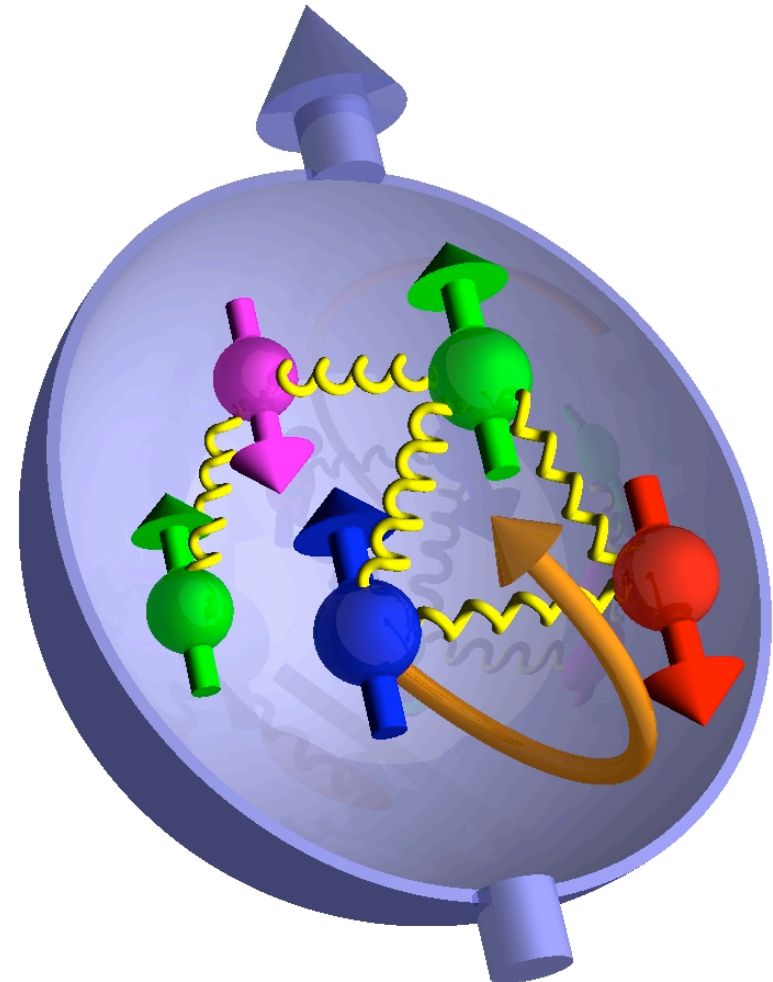


Nucleon Spin & Flavor Structure ... and Fragmentation too

N.C.R. Makins

Univ of Illinois at Urbana-Champaign

- **Quark Models:**
how to think about the proton?
- **Deep-Inelastic Scattering:**
*parton distribution functions
and fragmentation functions*
- **The Spin Puzzle**
and quark polarization
- **Single-Spin Asymmetries:**
*new structures within the proton
and the fragmentation process*



The Wacky World of Quarks

The Quark Model

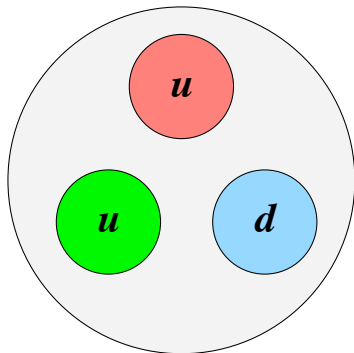
Hadrons are composed of **quarks** with :

- ❶ **flavor**: u,c,t (charge +2/3) d,s,b (charge -1/3) ❷ **color**: R,G,B ❸ **spin**: 1/2

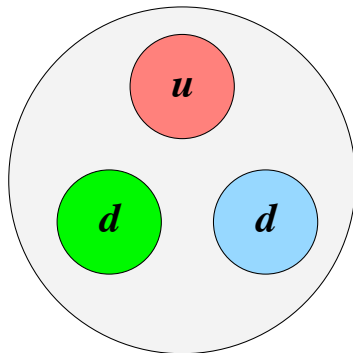
Each hadron observed in nature is **white** ("color singlet")

➤ **Baryons** 3-quark systems, with colors RGB

➤ **Mesons** quark + antiquark with colors CC



proton



neutron

The **spectrum** of observed hadrons is (roughly) explained:

Mesons: Spin 0

π^+ $u\bar{d}$
 π^- $d\bar{u}$
 π^0 $u\bar{u} \oplus d\bar{d}$
 K^+ $u\bar{s}$
 K^- $s\bar{u}$
 K^0 $d\bar{s}$
 \bar{K}^0 $s\bar{d}$
 η $u\bar{u} \oplus d\bar{d} \oplus s\bar{s}$
 η' $u\bar{u} \oplus d\bar{d} \oplus s\bar{s}$

Mesons: Spin 1

ρ^+ $u\bar{d}$
 ρ^- $d\bar{u}$
 ρ^0 $u\bar{u} \oplus d\bar{d}$
 K^{*+} $u\bar{s}$
 K^{*-} $s\bar{u}$
 K^{*0} $d\bar{s}$
 \bar{K}^{*0} $s\bar{d}$
 ϕ $s\bar{s}$
 ω $u\bar{u} \oplus d\bar{d} \oplus s\bar{s}$

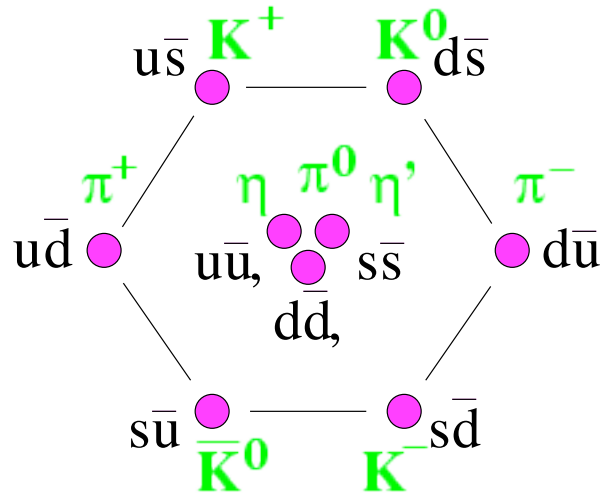
Baryons: Spin 1/2

p uud
 n udd
 Σ^+ uus
 Σ^0 uds
 Σ^- dds
 Λ uds
 Ξ^0 uss
 Ξ^- dss

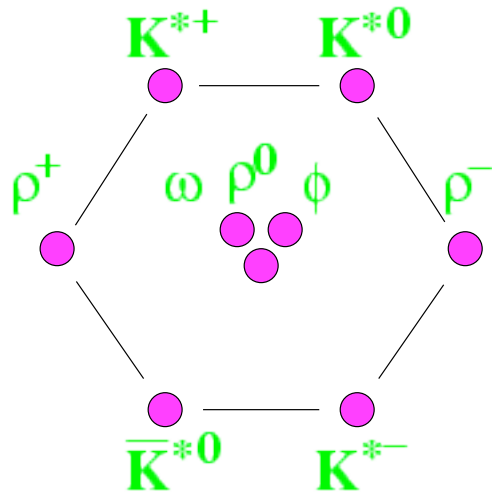
Hadronic Multiplets

• MESONS = $q\bar{q}$

SPIN 0 ($\uparrow\downarrow$)

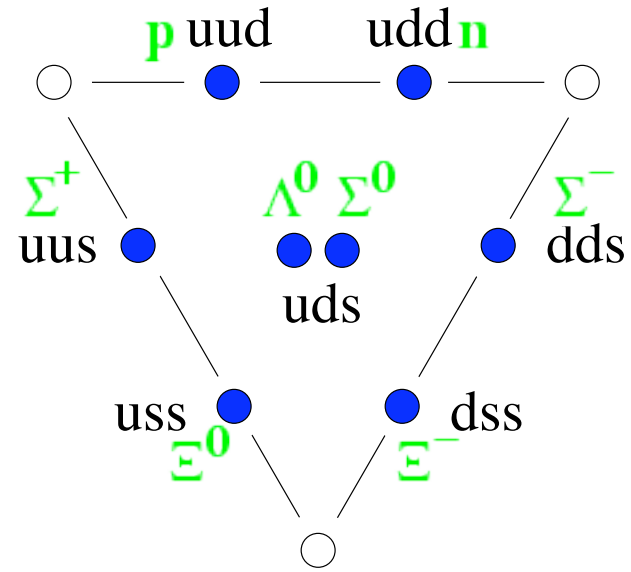


SPIN 1 ($\uparrow\uparrow$)

