

Nuclear & Particle Physics of Compact Stars

Madappa Prakash
Ohio University, Athens, OH

National Nuclear Physics Summer School

July 24-28, 2006, Bloomington, Indiana

The Role of the Equation of State in Binary Mergers

Madappa Prakash
Ohio University, OH

&

PALS

Saša Ratković (London)

James M. Lattimer (Stony Brook University)

PALS



Jim

Sasa

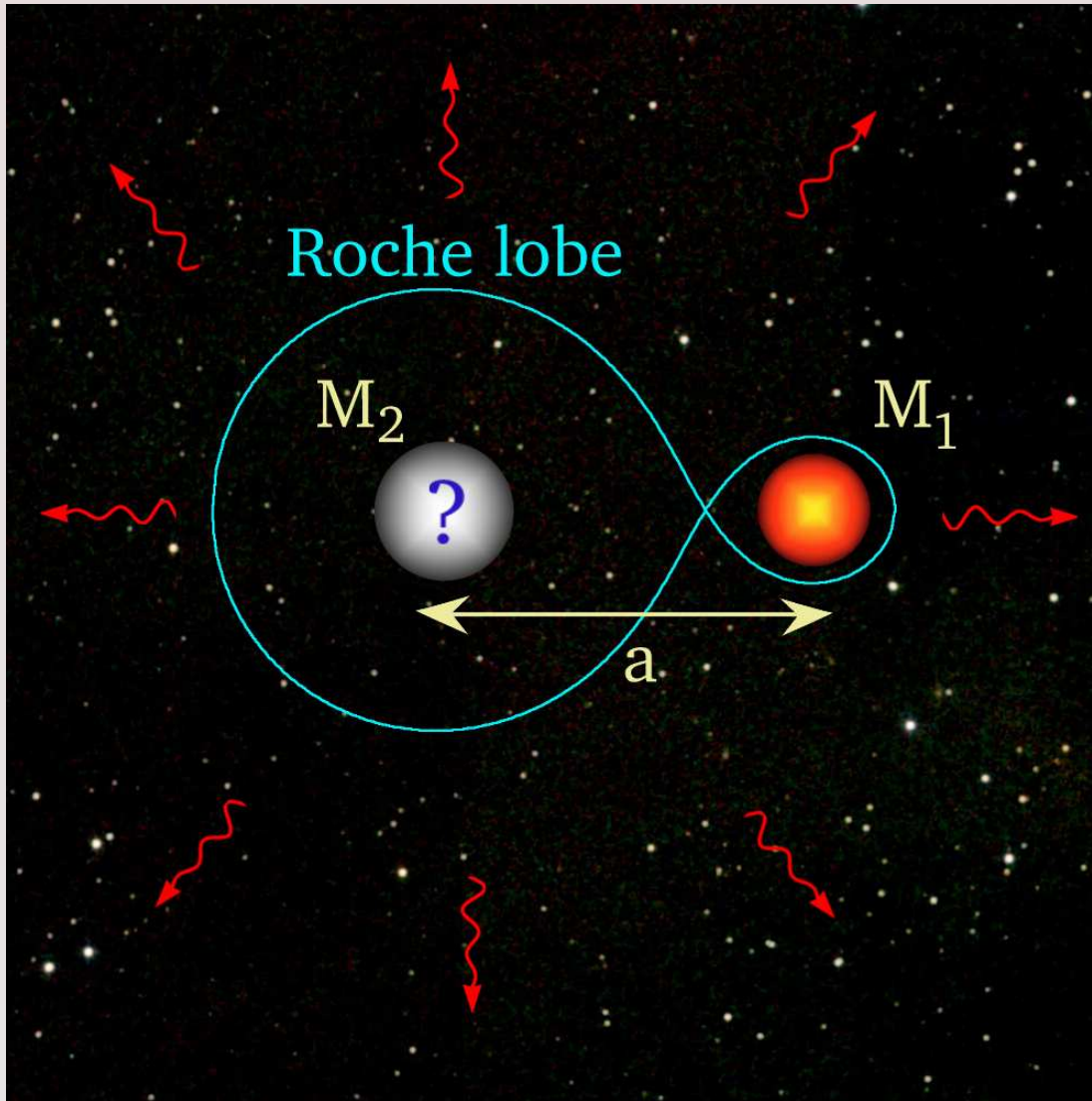


Our Thoughts on this Subject

1. M. Prakash,
Jl. Phys. G.: Nucl. Part. Phys. **30**, S451 (2003)
2. S. Ratkovic, M. Prakash & J. M. Lattimer,
Jl. Phys. G.: Nucl. Part. Phys. **30**, S1279 (2004)
3. S. Ratkovic, M. Prakash & J. M. Lattimer,
astro-ph/0512133; 0512136
ApJ (2006), To be published.

The Binary Merger Experience

The Ultimate Heavy-Ion Collision



- ▶ $M_1 \leq M_2$
- ▶ radial separation: $a(t)$
- ▶ M_1 - *NS* or *SQM*
- ▶ M_2 - *BH*, *NS*, ...
- ▶ GW emission \Rightarrow

$$\begin{aligned} L_{GW} &= \frac{1}{5} \frac{G}{c^5} \langle \ddot{I}_{jk} \ddot{I}_{jk} \rangle \\ &= \frac{32}{5} \frac{G^4}{c^5} \frac{M^3 \mu^2}{a^6} \end{aligned}$$

orbit shrinks

- ▶ Mass transfer

Einstein's General Relativity

$$G^{\alpha\beta} [g, \partial g, \partial^2 g] = 8\pi T^{\alpha\beta} [g]$$

- $G^{\alpha\beta}$: 2^{nd} -order nonlinear differential operator acting on $g_{\alpha\beta}$
- $T^{\alpha\beta}$: Stress-energy tensor of matter fields

Parametrized Post-Newtonian (PPN) Formulation

In weak field limit,

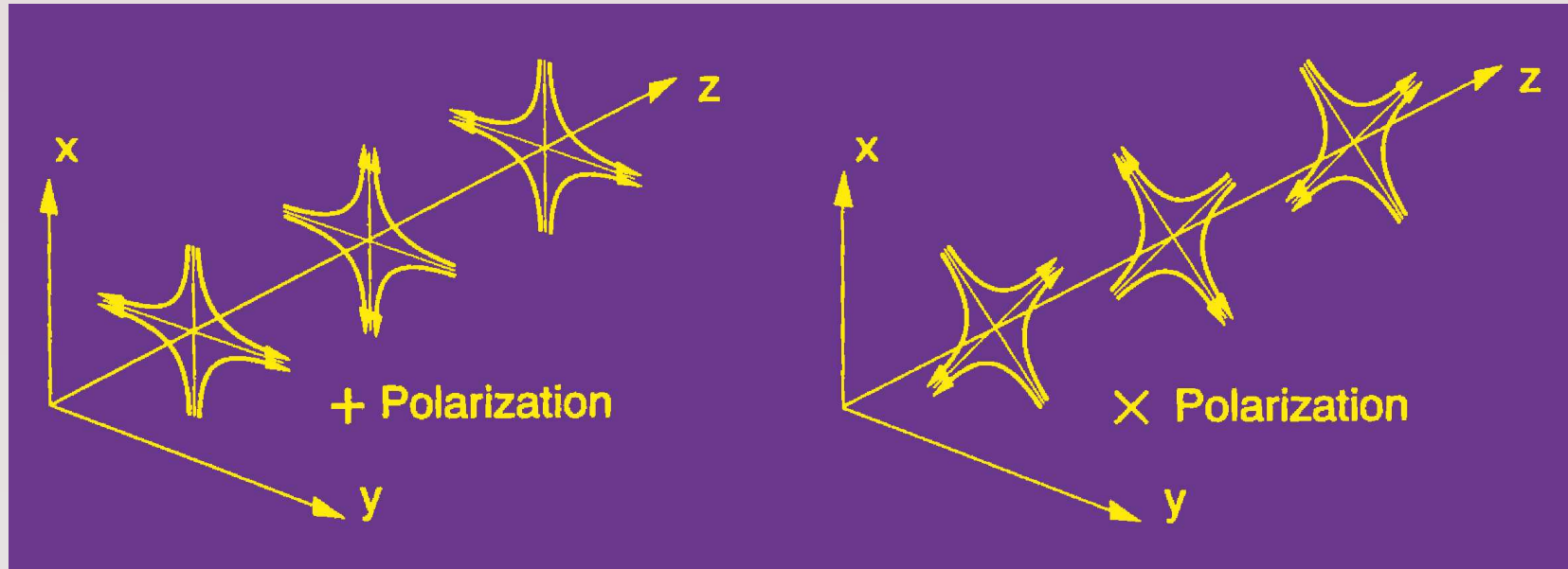
$$g_{\mu\nu}^{PPN} = \eta_{\mu\nu} + h_{\mu\nu}^{1PN}(M) + h_{\mu\nu}^{2PN}(M) + h_{\mu\nu}^{3PN}(M) + \dots$$

- $\eta_{\mu\nu}$: flat-space Minkowski metric
- M : incorporates dependence on matter fields
- $1PN, 2PN, \dots \Rightarrow [\mathcal{O}(v^2/c^2)]^\epsilon$ with $\epsilon = 1, 2, \dots$

