:

• In mid 70s, goal was to show $sin^2\theta_W$ was the same as in neutrino scattering

Early 90s: target couplings carry novel information about hadronic structure

 Now: precision measurements with carefully chosen kinematics can probe physics at the multi-TeV scale

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 g_V

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Experimental Technique



: copper

: quartz

electron flux

phototube

Optical pumping of a GaAs wafer
Rapid helicity reversal: change sign of longitudinal polarization ~ 100 Hz to minimize drifts (like a lockin amplifier)
Control helicity-correlated beam motion: under sign flip, keep beam stable at the submicron level

"Flux Integration": very high rates direct scattered flux to background-free region

Technical progress over 3 decades has enabled ppb systematic control

integrator

Parity-violating electron scattering has become a precision tool: Many-body nuclear physics: Neutron skin of ²⁰⁸Pb

- Nucleon structure: strangeness contribution to form factors
- Valence quark structure: Deep inelastic scattering at high-x
 Search for new TeV physics: Precision electroweak parameters Krishna Kumar
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Parity-Violating Electron-Electron (Møller) Scattering E158 @ SLAC



50 GeV at SLAC: ~ 150 ppb!

45 & 48 GeV Beam 85% longitudinal polarization



Purely leptonic reaction

4-7 mrad



End Station A at the Standord Linear Accelerator Center (SLAC)





g-2 spin precession
45 GeV: 14.0 revs
48 GeV: 14.5 revs



 A_{PV} = (-131 ± 14 ± 10) x 10⁻⁹ NNPSS 2010 Lecture 4

E158 Implications



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Neutrino Deep Inelastic Scattering

NuTeV Neutrino Experiment



- Most precise
 measurement of
 neutrino-quark
 coupling
- subtle quark
 physics effects
 can affect the
 result
- generated great interest in both nuclear and particle phenomenology

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Lepton-Quark Neutral Current Interactions

Conside

sider	$f_1 \bar{f}_1 \twoheadrightarrow f_2 \bar{f}_2$	or	$f_1 f_2 \rightarrow f_1 f_2$	f_1
$L_{f_1 f_2} =$	$= \sum_{i,j=L,R} \frac{4\pi}{\Lambda_{ij}^2} \eta_{ij} j$	$\bar{f}_{1i}\gamma_{\mu}f_{1}$	$_i\bar{f}_{2j}\gamma^\mu f_{2j}$	f_2



$$\delta(C_{1q}) \propto (+\eta_{RL}^{eq} + \eta_{RR}^{eq} - \eta_{LL}^{eq} - \eta_{LR}^{eq})$$

$$\delta(C_{2q}) \propto (-\eta_{RL}^{eq} + \eta_{RR}^{eq} - \eta_{LL}^{eq} + \eta_{LR}^{eq})$$

PV elastic e-p scattering, APV
 PV deep inelastic scattering



Qweak at Jefferson Lab



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Qweak at Jefferson Laboratory

A_{PV} in elastic *e*-*p* scattering:

$$A(Q^2 \to 0) = -\frac{G_F}{4\pi\alpha\sqrt{2}} \left[Q^2 Q_{weak}^p + Q^4 B(Q^2) \right]$$
$$Q_{weak}^p = 2C_{1u} + C_{1d} \propto 1 - 4\sin^2\vartheta_W$$

 $E = 1.165 \ GeV, \ \theta_{lab} \sim 9^{o},$ $Q^{2} = 0.026 \ GeV^{2}$

86 scientists from 25 institutions including U. Manitoba and TRIUMF

Design and construction over past several years
Installation nearly complete
First beam next week!
Data ~ 2010 thru mid-2012

Contains $G^{\gamma}_{E,M}$ and $G^{Z}_{E,M}$, Extracted using global fit



New, complementary constraints on lepton-quark interactions at the TeV scale

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Atomic Parity Violation

•6S → 7S transition in ¹³³Cs is forbidden within QED
•Parity Violation introduces small opposite parity admixtures
•Induce an E1 Stark transition, measure E1-PV interference
•5 sign reversals to isolate APV signal and suppress systematics
•Signal is ~ 6 ppm, measured to 40 ppb



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Electroweak and Hadron Physics Interplay

nuclei and nucleons are special laboratories to test electroweak interactions

these are bound states ultimately governed by QCD dynamics

a detailed knowledge of hadron dynamics is often needed to interpret the measurement and probe the TeV scale

conversely, the experimental techniques being developed lead to new insights on hadron structure. Some classic examples:

elastic electron-nucleon and -nuclear scattering

nuclear beta decay and muon capture
 Very
 Lamb Shift (1s-2s transition)
 in muonic hydrogen
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Strange Quarks in the Nucleon



Late 1980's

Strange quarks carry nucleon momentum: Other external properties affected?