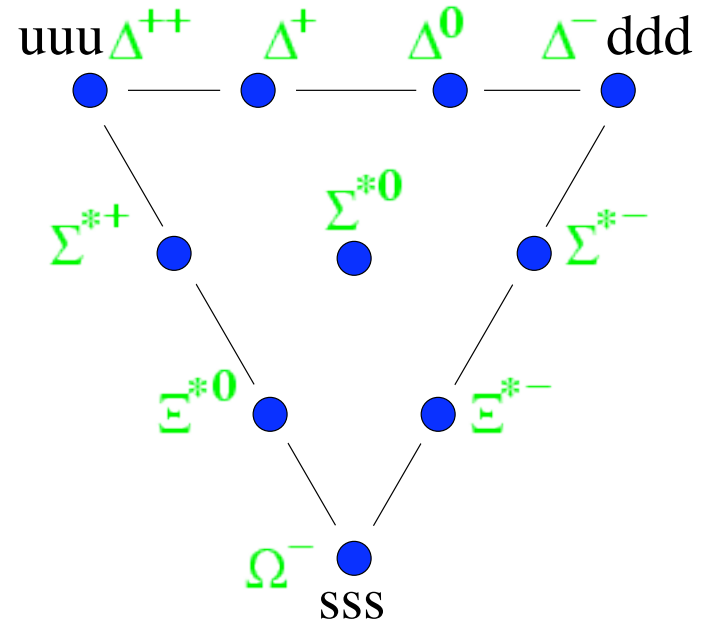


• BARYONS = qqq or $\bar{q}\bar{q}\bar{q}$

SPIN 1/2 ($\uparrow\downarrow\uparrow$)



SPIN 3/2 ($\uparrow\uparrow\uparrow$)



Murray Gell-Mann, 1964:

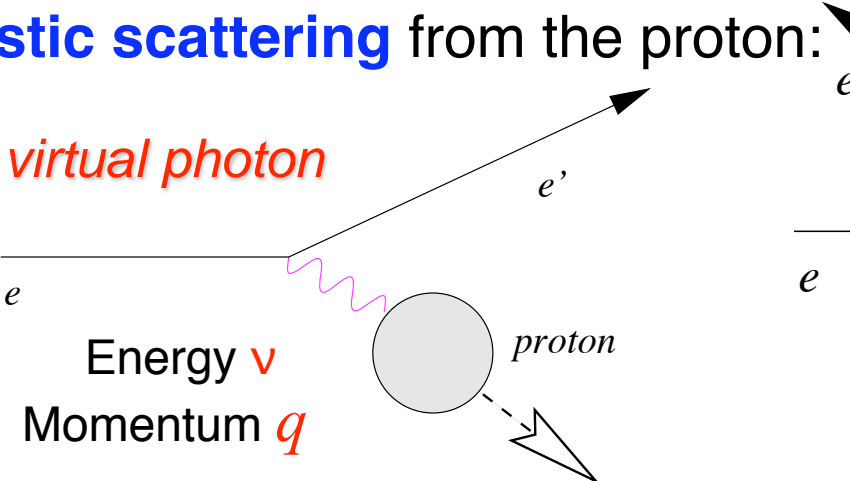
“A search for stable quarks ... at the highest energy accelerators would help to reassure us of the non-existence of real quarks.”

Electron Scattering and Scaling

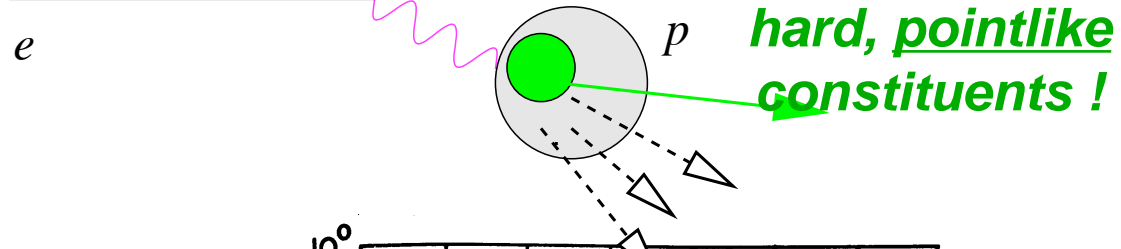
Elastic scattering from the proton:

Deep-Inelastic scattering (DIS):

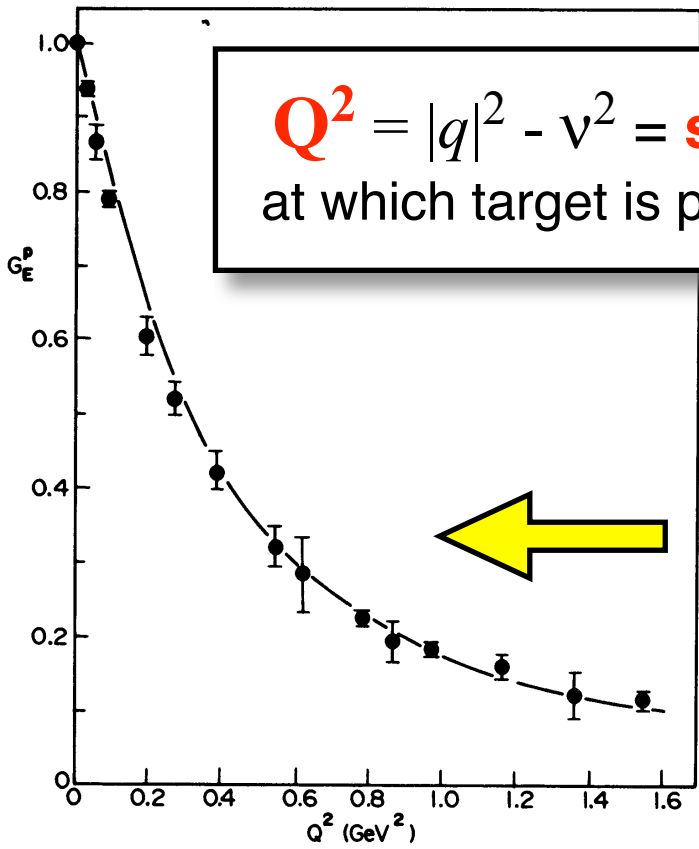
virtual photon



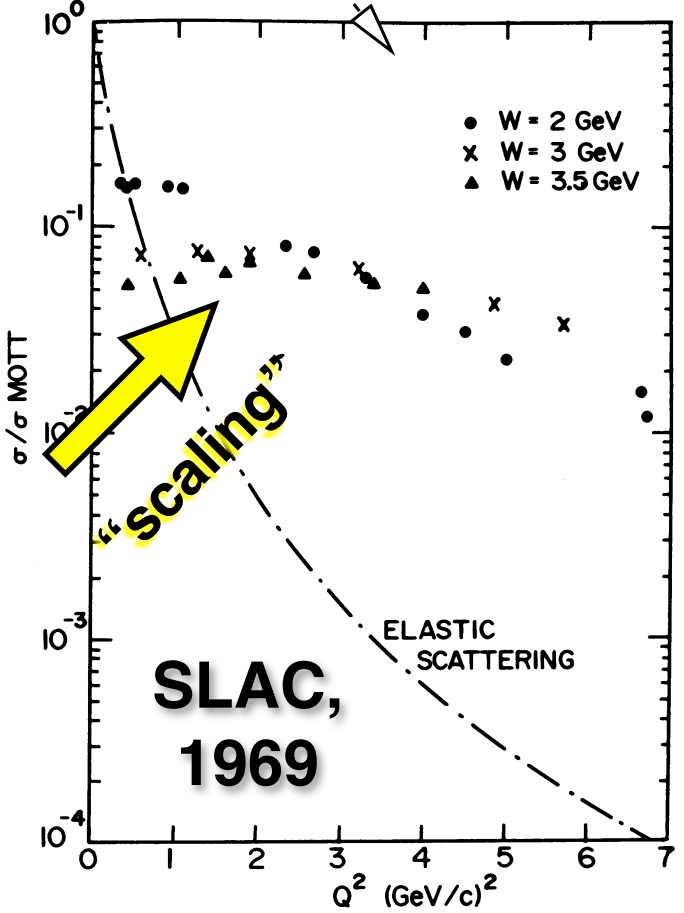
at **high energies** you see ...