For vacuum gravitational fields (in transverse traceless gauge),

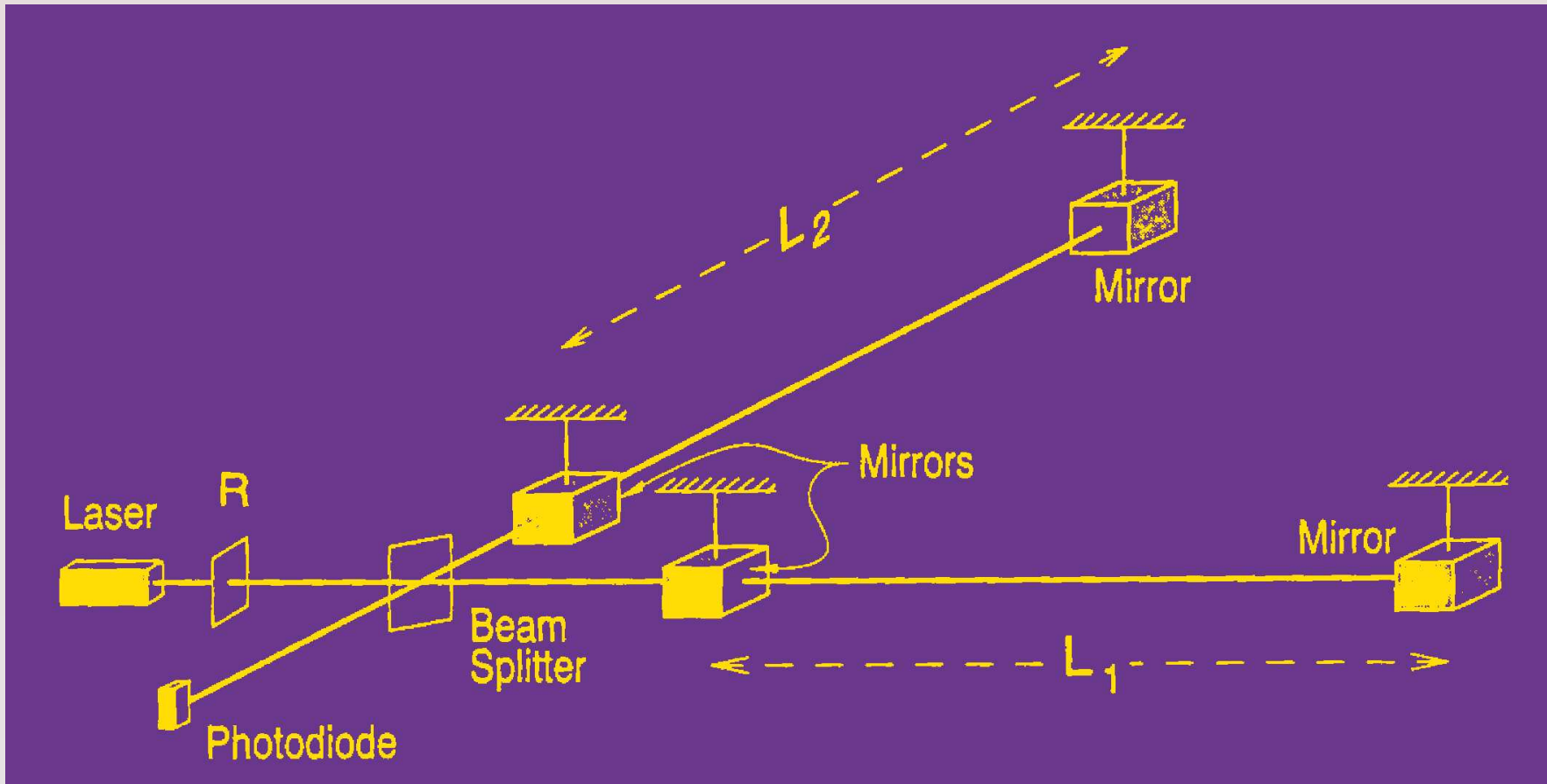
$$\left(-\frac{\partial^2}{\partial t^2} + \nabla^2 \right) h_{\times/+} = 0$$

GW Lines of Force



GW's have two transverse polarizations, h_+ & h_x .

Laser Interferometer GW Detector



For a readable account,
see K. Thorne, arXiv:gr-qc/9506084

Gravitational Wave Detection

► GW Strain : $h(t) = F_{\times} h_{\times}(t) + F_{+} h_{+}(t)$

- $F_{\times,+}$: Constants of order unity

- $h_{\times,+} \sim \frac{\delta L}{L_0} \sim \frac{1}{c^2} \frac{4G(E_{kin}^{ns}/c^2)}{r}$: Gravitational waveforms

- L_0 : Unperturbed length of detector arm

- δL : Relative change in length

- E_{kin}^{ns} : Nonspherical part of the internal kinetic energy

- ELF : 10^{-15} - 10^{-18} Hz VLF : 10^{-7} - 10^{-9} Hz*

- LFB : 10^{-4} Hz - 1 Hz, HFB : 1 Hz - 10^4 Hz

► Astrophysical Sources Radiating GW's in the HFB

Supernovae	at 10 Mpc	$h \geq 10^{-25}$
Supernovae	Milky Way	$h \sim 10^{-18}$
1.4M _⊙ NS Binaries	at 10 Mpc	$h \sim 10^{-20}$
10M _⊙ BH Binaries	at 150 Mpc	$h \sim 10^{-20}$

Inspiral Waveform

Chirp signal:

$$h_+ \propto \frac{\mathcal{M}^{5/3}}{r} f^{2/3} \cos(2\pi ft)$$

$$h_\times \propto \frac{\mathcal{M}^{5/3}}{r} f^{2/3} \sin(2\pi ft)$$

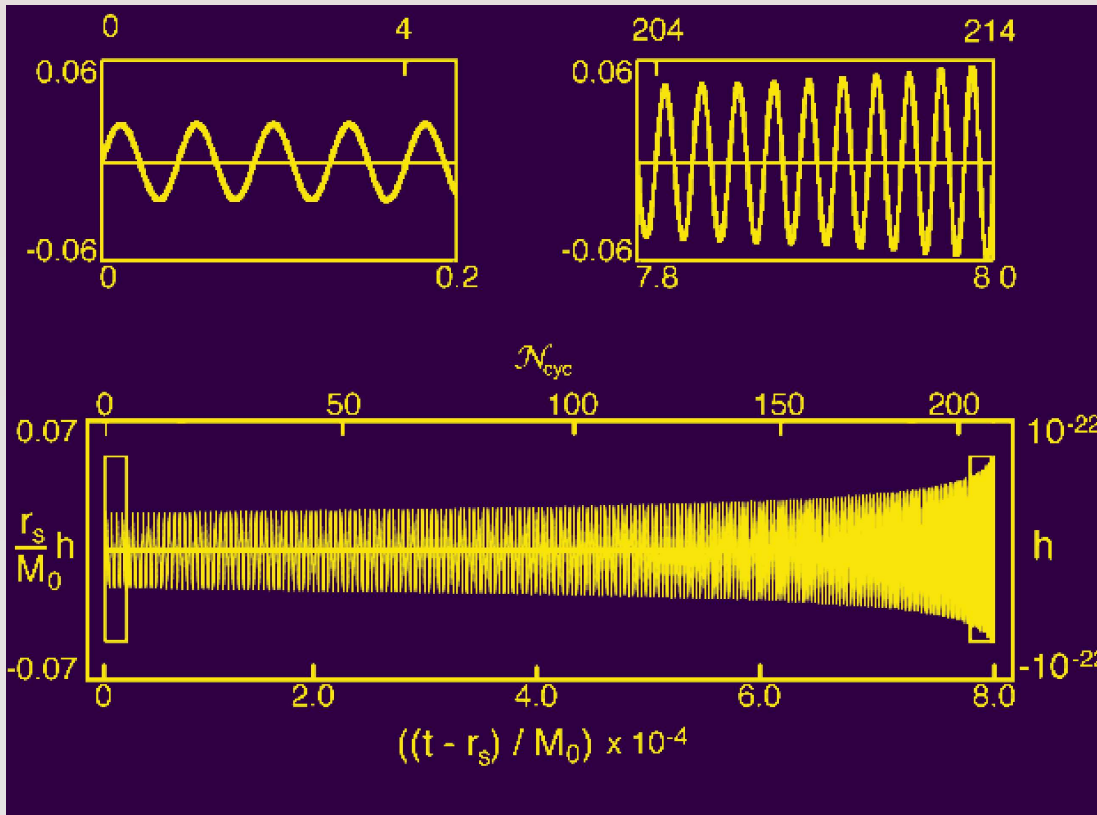
$$f = K_0 \mathcal{M}^{-5/8} (t_c - t)^{-3/8}$$

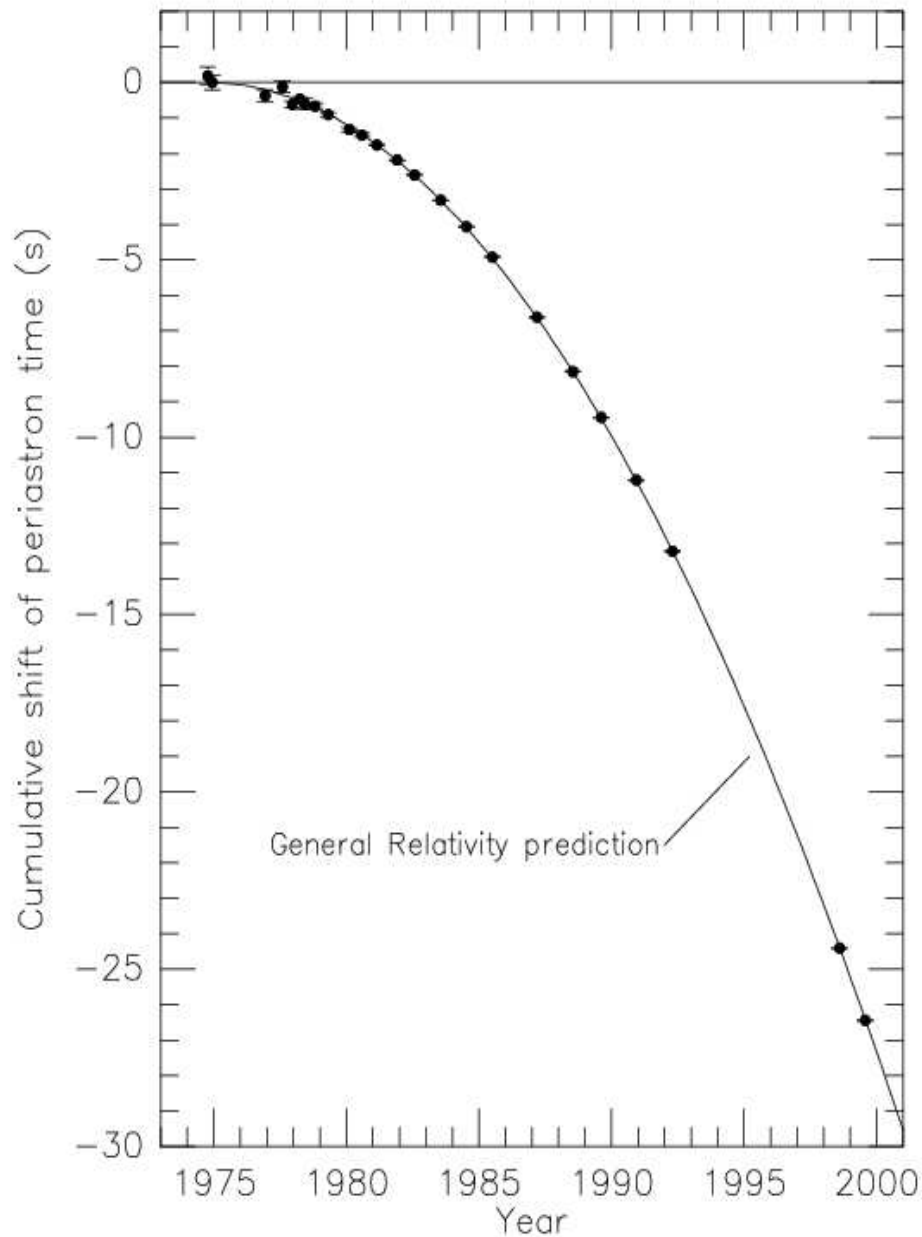
with the “chirp mass”:

$$\mathcal{M} = (M_1 M_2)^{3/5} M^{-1/5}$$

and the constant:

$$K_0 = \frac{5^{3/8}}{8\pi} \left(\frac{c^3}{G} \right)^{5/8}$$





- ▶ Binary pulsar PSR 1913+16
- ▶ Period: 7 h 45 min
- ▶ $M_{NS} = 1.4408 \pm 0.0003 M_{\odot}$
- ▶ $M_c = 1.3873 \pm 0.0003 M_{\odot}$
- ▶ Distance: 7.13 kpc
- ▶ Merger in 300 Myr