PREX at Jefferson Laboratory

- Neutron star has solid crust over liquid core.
- Heavy nucleus has neutron skin. crust



 $R_p \sim 5.5 \ fm$ $R_n - R_p \sim 0.1 \ to \ 0.3 \ fm?$

Both neutron skin and neutron star crust are made out of neutron rich matter at similar densities.

•A neutron skin is expected: how thick is it?

•The extent of the skin constrains the transition density from solid crust to liquid core in a neutron star

•The density dependence of the symmetry energy constrains the composition of the neutron star core: important implications for rate of neutron star cooling

An experimental clean measurement of skin is now viable:

PREX and Neutron Stars



A technically demanding measurement:Installation took place
over the last 6 monthsRate ~ 2 GHzSeparate excited state at 2.4 MeVPhysics run April - mid JuneStat. Error ~ 15 ppbexpect to have < 3% measurement of neutron radius</td>

Result highly prized by nuclear structure and nuclear astrophysics communities

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MOLLER at Jefferson

Laboratory Measurement of Lepton-Lepton Electroweak Reaction

A_I(SLD) $E_{beam} = 11 \text{ GeV} \quad 75 \ \mu A \quad 80\% \text{ polarized} \qquad \xrightarrow{\sim 38 \text{ weeks}} \delta(A_{PV}) = 0.73 \text{ ppb}$ $A_{fb}^{0,b}$ $\delta(Q^{e}_{W}) = \pm 2.1 \text{ (stat.) } \pm 1.0 \text{ (syst.) }\%$ $A_{PV} = 35.6 \ ppb$ Average 10[°] $\delta(\sin^2\theta_W) = \pm 0.00026 \text{ (stat.)} \pm 0.00012 \text{ (syst.)} \implies \sim 0.1\%$ [GeV] Project design, construction and $\blacksquare \square \alpha_{\text{bod}}^{(5)} = 0.02758 \pm 0.00035$ installation will take 4-5 years 0.23 0.232 lept Jefferson Lab 12 GeV Upgrade $\mathcal{L}_{e_1e_2} = \sum_{\mathbf{a}_1 \in \mathbf{a}_2} \frac{\mathbf{g}_{\mathbf{ij}}^2}{2\Lambda^2} \mathbf{\bar{e}_i} \gamma_{\mu} \mathbf{e_i} \mathbf{\bar{e}_j} \gamma^{\mu} \mathbf{e_j} \qquad \stackrel{A_{PV}}{\longrightarrow} \frac{\Lambda}{\sqrt{|\mathbf{g}_{RR}^2 - \mathbf{g}_{LL}^2|}} = 7.5 \text{ TeV}$ Best current limits on 4-electron contact interactions: LEPII at 200 GeV $\frac{\Lambda}{\sqrt{|\mathbf{g}_{\mathbf{RR}}^2 + \mathbf{g}_{\mathbf{LL}}^2|}} = 4.4 \text{ TeV} \quad \mathbf{OR} \quad \frac{\Lambda}{\mathbf{g}_{\mathbf{RL}}} = 5.2 \text{ TeV}$ insensitive to $|\mathbf{g}_{RR}^2 - \mathbf{g}_{LL}^2|$

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 ± 0.00029

 0.23098 ± 0.00026

 0.23221 ± 0.00029

 0.23153 ± 0.00016

 χ^2 /d.o.f.: 11.8 / 5

0.234

Precision Tests of the Weak Charged Current

$$L_{CC} = \frac{g}{2\sqrt{2}} W^+_{\mu} \left[\overline{U_i} \gamma^{\mu} (1 - \gamma^5) \mathbf{V_{ij}} D_j + \overline{\nu_k} \gamma^{\mu} (1 - \gamma^5) l_k \right] + \text{h.c.}$$

$$M^{CC} \approx \frac{g^2}{8M_W^2} (V - A) \otimes (V - A)$$

$$U = \begin{pmatrix} u \\ c \\ t \end{pmatrix} \qquad D = \begin{pmatrix} d \\ s \\ b \end{pmatrix} \qquad \nu = \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} \qquad l = \begin{pmatrix} e \\ \mu \\ \tau \end{pmatrix}$$