$Q^2 = |q|^2 - \nu^2 = \text{scale}$
at which target is probed

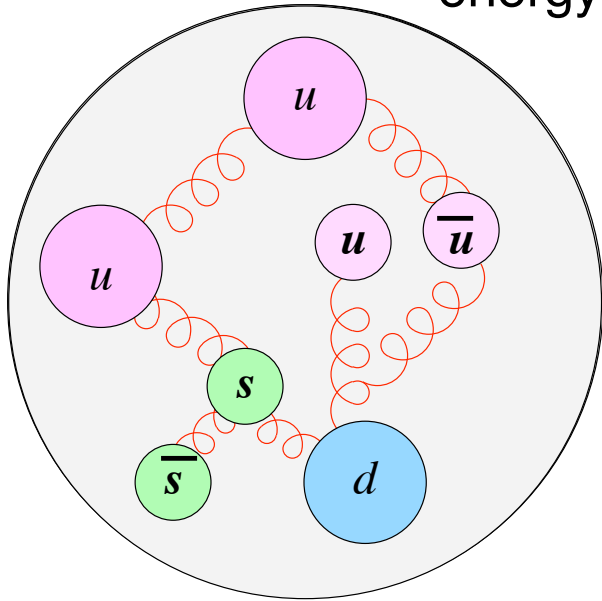


$$\frac{\sigma(Q^2)}{\sigma_{\text{point}}(Q^2)}$$



Parton Distribution Functions

Let's look *inside* the proton: **Deep-Inelastic Scattering** (DIS) with high energy beams \Rightarrow a rich substructure is revealed!



3 **constituent quarks**
of mass ≈ 350 MeV



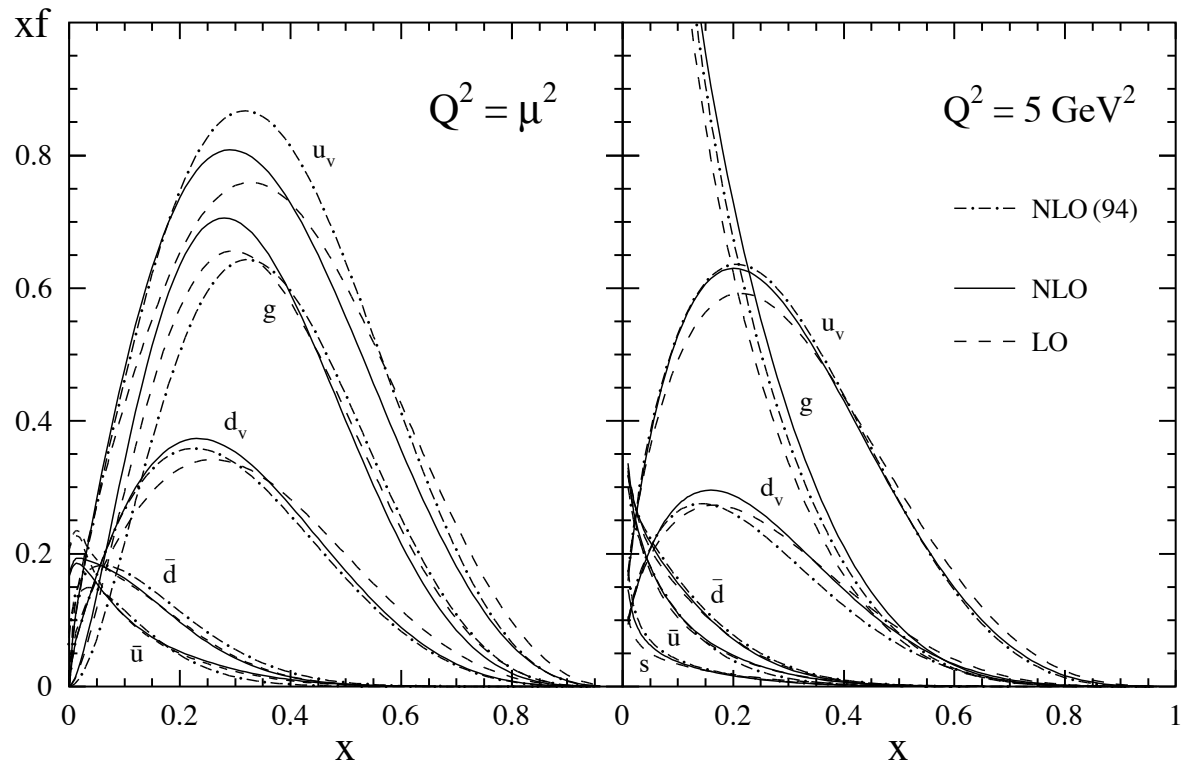
∞ many **current quarks**
with bare masses ≈ 5 MeV

sea quarks : virtual quark-antiquark pairs that fluctuate in and out of the vacuum!

gluons : carriers of the strong force

x fraction of proton momentum carried by struck quark

$q(x)$ parton distribution funcⁿ
(number density for quark flavor q)



Quantum Chromodynamics

The Theory of the Strong Interaction

$$\mathcal{L}_{\text{QCD}} = -\bar{\Psi} \left\{ \gamma_{\mu} \left[\partial_{\mu} - \frac{i}{2} g \lambda^a A_{\mu}^a(x) \right] + M \right\} \Psi - \frac{1}{4} \mathcal{F}_{\mu\nu}^a \mathcal{F}_{\mu\nu}^a$$

The End.

Bound States in QED and QCD

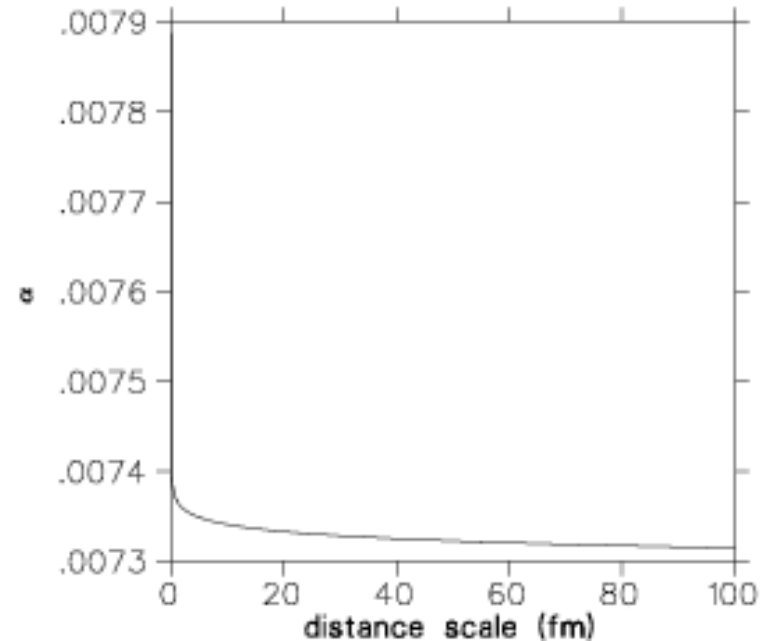
QED

Coupling $\alpha = 1/137$ is weak at relevant scales

✓ **Perturbation theory** works very well

✓ **Non-relativistic** quantum mechanics ok

e.g. Hydrogen: binding $E = 13.6 \text{ eV} \ll M_{\text{elec}} = 511 \text{ keV}$