Merger Rates of Binary Systems

Author(s)	Information	Type	Merger Rate
Phinney (1991)	pulsar lifetimes,	cons.	5×10^{-8}
	distributions	bguess	7×10^{-6}
Van den Heuval & Lorimar (1996)	pulsar detectability,	cons.	3×10^{-7}
	distribution	bguess	8×10^{-6}
Bailes (1996)	galactic pulsar	lbound	10^{-7}
	birth rates	ubound	10^{-5}
Potegies Zwart & Yungelson (1998)	“scenario machine” w/ supernova kicks		$0.2 - 3$ $\times 10^{-5}$
Bethe & Brown (1998)	common envelope hypercritical accretion	ubound	10^{-5}

Rates in $\text{yr}^{-1} \text{Mpc}^{-3}$

$1 \text{ pc} = 3 \times 10^{18} \text{ cm.}$

Discovery of Double-Pulsar System

Pulsar	PSR J0737-3039A	PSR J0737-3039B
Pulse Period P (ms)	22.69937855615(6)	2773.4607474(4)
Period derivative \dot{P}	$1.74(5) \times 10^{-18}$	$0.88(13) \times 10^{-15}$
Orbital period P_b (day)	0.102251563(1)	—
Eccentricity e	0.087779(5)	—
Characteristic age (My)	210	50
Magnetic field B_s	6.3×10^9	1.6×10^{12}
Spin-down luminosity \dot{E} (erg/s)	5.8×10^{33}	1.6×10^{30}
Distance (kpc)	~ 0.6	—
Stellar mass	1.337(5)	1.250(5)

Merger expected in 85 Myr, a factor 3.5 shorter than PSR 1913+16

A.G Lyne et al., *Science*, **303**, 1153 (2004)

Kalogera et al. (2004): Revisions w/ PSR J037-3039 imply 1 event per 1.5 yr for initial LIGO (for advanced LIGO, 20-1000 events per yr).

PSR J0737 3039 and LIGO

► Merger rate $R \propto N/\tau$

► Binary pulsar lifetime: $\tau = \tau_{BIRTH} + \tau_{COAL}$.

$$\frac{\tau_{1913}}{\tau_{0737}} = \frac{365 \text{ Myr}}{185 \text{ Myr}} \approx 2$$

► scaling factor $N \propto L_{400}^{-1}$

$$\frac{N_{0737}}{N_{1913}} = \frac{L_{1913}}{L_{0737}} = \frac{200 \text{ mJy kpc}^2}{30 \text{ mJy kpc}^2} \approx 6$$

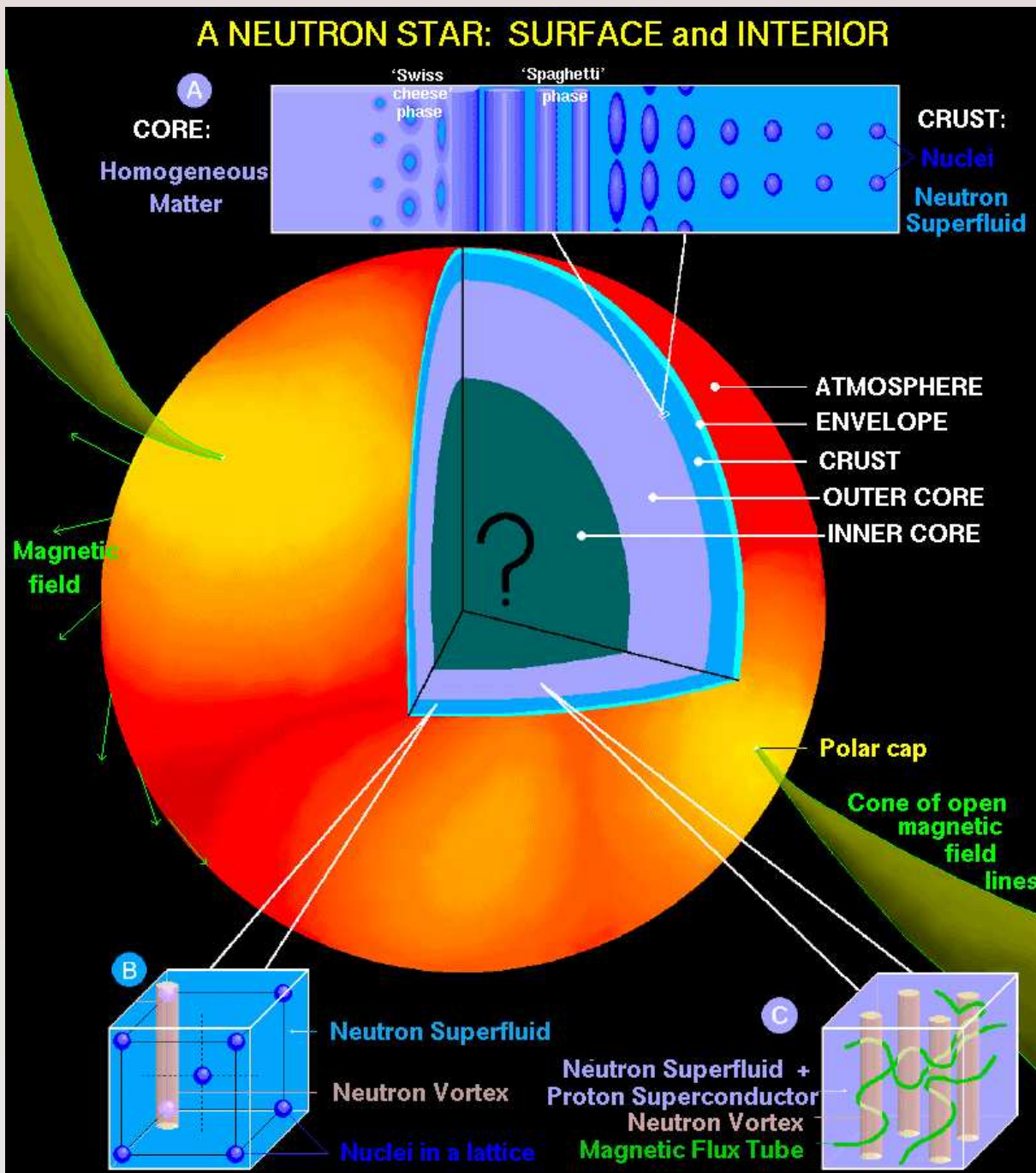
$2 \times 6 = 12 \Rightarrow$ an order of magnitude increase of merger rates!

GW Detectors & Expected Gains

- ▶ Ground-Based Laser Interferometers
 - LIGO, VIRGO, GEO, TAMA, ...
- ▶ The Laser Interferometer Space Antenna (LISA)
- ▶ GW's provide valuable new information “orthogonal” to electromagnetic observations
 - First direct test of GR
 - Precise (\pm a few %) determination of Hubble's constant H_0
 - Calibration of distance measurements
 - Masses of NS, BH (large scale structure formation)
 -

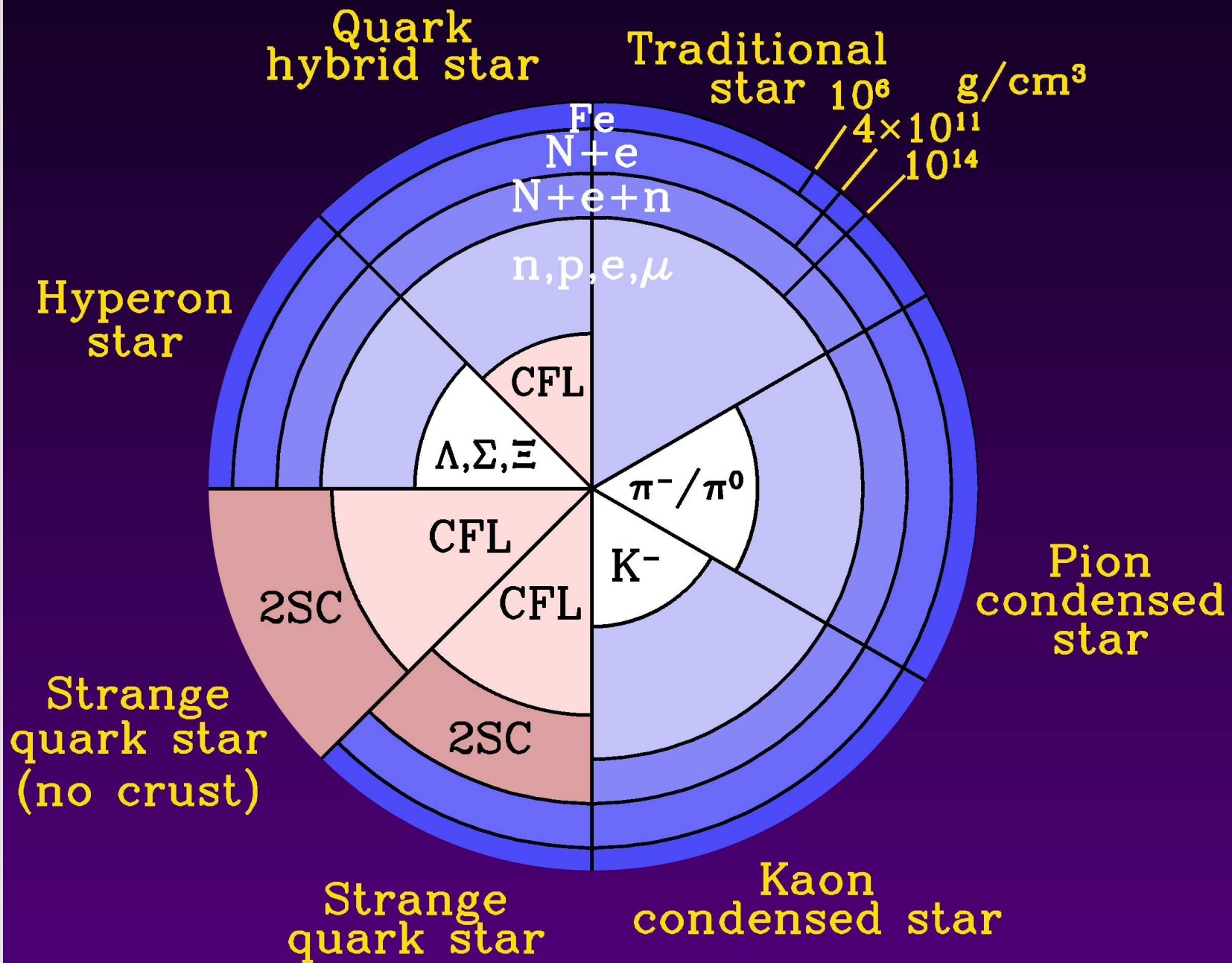
Objectives

- ▶ Explore EOS dependence of GW signals from mergers.
 - Specifically, look at differences between “normal” stars and “self-bound” (e.g., SQM) stars.
 - EOS parameter : $\alpha(M_1) \equiv d \ln(R_1)/d \ln(M_1)$
 - $\alpha_{NS} \leq 0$, while $\alpha_{SQM} \geq 0$ ($\approx 1/3$)
- ▶ Incorporate analysis to include GR (2PN, ...) orbital dynamics.
 - Extend the Roche lobe analysis from Newtonian to 2PN, ...
GR makes stable mass transfer easier.
 - Utilize pseudo-GR potential to account for innermost circular orbit changes as a function of mass ratio. Study effects on results for existence of stable mass transfer.
- ▶ Explore astrophysical consequences of differences in $\alpha(M_1)$ in (1) merger time scales and (2) GW signals.