Fermi Constants

 $\mu \text{ decay} \qquad \frac{G_F^{\mu}}{\sqrt{2}} = \frac{g^2}{8M_w^2} \qquad \left(1 + \Delta r_{\mu}\right)$

 $\mathbf{V} = \left(egin{array}{ccc} V_{ud} & V_{us} & V_{ub} \ V_{cd} & V_{cs} & V_{cb} \ V_{ud} & V_{ud} & V_{ud} \end{array}
ight)$

$$\mathbf{V}\mathbf{V}^{\dagger} = \mathbf{V}^{\dagger}\mathbf{V} = I$$
 \downarrow
 $\sum_{eta} V_{lphaeta} V_{\gammaeta}^{*} = \delta_{lpha\gamma}$

 $\beta \text{ decay} \qquad \frac{G_F^{\beta}}{D} = \frac{g^2}{8M_{\odot}^2} V_{ud} \left(1 + \Delta r_{\beta}\right)$

 $g^2/8M_W^2$ is universal

New physics

Universality obscured by $G_F^{\beta} / G_F^{\mu} = V_{ud} \left(1 + \Delta r_{\beta} - \Delta r_{\mu} \right)$

Super-Allowed Beta



Beta Decay Correlation Coefficients



$$\mathcal{N}(E_e) \left\{ 1 + a \frac{\vec{p_e} \cdot \vec{p_\nu}}{E_e E_\nu} + b \frac{\Gamma m_e}{E_e} + \langle \vec{J} \rangle \cdot \left[A \frac{\vec{p_e}}{E_e} + B \frac{\vec{p_\nu}}{E_\nu} + D \frac{\vec{p_e} \times \vec{E_e}}{E_e E_\nu} \right] \right\}$$

Jackson, Treiman, Wyld

• Unique sensitivity to S,T operators (via interference terms $\propto m_e/E_e$)

 Example: limit on b from 0⁺ → 0⁺ transitions corresponds to

$$\frac{\Lambda}{\sqrt{C_S}} \ge 7 \,\mathrm{TeV}$$



 $(Z\alpha)^2$

 $\Gamma = \sqrt{2}$

The Neutron

 $rac{{
m dw}}{{
m dE_e}{
m d}\Omega_{
m e}{
m d}\Omega_{
u}}\simeq {
m k_e}{
m E_e}({
m E_0}-{
m E_e})^2$ $\times \left[1 + \mathbf{a} \frac{\vec{\mathbf{k}}_{e} \cdot \vec{\mathbf{k}}_{\nu}}{\mathbf{E}_{e} \mathbf{E}_{\nu}} + \mathbf{b} \frac{\mathbf{m}}{\mathbf{E}_{e}} + \langle \vec{\sigma}_{n} \rangle \cdot \left(\mathbf{A} \frac{\vec{\mathbf{k}}_{e}}{\mathbf{E}_{e}} + \mathbf{B} \frac{\vec{\mathbf{k}}_{\nu}}{\mathbf{E}_{\nu}} + \mathbf{D} \frac{\vec{\mathbf{k}}_{e} \times \vec{\mathbf{k}}_{\nu}}{\mathbf{E}_{e} \mathbf{E}_{\nu}} \right) \right]$

with:

$=rac{1- oldsymbol{\lambda} ^2}{1+3 oldsymbol{\lambda} ^2}$	$A = -2rac{ m{\lambda} ^2 + \operatorname{Re}(m{\lambda})}{1 + 3 m{\lambda} ^2}$
_	

$$\mathsf{B} = 2 \frac{|\boldsymbol{\lambda}|^2 - \mathsf{Re}(\boldsymbol{\lambda})}{1 + 3|\boldsymbol{\lambda}|^2} \qquad \mathsf{D} = 2 \frac{\mathsf{Im}(\boldsymbol{\lambda})}{1 + 3|\boldsymbol{\lambda}|^2}$$