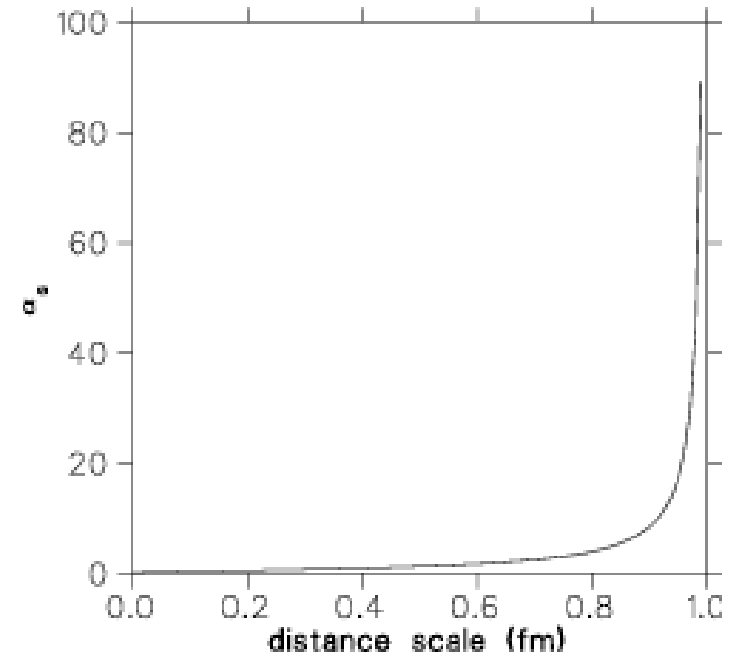
QCD

Coupling α_s ***blows up*** at relevant scales !

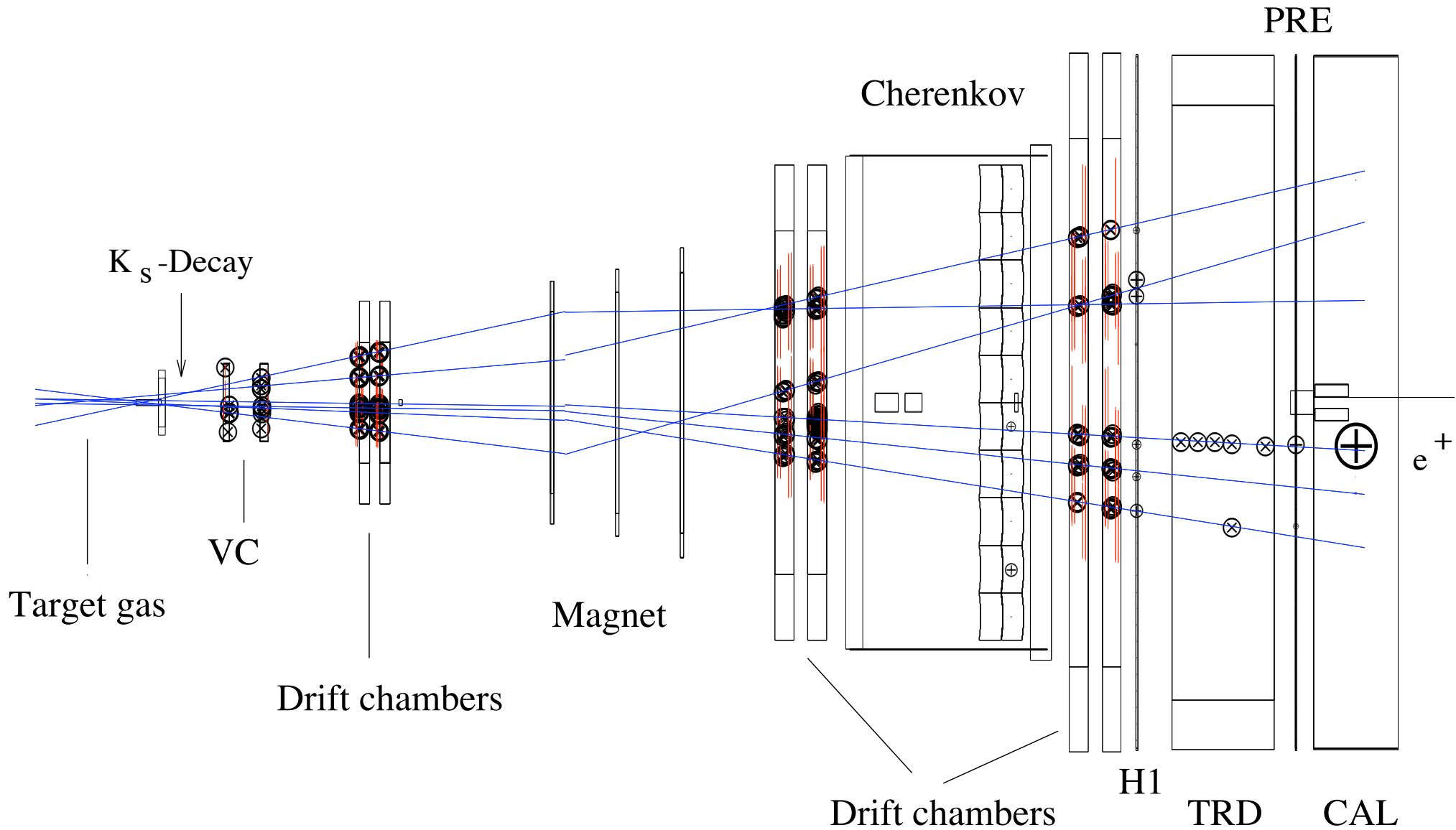
✗ **Perturbation theory impossible**

✗ Bound systems inherently **relativistic**

e.g. Proton: Mass = 938 MeV \gg
bare $m_{\text{quark}} = 5 \text{ MeV}$!

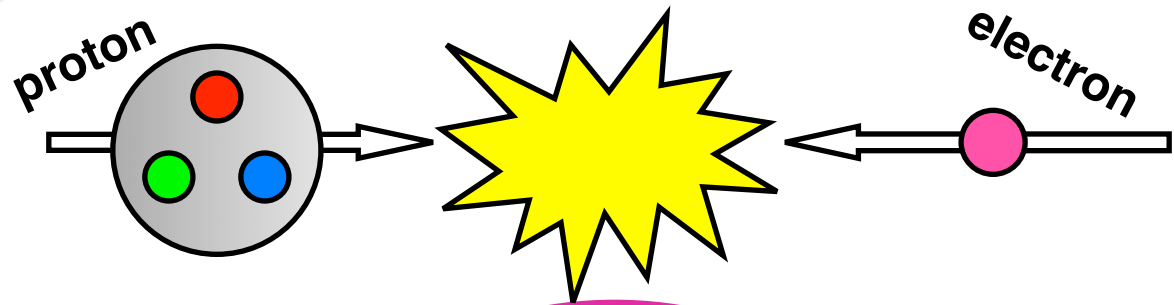


And here's something else we can't calculate ...



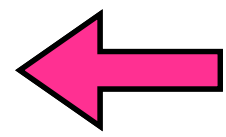
What Happens in a High Energy Collision

Lund String Model

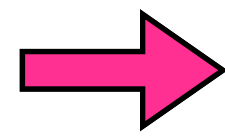


HADRONS
are formed, in
"JETS"

n π^+ Δ^+ K
 ω K^- $+$ $+$
 ϕ η



K^0 π^0 ρ
 Λ^0 ρ^0 Σ^-
 π^-



$$z \equiv E_h / E_q$$

Confinement at Work !