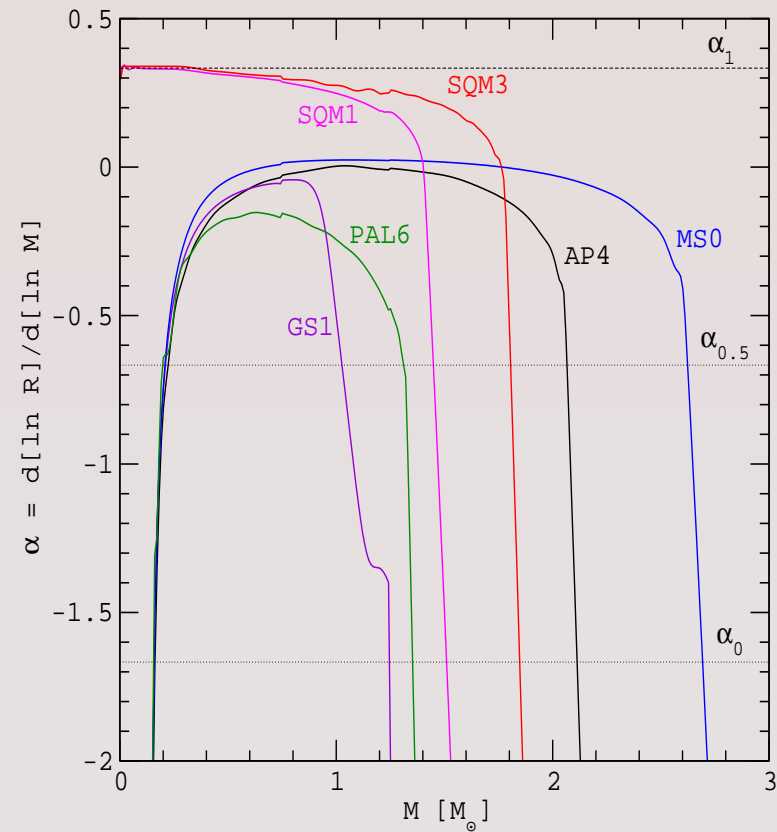
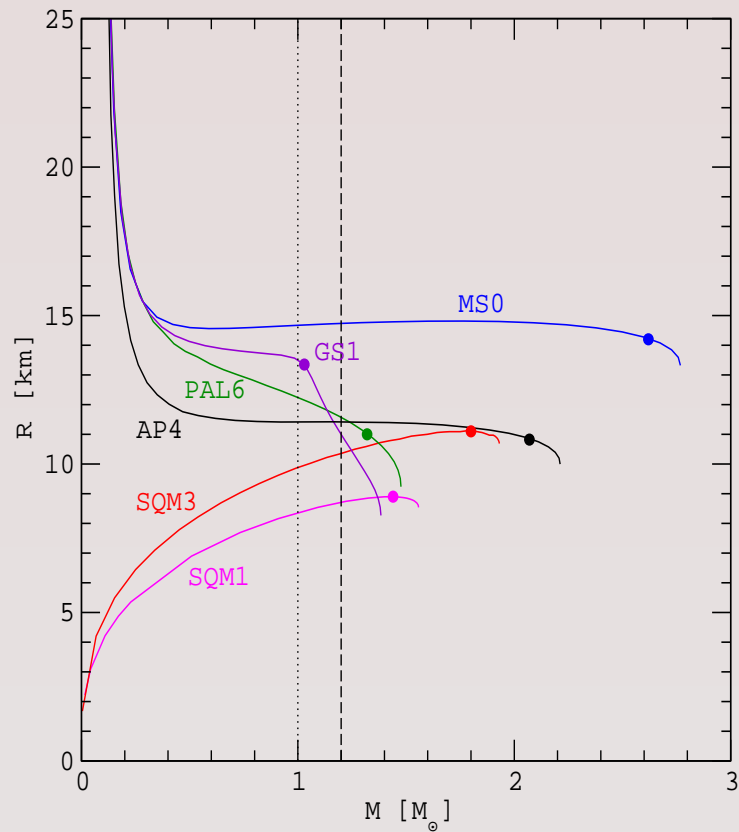


- ▶ $M \sim (1 - 2)M_{\odot}$
 $M_{\odot} \simeq 2 \times 10^{33} \text{ g.}$
- ▶ $R \sim (8 - 16) \text{ km}$
- ▶ $\rho > 10^{15} \text{ g cm}^{-3}$
- ▶ $B_s = 10^9 - 10^{15} \text{ G.}$
- ▶ Tallest mountain:
 $\sim \frac{E_{liq}}{Am_p g_s} \sim 1 \text{ cm}$
- ▶ Atmospheric height:
 $\sim \frac{RT}{\mu g_s} \sim 1 \text{ cm}$

Lattimer & Prakash , Science 304, 536 (2004).



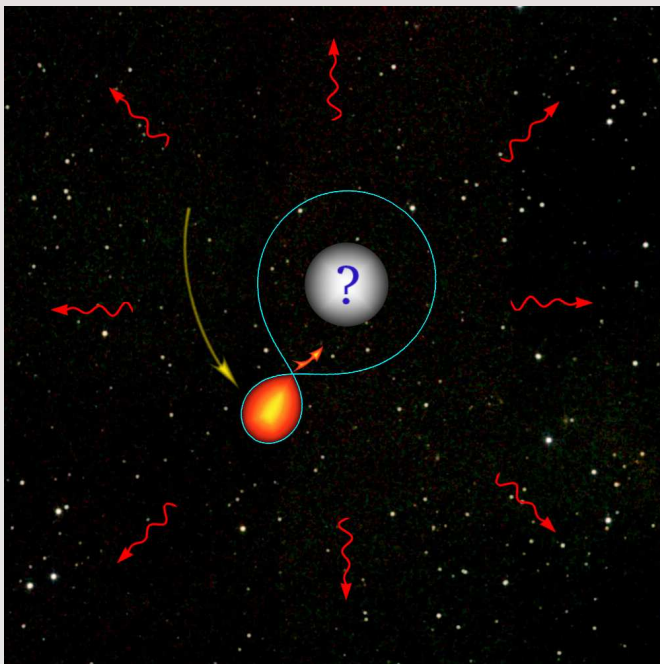
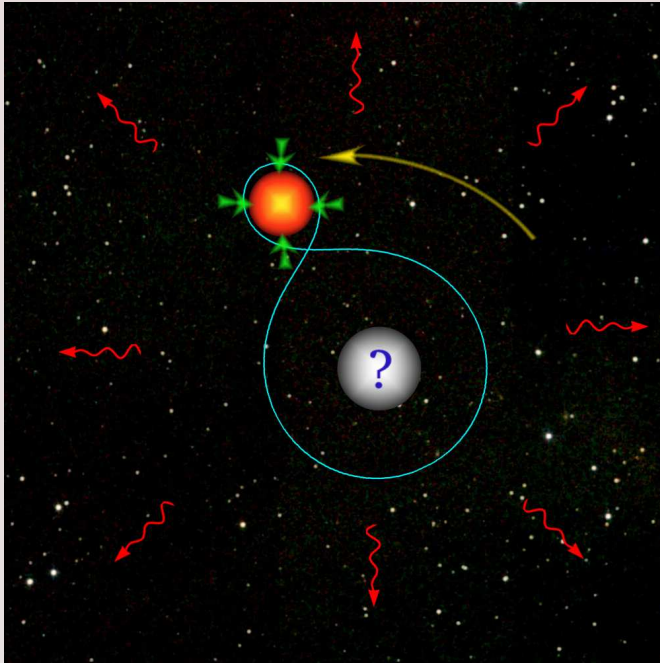
Equation of State: $\alpha(M)$



► $\alpha_{NS} \leq 0$

► $\alpha_{SQM} \geq 0$
 ($\approx 1/3$)

Roche Lobe Overflow



- ▶ Energy Loss

$$L_{GW} = \frac{1}{5} \langle \ddot{\mathcal{I}}_{jk} \ddot{\mathcal{I}}_{jk} \rangle = \frac{32}{5} a^4 \mu^2 \omega^6$$

- ▶ Angular Momentum Loss

$$\left(\dot{J}_{GW} \right)_i = \frac{2}{5} \epsilon_{ijk} \langle \ddot{\mathcal{I}}_{jm} \ddot{\mathcal{I}}_{km} \rangle = \frac{32}{5} a^4 \mu^2 \omega^5$$

- ▶ $a(t)$ and V_{Roche} shrink!

- ▶ $R_1 = r_{Roche}$

⇒ Mass transfer begins!

- ▶ To merge or not to merge?