 $\lambda = \frac{G_A}{G_M}$ (with $\tau_n \Rightarrow \text{CKM } V_{ud}$)

a

Independent measure of Vud Other terms sensitive TeV scale BSM physics

Major initiatives in Canada, USA, Europe and Japan



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Semi-Leptonic Decays



Global Fit to V_{ud} and V_{us}



Fit result

$$V_{ud} = 0.97425 (22)$$

 $V_{us} = 0.2252 (9)$

$$\chi^2/dof = 0.65/1$$

$$|V_{ud}|^2 + |V_{us}|^2 = 0.9999(6)$$

Error equally shared between V_{ud} and V_{us}

Remarkable agreement with Cabibbo universality: $\Delta_{CKM} = -(1 \pm 6) * 10^{-4}$

Marciano-Sirlin

- Confirms large EW rad. corr. $(2 \alpha/\pi \log(M_Z/M_p) = +3.6\%)$
- It would naively fit $M_Z = (90 \pm 7) \text{ GeV}$ Krishna Kumar 23

Precision Muon Decay Measurements

TWIST at TRIUMF



^{(©} Muon decay ("Michel") parameters ρ , η, **P**_μξ, δ

$$\begin{split} & \textcircled{O} \text{ muon differential decay rate } vs. \text{ energy and angle:} \\ & \frac{d^2\Gamma}{dx\;d\cos\theta} \;=\; \frac{1}{4}m_\mu W^4_{\mu e}G^2_F\sqrt{x^2-x_0^2} \cdot \\ & \quad \{\mathcal{F}_{IS}(x,\rho,\eta)+\mathcal{P}_\mu\cos\theta\cdot\mathcal{F}_{AS}(x,\boldsymbol{\xi},\boldsymbol{\delta})\}+R.C. \end{split}$$

TWIST Results



Important new limits on righthanded currents

 $W_L = W_1 \cos \zeta + W_2 \sin \zeta, \quad W_R = e^{i\omega} (-W_1 \sin \zeta + W_2 \cos \zeta)$







Muon g-2

$$\vec{\mu} = g \frac{e}{2mc} \vec{s}, \quad \vec{s} = \frac{\hbar}{2} \vec{\sigma}$$

- Dirac predicts in g=2 in 1928
- 1947 : Measurements of Kusch and Foley found g_e deviates from 2
- Schwinger calculated :



Radiative Corrections



g-2 Experiment at BNL



• Inject polarized muons at 3.094 GeV/c into superferric storage ring, radius = 711.2 cm

ullet Muon spin precesses in homogeneous 1.45 T field, time dilated lifetime of 64.4 μ s, measure for 700 μ s

$$\vec{\omega}_{a} = \vec{\omega}_{s} - \vec{\omega}_{c} \qquad : \qquad \text{difference between spin and cyclotron frequencies} \\ \vec{\omega}_{a} = -\frac{q}{mc} \left[a_{\mu} \vec{B} - \left(a_{\mu} - \frac{1}{\gamma^{2} - 1} \right) \vec{\beta} \times \vec{E} \right] \qquad \Rightarrow \quad \text{at} \quad \gamma = 29.3 \quad \Rightarrow \quad \vec{\omega}_{a} = -\frac{q}{mc} \left[a_{\mu} \vec{B} \right] \\ \Rightarrow \quad \text{To determine } a_{\mu}, \text{ need to measure } \omega_{a} \text{ and } B$$

• Muon spin direction correlated with decay electron direction, and $E_{\text{lab}} \approx \gamma \ E^* \left(1 + \cos \theta^*\right)$

g-2 Data and Plans



Summary

- Weak Neutral Current Measurements are an important complement to search for TeV-scale flavor diagonal interactions
- Electroweak experiments probe novel aspects of hadron structure
- Charged Current Interactions search for new physics in sectors often not accessible at colliders
- muon g-2 is a very sensitive indirect search for SUSY

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Conclusions and Perspectives

I hope you had at least some fraction of the fun I had in preparing these lectures

Nuclear theory & experiments will continue to explore fundamental symmetries and interpret/complement collider experiments uncover the underlying theories of nature

We need you all to develop the next generation of clever ideas that will move this subfield forward into the future

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