Creation of **hadrons** from **struck quark**: the "fragmentation process"

Fragmentation Functions

$$D_q^h(z)$$

describe **number density** of **hadrons** of type h and energy-fraction z produced from a **struck quark** of flavor q

$$z \equiv E_h/E_q$$

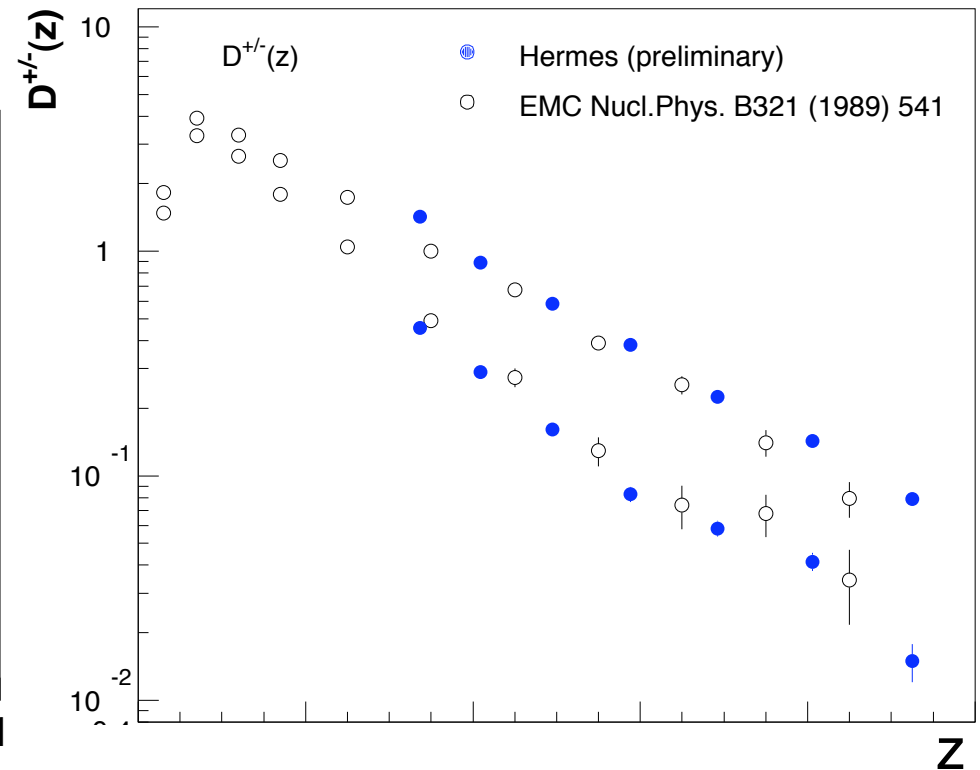
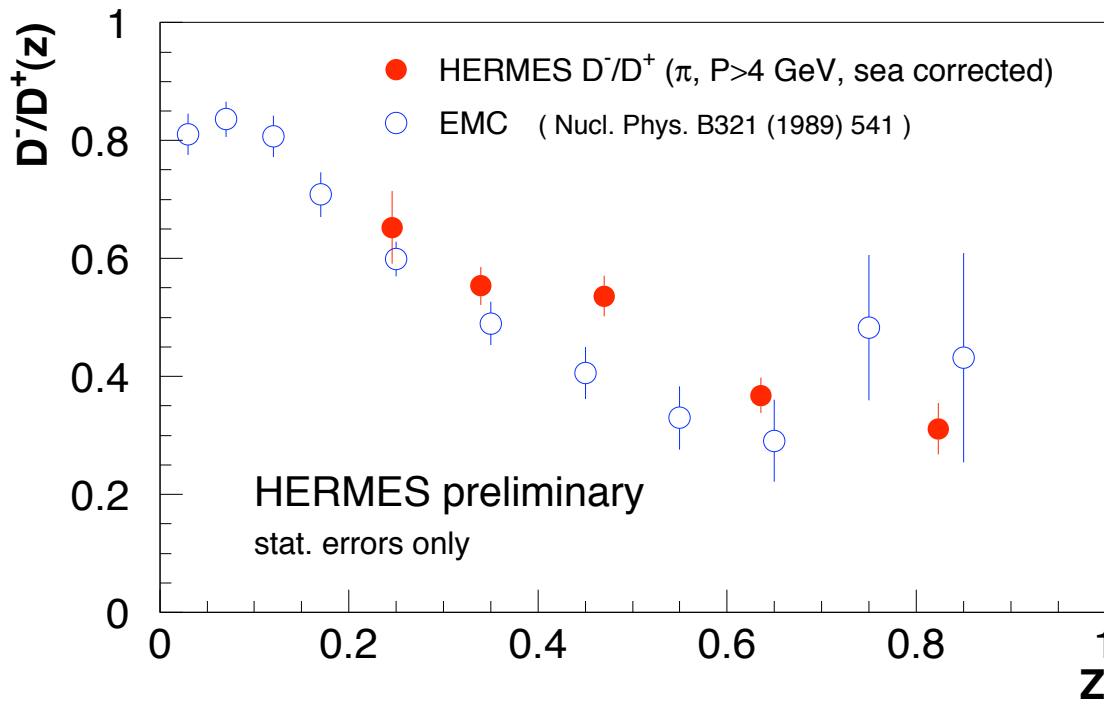
Symmetries: favored / disfavored FF's for pions:

$$D_{\text{fav}}(\text{or } D^+) \equiv D_u^{\pi^+} = D_d^{\pi^-} = D_{\bar{d}}^{\pi^+} = D_{\bar{u}}^{\pi^-}$$

$$D_{\text{dis}}(\text{or } D^-) \equiv D_u^{\pi^-} = D_d^{\pi^+} = D_{\bar{d}}^{\pi^-} = D_{\bar{u}}^{\pi^+}$$

$$D_s \equiv D_s^{\pi^+} = D_s^{\pi^-} = D_{\bar{s}}^{\pi^+} = D_{\bar{s}}^{\pi^-}$$

(based on *charge-conjugation* and *isospin symmetry*)

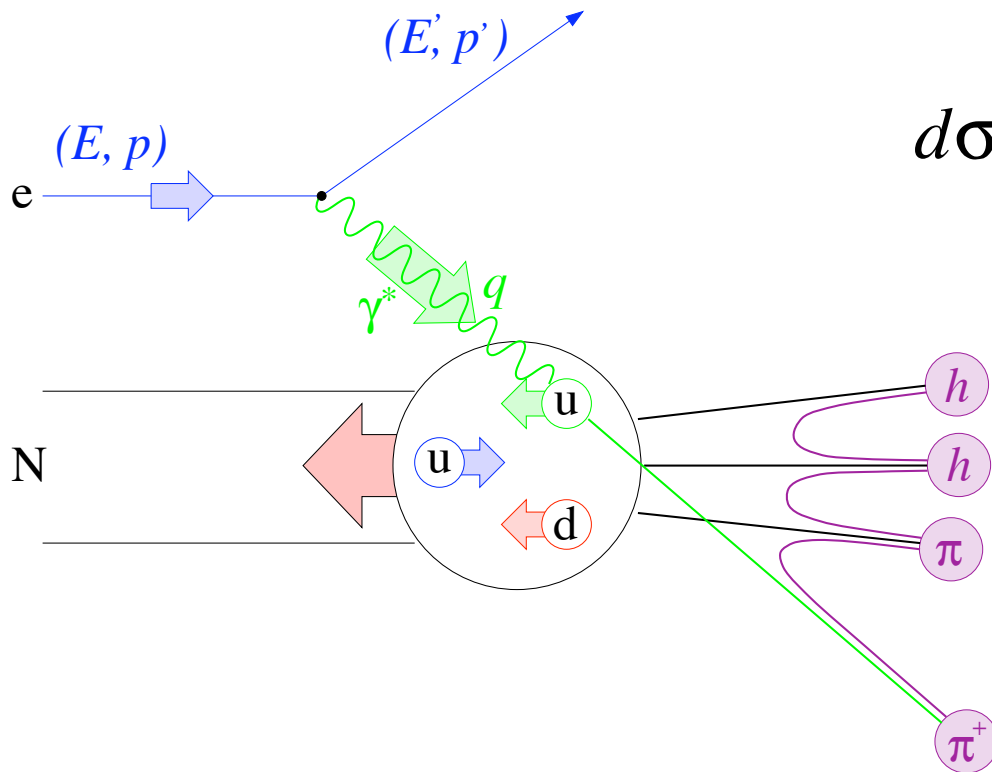


Semi-Inclusive Deep-Inelastic Scattering (SIDIS)

In SIDIS, a **hadron h** is detected in **coincidence** with the scattered lepton:

Factorization of the cross-section:

$$d\sigma^h \sim \sum_q e_q^2 \underbrace{q(x)}_{\text{green}} \cdot \underbrace{\hat{\sigma}}_{\text{blue}} \cdot \underbrace{D^{q \rightarrow h}(z)}_{\text{pink}}$$



● The perturbative part

Cross-section for elementary photon-quark **subprocess**

Large energies \rightarrow asymptotic freedom
 \rightarrow can calculate!