Pseudo-GR Potentials

- ▶ Paczyński-Wiita (accretion disks)

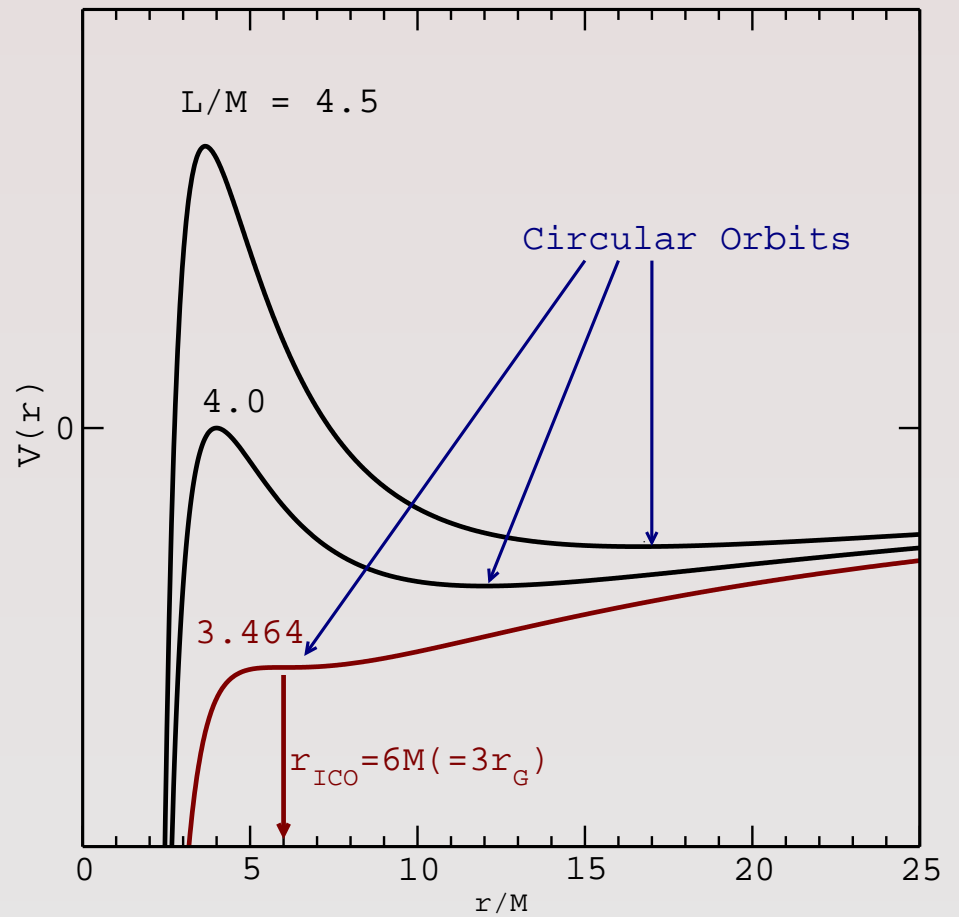
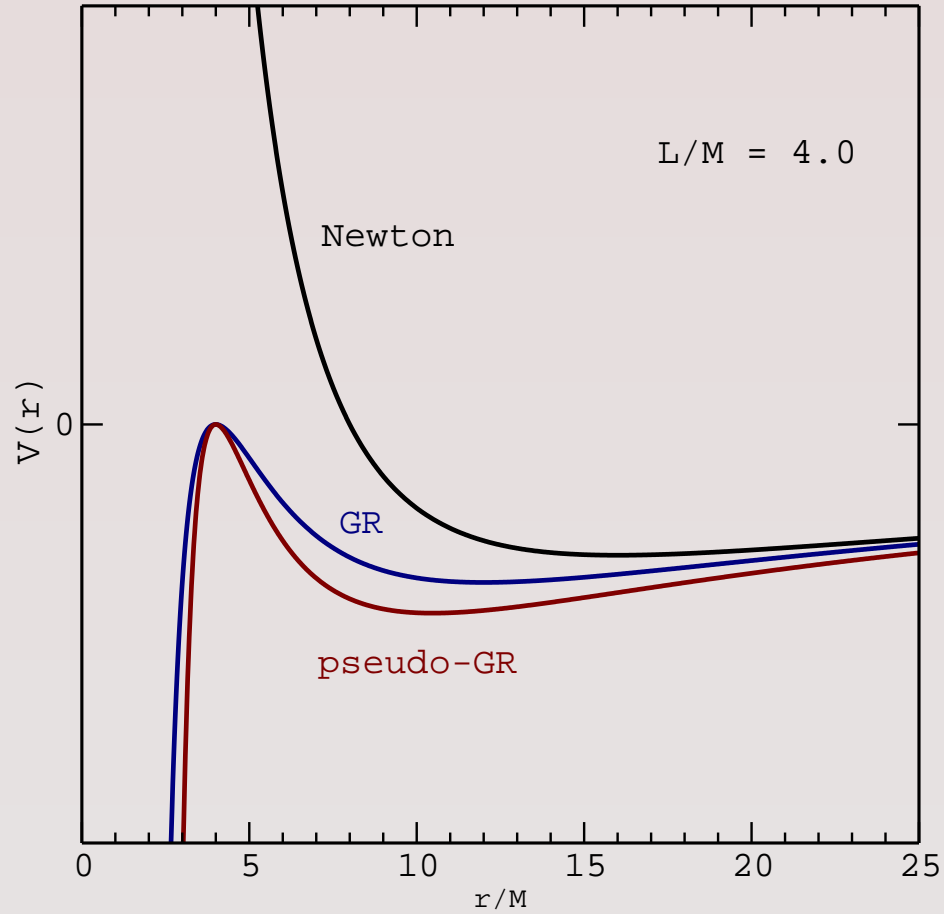
$$\phi_N(r) = -\frac{M}{r} \quad \rightarrow \quad \phi_{PW}(r) = -\frac{M}{r - r_G}$$

- ▶ Innermost Circular Orbit (ICO) at $r_{ICO} = 3r_G$; $r_G = 2M$
- ▶ Post-Newtonian (PN) : $r_{ICO} < 3r_G$ for $q \neq 0$
- ▶ Pseudo-GR or Hybrid Potential :

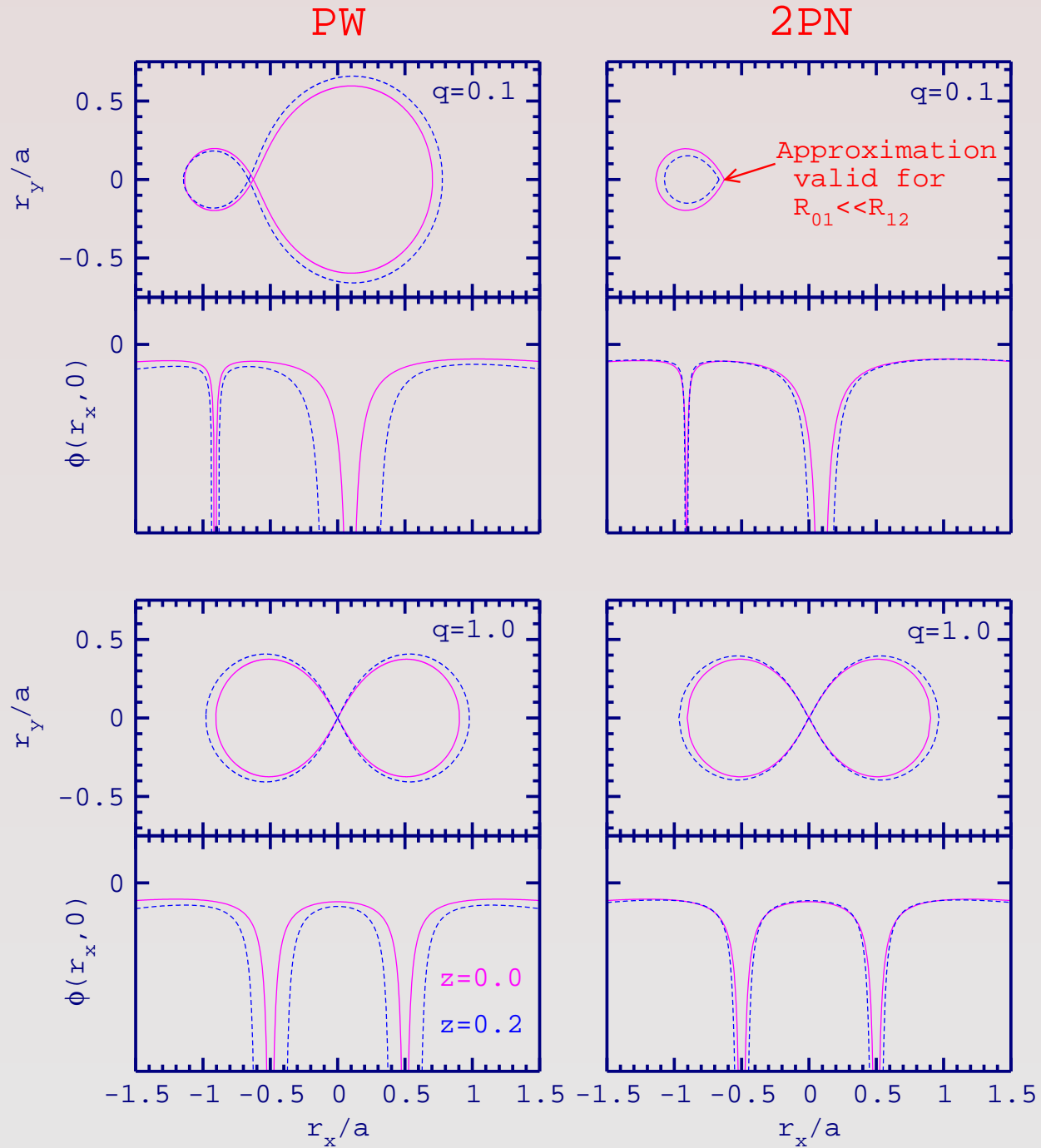
$$\phi_H(r) = -\frac{M}{r - \zeta(q)r_G}; \quad q = M_1/M_2$$