● The Distribution Function

momentum **distribution of quarks q**
 within their proton bound state

\rightarrow **lattice QCD** progressing steadily

● The Fragmentation Function

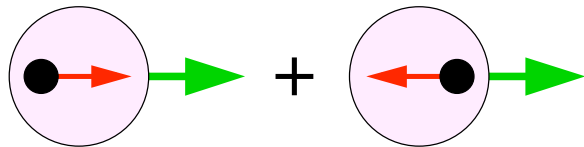
momentum **distribution of hadrons h**
 formed from quark q

\rightarrow not even lattice can help ...

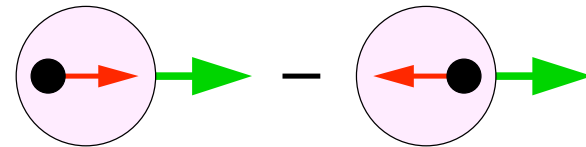
The Spin Puzzle

A particular puzzle: Where does the proton spin come from?

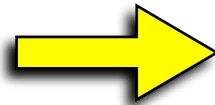
$$q(x) = q^\uparrow(x) + q^\downarrow(x)$$



$$\Delta q(x) = q^\uparrow(x) - q^\downarrow(x)$$



only three possibilities



$$\frac{1}{2} = \frac{1}{2}\Delta\Sigma + \Delta G + L_q + L_g$$

1 Quark polarization

$$\Delta\Sigma \equiv \int dx (\Delta u(x) + \Delta d(x) + \Delta s(x) + \Delta \bar{u}(x) + \Delta \bar{d}(x) + \Delta \bar{s}(x)) \approx 20\% \text{ only}$$

2 Gluon polarization

$$\Delta G \equiv \int dx \Delta g(x) \quad ?$$

In friendly, **non-relativistic** bound states like atoms & nuclei (& constituent quark model), particles are in **eigenstates of L**

3 Orbital angular momentum

$$L_z \equiv L_q + L_g \quad ?$$

?

Not so for bound, **relativistic Dirac particles ...**
Noble " l " is **not a good quantum number**

Flavor Structure of the Proton

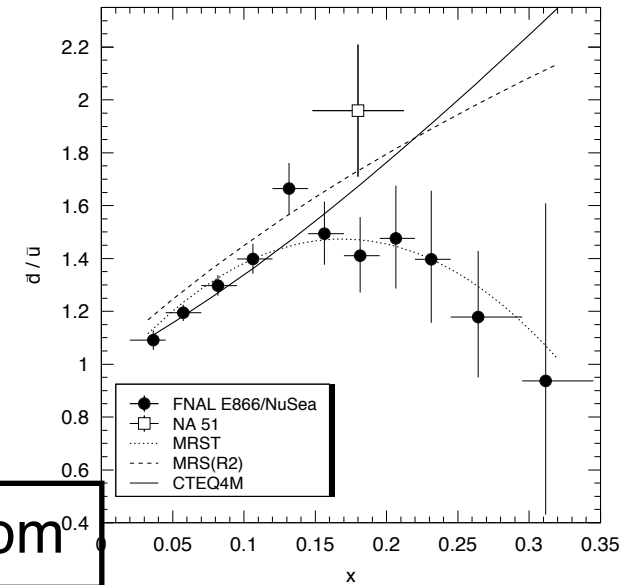
E866: $\bar{d}/\bar{u} > 1$

Constituent Quark Model

Pure valence description: proton = $2u + d$

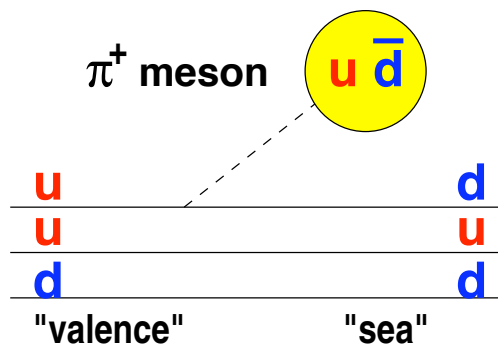
Perturbative Sea Sea quark pairs from $g \rightarrow q\bar{q}$ should be flavor symmetric:

$$\bar{u} = \bar{d}$$



Non-perturbative models: alternate deg's of freedom

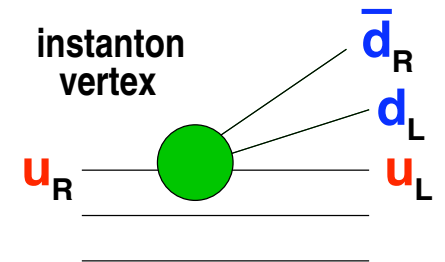
Meson Cloud Models



Quark sea from cloud of 0^- mesons: $\bar{d} > \bar{u}$

Chiral-Quark Soliton Model

- quark degrees of freedom in a pion mean-field
- nucleon = chiral soliton
- one parameter: dynamically-generated quark mass
- expand in $1/N_c$



'tHooft instanton vertex

$$\bar{d} > \bar{u}$$

$$\sim \bar{u}_R u_L \bar{d}_R d_L$$

Spin Structure: SU(6) Proton Wave Function in CQM

The 3 quarks are **identical fermions** $\Rightarrow \psi$ **antisymmetric** under exchange

$$\psi = \psi(\text{color}) * \psi(\text{space}) * \psi(\text{spin}) * \psi(\text{flavor})$$

1 Color: All hadrons are color singlets = **antisymmetric**

$$\psi(\text{color}) = 1/\sqrt{6} (\text{RGB} - \text{RBG} + \text{BRG} - \text{BGR} + \text{GBR} - \text{GRB})$$

2 Space: proton has $l = l' = 0 \rightarrow \psi(\text{space}) = \mathbf{symmetric}$

3 Spin: $2 \otimes 2 \otimes 2 = (3_S \oplus 1_A) \otimes 2 = 4_S \oplus 2_{MS} \oplus 2_{MA}$

- 4_S symmetric states have spin 3/2, e.g. $\left| \frac{3}{2}, +\frac{3}{2} \right\rangle = \uparrow\uparrow\uparrow$

- 2_{MS} and 2_{MA} have spin 1/2 and **mixed symmetry**:

S or A under exchange of first 2 quarks only, e.g.

$$\left| \frac{1}{2}, +\frac{1}{2} \right\rangle_{MS} = (\uparrow\downarrow\uparrow + \downarrow\uparrow\uparrow - 2\uparrow\uparrow\downarrow)/\sqrt{6}$$

$$\left| \frac{1}{2}, +\frac{1}{2} \right\rangle_{MA} = (\uparrow\downarrow\uparrow - \downarrow\uparrow\uparrow)/\sqrt{2}$$

④ **Flavor**: symmetry groups SU(2)-spin and SU(3)-color are exact ...

- strong force is **flavor blind**
- constituent q masses **similar**: $m_u, m_d \approx 350$ MeV, $m_s \approx 500$ MeV

→ SU(3)-flavor is **approximate** for u, d, s

$$\text{SU(3)-flavor gives } 3 \otimes 3 \otimes 3 = 10_S \oplus 8_{MS} \oplus 8_{MA} \oplus 1_A$$

➤ **Proton**: $\psi(s=1/2)$ from spin $2_{MS}, 2_{MA}$ \otimes $\psi(uud)$ from flavor $8_{MS}, 8_{MA}$

$$|p^\uparrow\rangle = (u^\uparrow u^\downarrow d^\uparrow + u^\downarrow u^\uparrow d^\uparrow - 2u^\uparrow u^\uparrow d^\downarrow + 2 \text{ permutations}) / \sqrt{18}$$

➤ Count the number of quarks with spin up and spin down:

	$\langle p^\uparrow \hat{N}(u^\uparrow) p^\uparrow \rangle = \frac{30}{18} = \frac{5}{3}$	$\langle p^\uparrow \hat{N}(d^\uparrow) p^\uparrow \rangle = \frac{6}{18} = \frac{1}{3}$
	$\langle p^\uparrow \hat{N}(u^\downarrow) p^\uparrow \rangle = \frac{6}{18} = \frac{1}{3}$	$\langle p^\uparrow \hat{N}(d^\downarrow) p^\uparrow \rangle = \frac{12}{18} = \frac{2}{3}$

➤ Quark contributions to proton spin are:

	$\Delta u = N(u^\uparrow) - N(u^\downarrow) = +\frac{4}{3}$	$\Delta d = N(d^\uparrow) - N(d^\downarrow) = -\frac{1}{3}$
--	---	---

$$\Rightarrow \Delta\Sigma = \Delta u + \Delta d + \Delta s = 1$$

All spin present & accounted for!

CQM / SU(6) Scorecard

✓ Baryon Magnetic Moments

$$\mu_B = \sum_q \mu_q \Delta q \quad \text{where} \quad \mu_q \sim e_q/m_q$$

- take constituent quark masses
- take $\mu_u = -2\mu_d$, $\mu_s = 2\mu_d/3$
and fit μ_d to data

Note: $\mu_B \sim (e_q/m_q)\Delta q \sim |e_q|(\Delta q - \Delta \bar{q})$

\Rightarrow observable sensitive to **valence quarks**

B	Magnetic Moment		Expt
p	$(4\mu_u - \mu_d)/3$	2.7	2.79
n	$(4\mu_d - \mu_u)/3$	-1.8	-1.91
Σ^+	$(4\mu_u - \mu_s)/3$	2.6	2.48
Σ^-	$(4\mu_d - \mu_s)/3$	-1.0	-1.16
Ξ^0	$(4\mu_s - \mu_u)/3$	-1.4	-1.25
Ξ^-	$(4\mu_s - \mu_d)/3$	-0.5	-0.68
Λ	μ_s	-0.6	-0.61
$\Lambda\Sigma^0$	$(\mu_d - \mu_u)/\sqrt{3}$	-1.6	-1.60

✗ Hyperon β -Decay

- parity-violating weak decay
- decay products parallel to spin
- sensitive to $\sum_q (\Delta q + \Delta \bar{q})$

Decay Parameter	SU(6)	Expt
$F = (\Delta u - \Delta s)/2$	0.67	0.46
$D = (\Delta u - 2\Delta d + \Delta s)/2$	1.00	0.80

\Rightarrow Constituent Quark Model **lacks sea quarks**

Spin Structure of the Proton

$$\frac{1}{2} = \frac{1}{2}\Delta\Sigma + \Delta G + L_q + L_g$$

Parton Distribution Functions

unpolarized: $q(x) = q^\uparrow(x) + q^\downarrow(x)$

polarized: $\Delta q(x) = q^\uparrow(x) - q^\downarrow(x)$

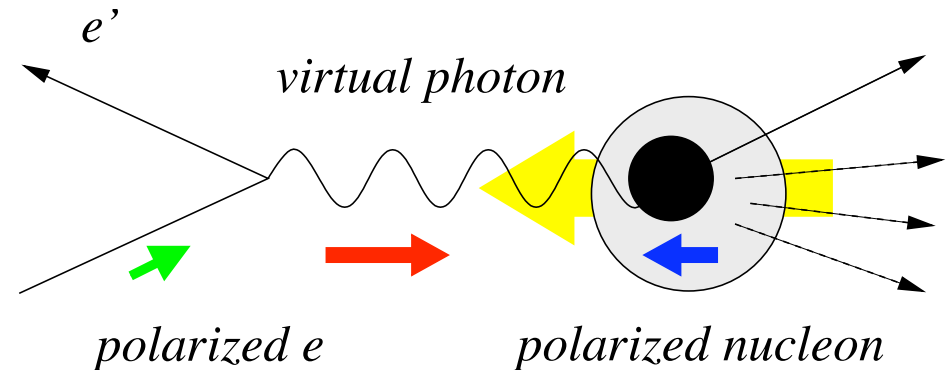
Constituent Quark Model

$$\Delta u = +\frac{4}{3}, \quad \Delta d = -\frac{1}{3} \quad \rightarrow \quad \Delta\Sigma = 1$$

Relativistic Quark Model

$$\Delta\Sigma \simeq 0.60 - 0.75 \quad L_q = \frac{1}{2}(1 - \Delta\Sigma)$$

Polarized Deep-Inelastic Scattering



From NLO-QCD analysis of inclusive DIS measurements + hyperon β -decay ...

- $\Delta\Sigma = 0.19 \pm 0.07$ ($\overline{\text{MS}}$ scheme)

The Spin Crisis!

Puzzle

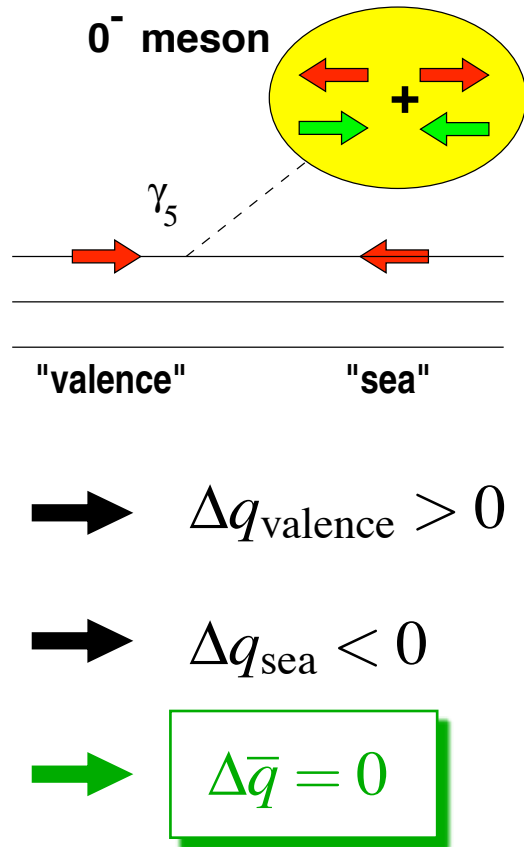
- $\Delta G = 1.0_{-0.6}^{+1.9}$ (AB scheme)

→ barely constrained, value > 0 favored

Anti-quark Spin in the Proton Sea

Meson Cloud Models

Li, Cheng, hep-ph/9709293

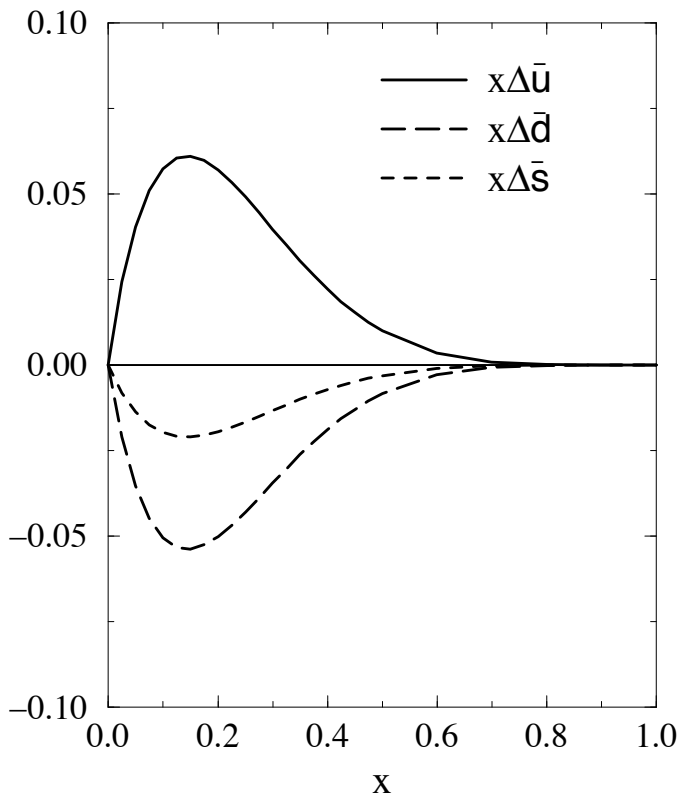


"Higher-order" cloud of vector mesons can generate a small polarization.