- ▶ $\zeta(q)$ - Mimics 2PN, 3PN Corrections to ICO

Test Particle Effective Potentials and ICO

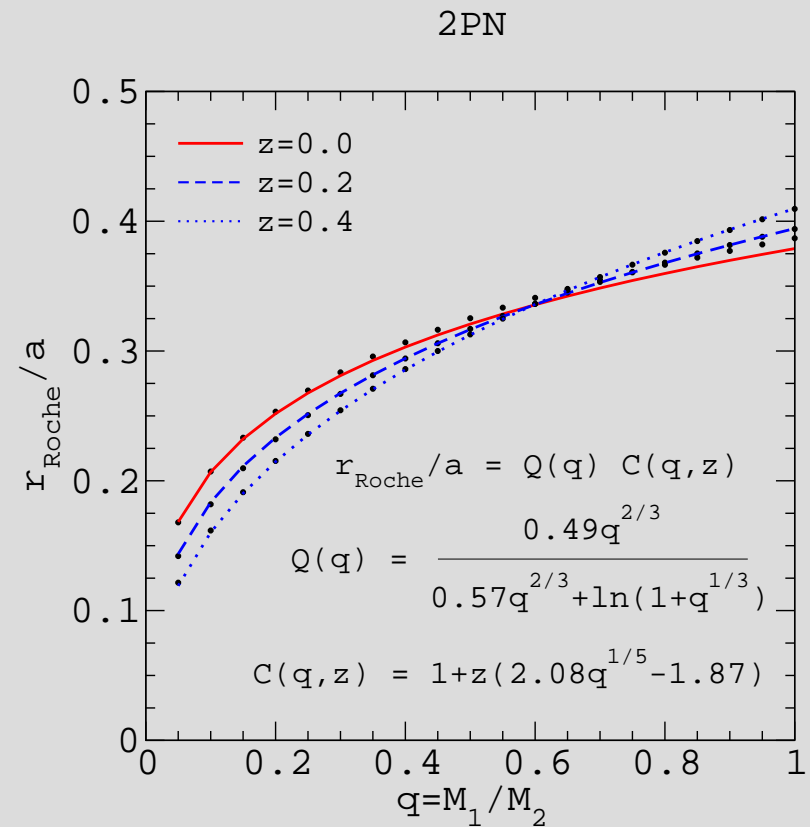
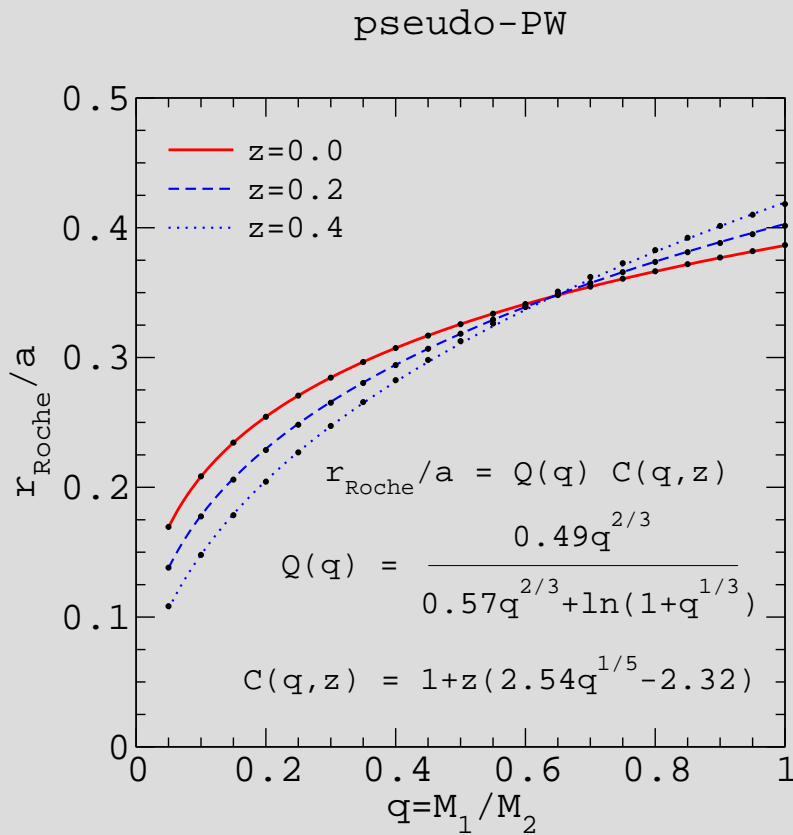


Roche Lobes: PW vs. 2PN



Effective Roche Lobe Radii

Ratković, Prakash,
& Lattimer (2005)



Orbital Evolution

► Angular Momentum Loss :

$$\left[\frac{1-q}{1+q} + \frac{r_G q \zeta'(q)}{a - \zeta(q)r_G} \right] \frac{\dot{q}}{q} + \frac{a - 3\zeta(q)r_G}{2(a - \zeta(q)r_G)} \frac{\dot{a}}{a} = - \frac{\dot{J}_{GW}}{J_{BS}} = - \frac{32}{5} a^2 \mu \omega^4$$

► Roche Lobe :

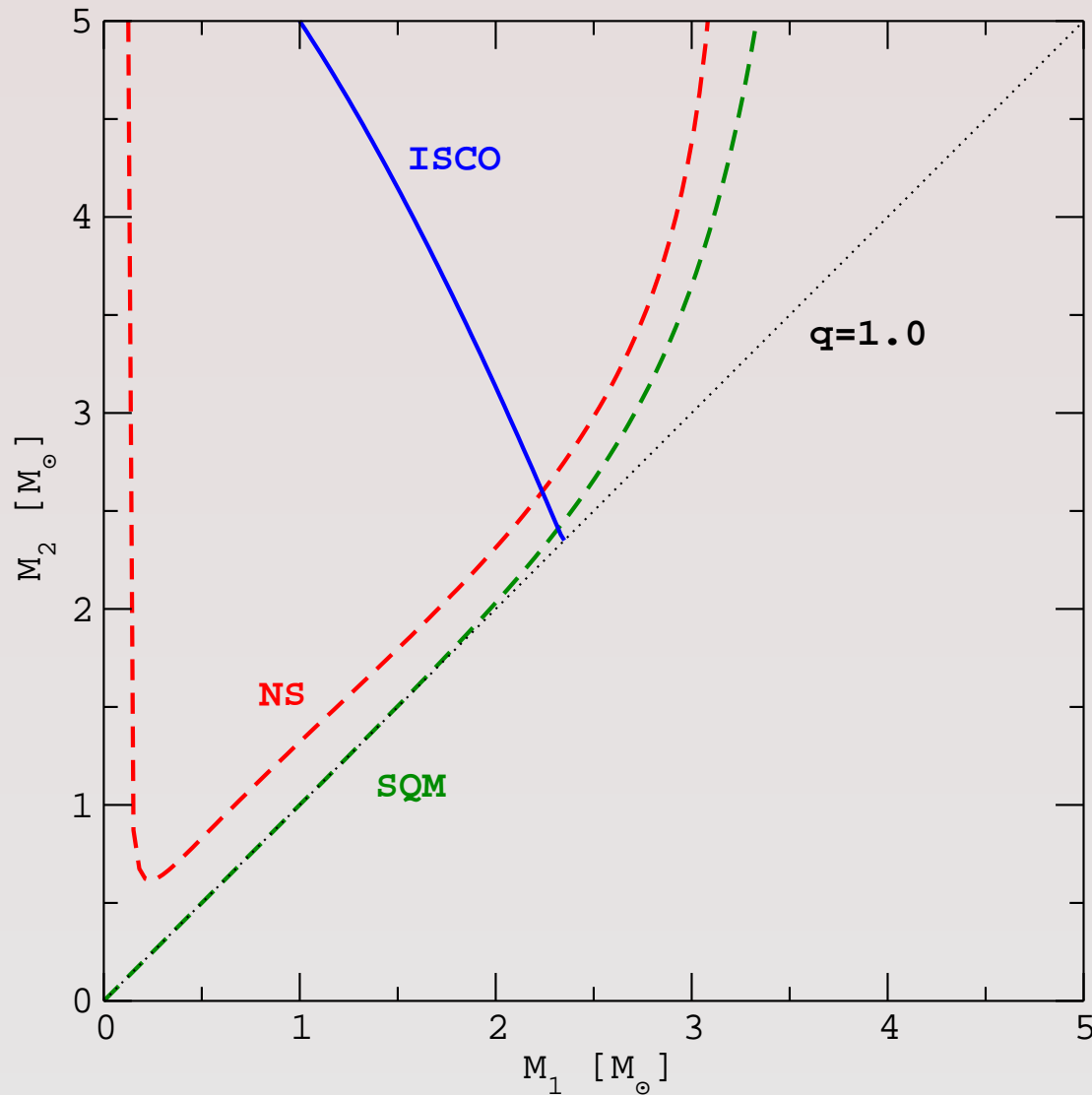
$$\frac{\dot{q}}{q} = \frac{1 - \frac{\partial \ln C(q, z)}{\partial \ln z}}{\frac{\alpha(M_1)}{1+q} - \frac{\partial \ln Q(q)C(q, z)}{\partial \ln q}} \times \frac{\dot{a}}{a}$$

► Connection to the dense matter EOS through

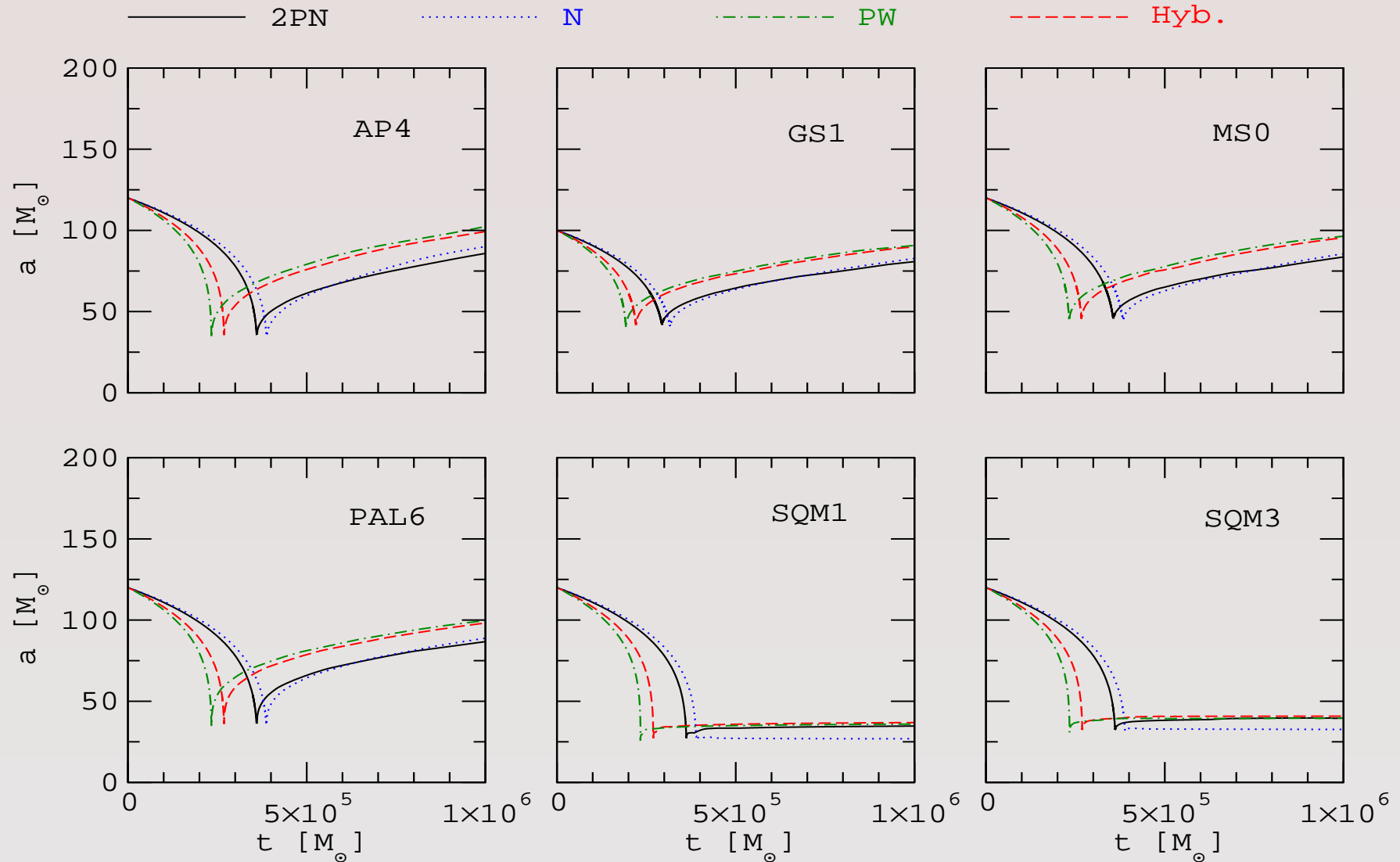
$$\alpha(M_1) \equiv \frac{d \ln(R_1)}{d \ln(M_1)}$$

Regions of stable mass transfer

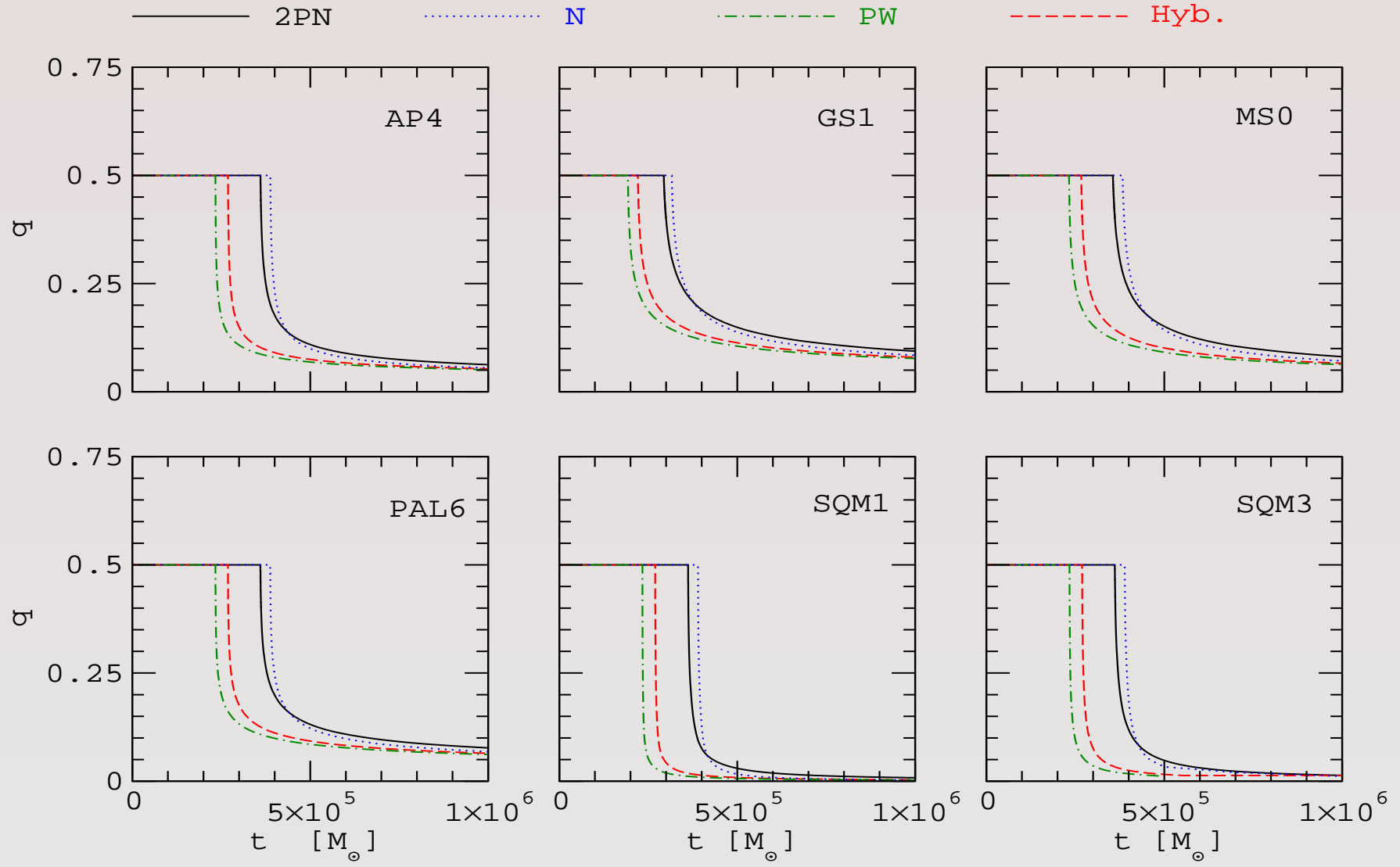
- Dashed curves: Lower mass limit to M_2 for stable mass transfer.
- Solid curve: Upper boundary for transfer beginning outside the ISCO.



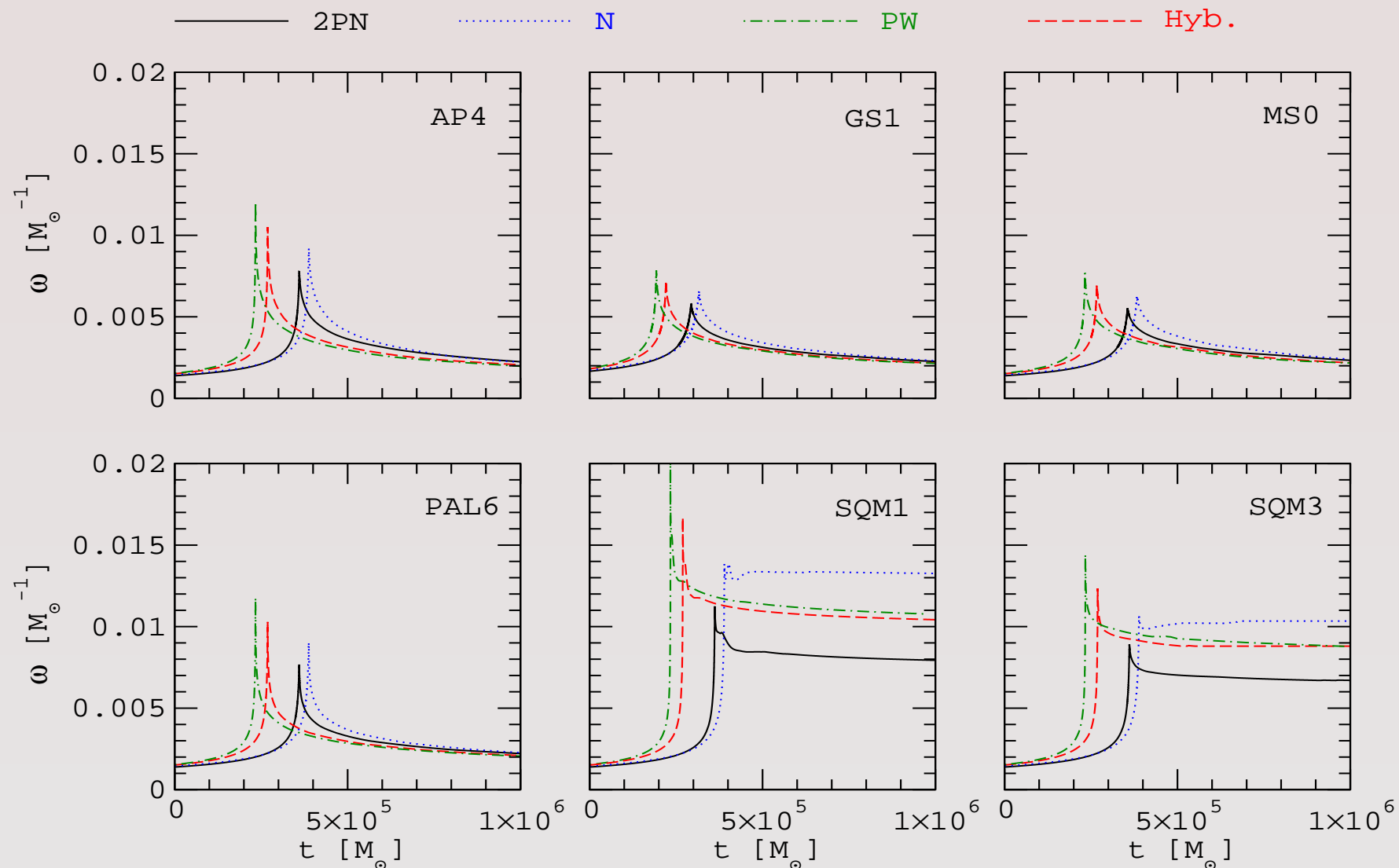
Evolution: Orbit Separation a



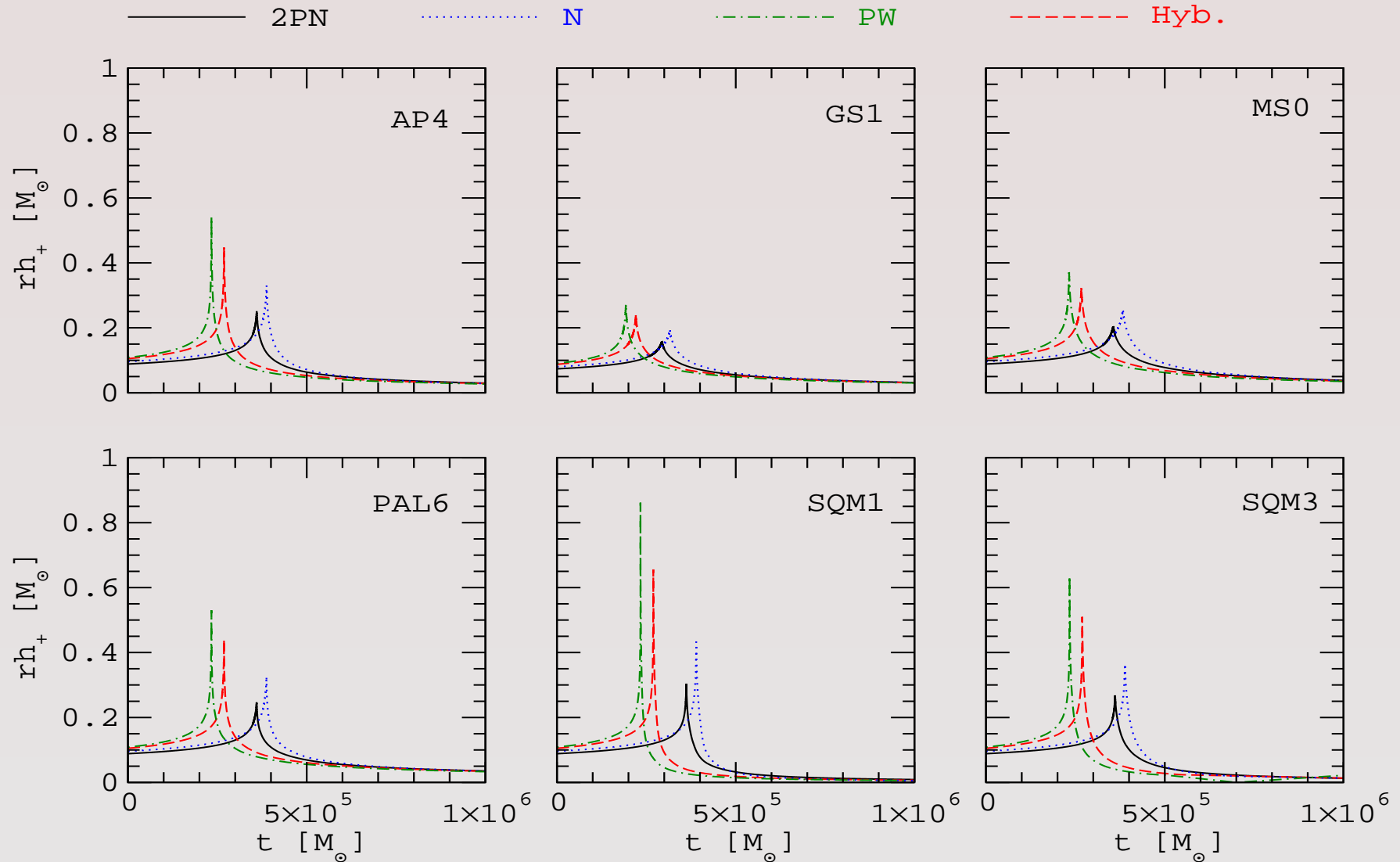
Evolution: Mass Ratio q



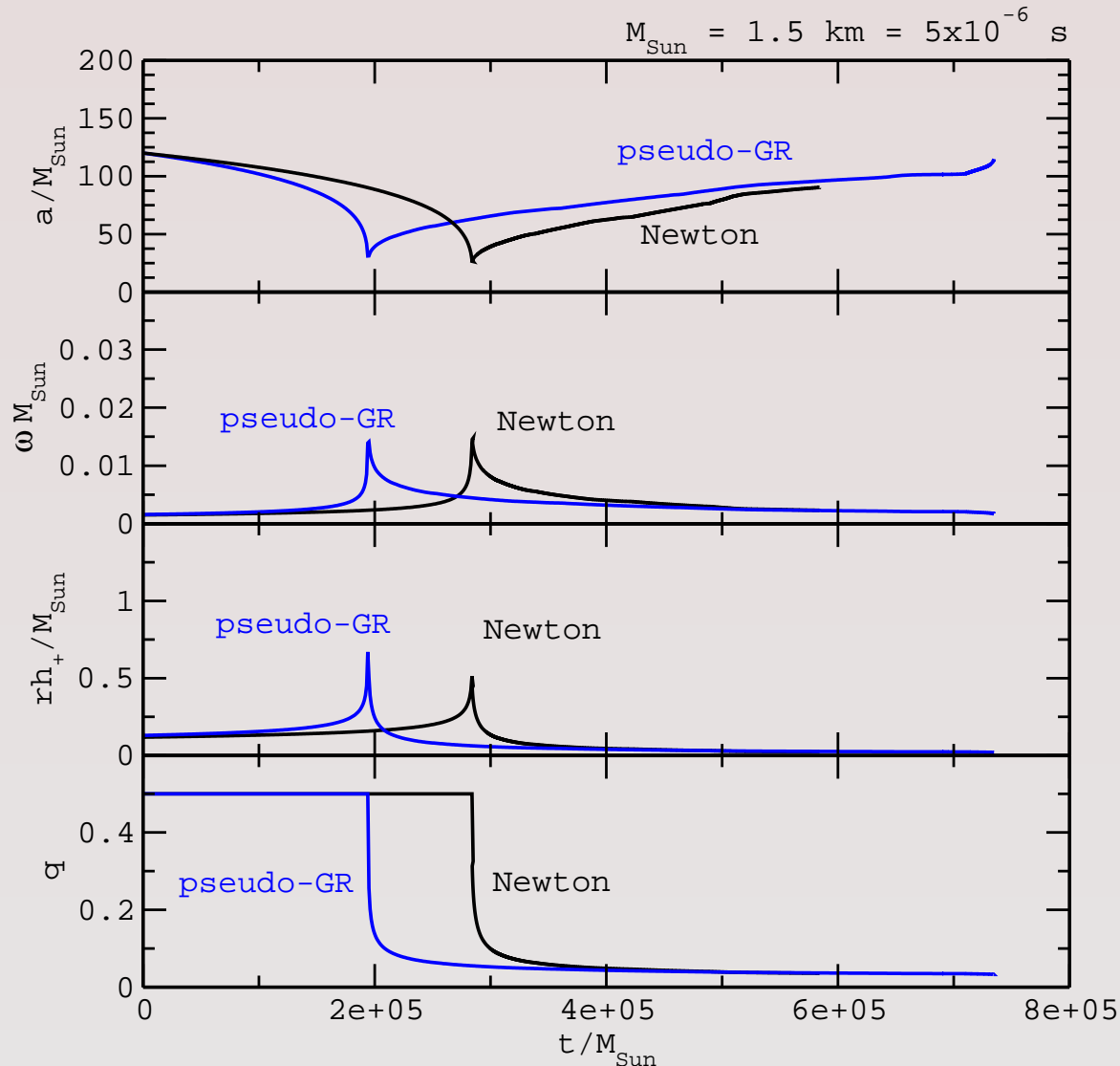
Evolution: Angular Frequency ω



Evolution: Distance \times Gravitational Amplitude rh_+



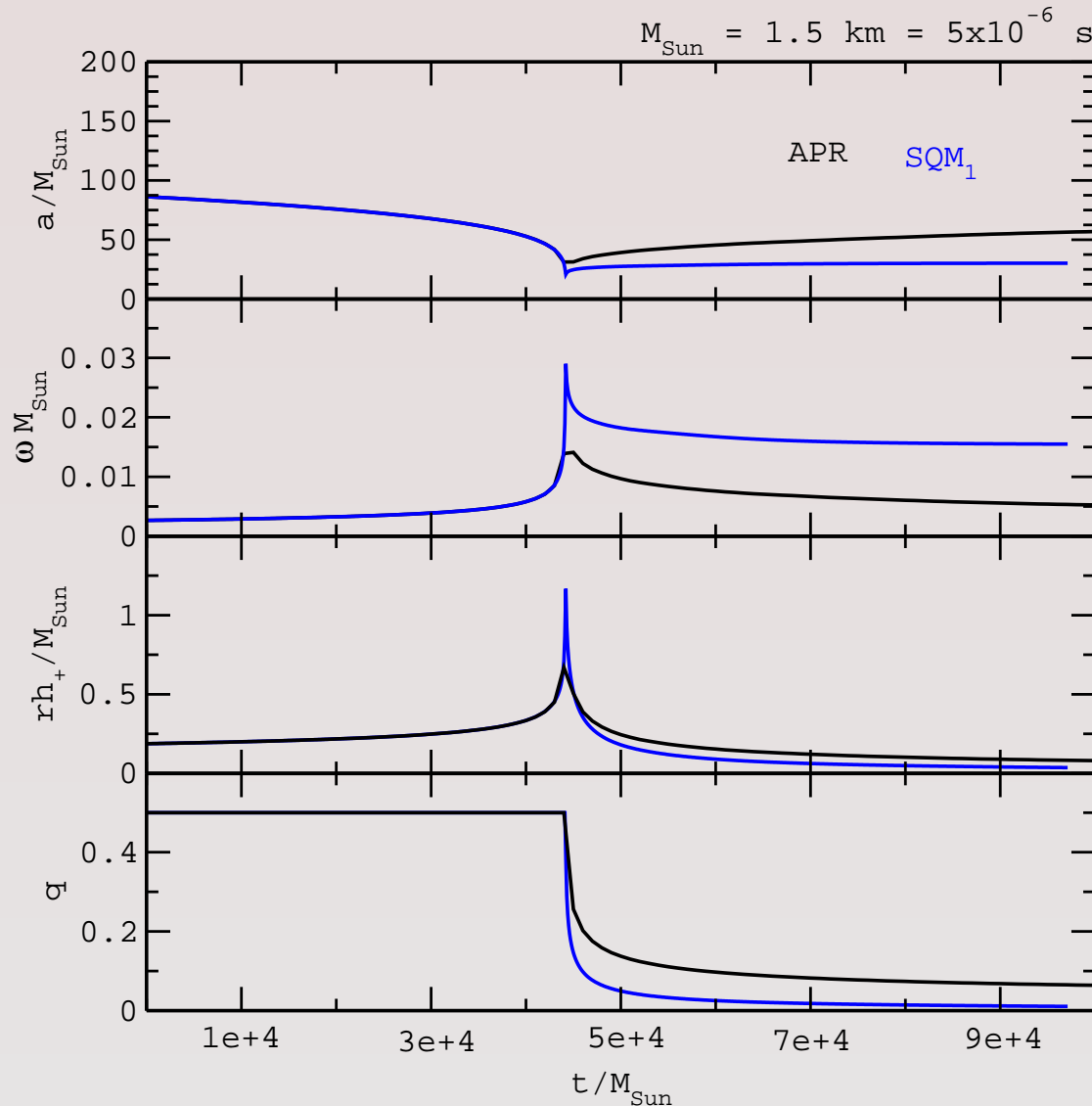
Evolution: Normal Star (*APR*)



$$h_+ = \frac{4}{r} \omega^2 a^2 \mu \cos(2\omega t)$$

- ▶ $M = 4M_{\odot}$, $q_{\text{ini}} = 1/3$
- ▶ GR speeds up evolution
- ▶ $a(t)$ increases after “touchdown”
- ▶ $\omega(t)$ stabilizes at long times
- ▶ Little variation among EOS’s of normal stars.
- ▶ M_1 approaches the NS minimum mass; subsequent plunge (timescale \sim a few minutes) yields a second spike in the GW signal !

Evolution: *SQM* Star



$$h_+ = \frac{4}{r} \omega^2 a^2 \mu \cos(2\omega t)$$

- ▶ $M = 4M_{\odot}$, $q_{\text{ini}} = 1/3$
- ▶ $a(t)$: “hovers” after “touchdown”
- ▶ $\omega(t)$: relaxes to $\gg \omega_{\text{initial}}$
- ▶ $h_{+/\times}(t)$ & $q(t)$: exponential decay unlike for a *NS*
- ▶ $M_{1,\text{final}} \rightarrow M_{\text{nugget}}^{\text{SQM}}$ unlike for a normal star; time to tiny $M_{1,\text{final}}$ is very long!

Major Results

- ▶ Incorporating GR into orbital dynamics leads to an evolution that is faster than the Newtonian evolution.
- ▶ Large differences exist between mergers of “normal” and “self-bound (SQM)” stars.
 - SQM stars penetrate to smaller orbital radii; stable mass transfer is more difficult than for normal stars.
 - For stable mass transfer, $q = M_1/M_2$ and $M = M_1 + M_2$ limits on SQM stars are more restrictive than for normal stars.
 - The SQM case has exponentially decaying signal and mass, while normal star evolution is slower.

Future Tasks

- ▶ Evolution of normal & self-bound star-black hole mergers including the effects of
 - non-conservative mass transfer,
 - tidal synchronization,
 - the presence accretion disk, etc.
- ▶ Calculation of templates of expected GW signals





That's All Folks!