Chiral-Quark Soliton Model

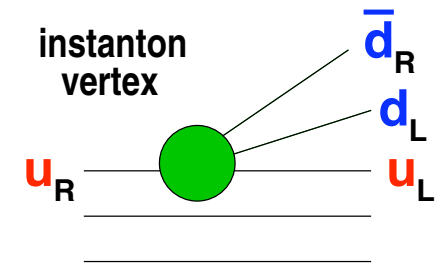
Light sea quarks polarized:

$$\Delta \bar{u} \simeq -\Delta \bar{d} > 0$$



Goeke et al, hep-ph/0003324

Instanton Mechanism

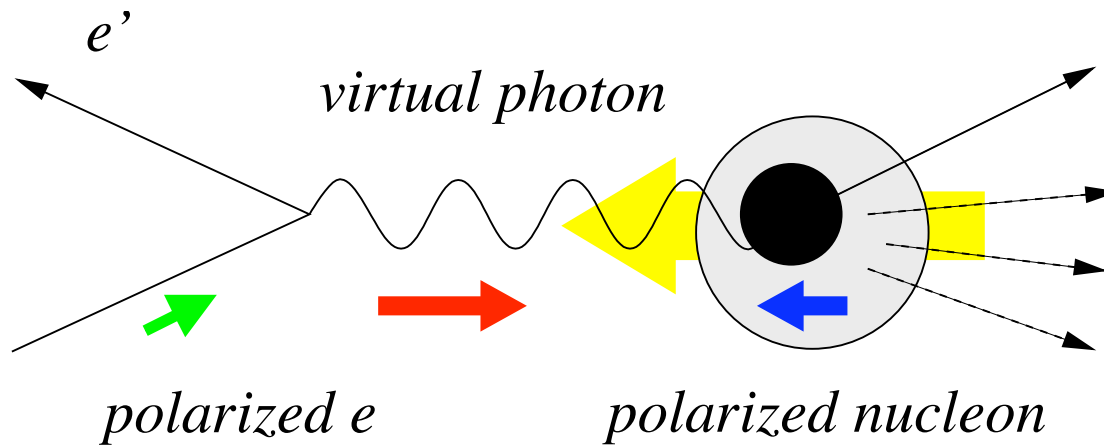


'tHooft instanton vertex $\sim \bar{u}_R u_L \bar{d}_R d_L$ transfers helicity from valence u quarks to $d\bar{d}$ pairs

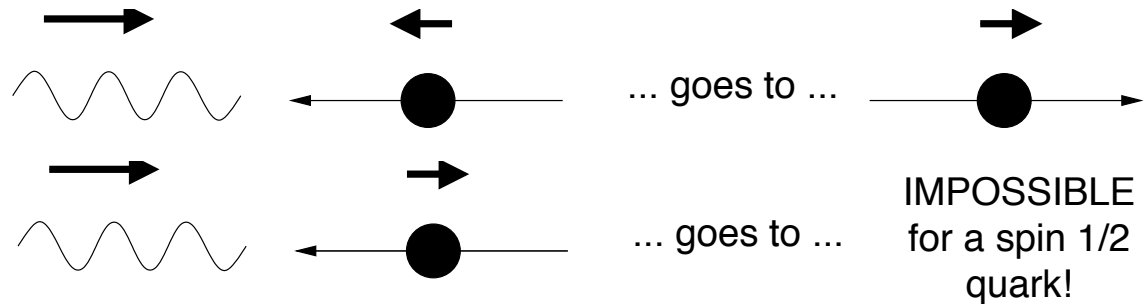
Quark Helicity
Distributions $\Delta q(x)$:
Results

Spin-Dependent Deep Inelastic Scattering (DIS)

Polarized lepton beams incident on polarized nucleon targets



The **polarized virtual photon** selects certain quark polarizations :



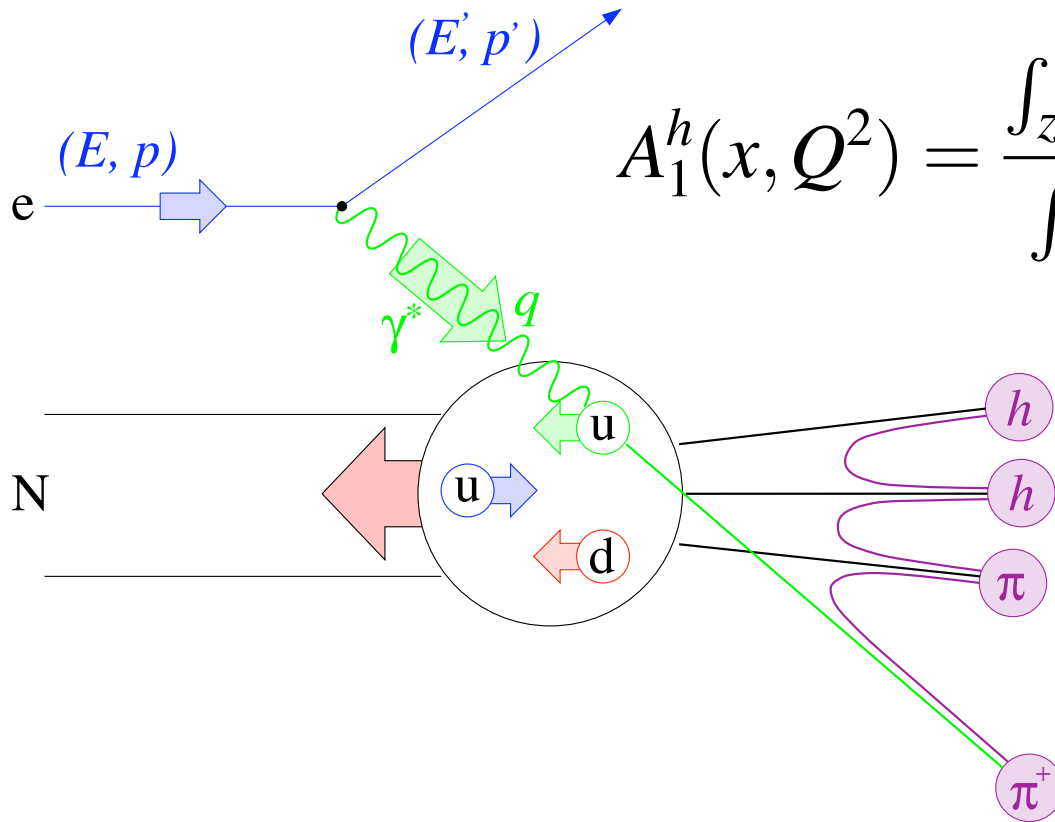
Double spin asymmetries are measured :

$$A_1 = \frac{\sigma_{1/2} - \sigma_{3/2}}{\sigma_{1/2} + \sigma_{3/2}} \simeq \frac{g_1}{F_1} = \frac{\sum_q e_q^2 \Delta q(x, Q^2)}{\sum_q e_q^2 q(x, Q^2)}$$

Polarized Semi-Inclusive DIS (SIDIS)

In SIDIS, a hadron h is detected in coincidence with the scattered lepton:

Flavor Tagging: Flavor content of observed hadron h is related to flavor of **struck quark q** via the **fragmentation functions $D(z)$**



$$A_1^h(x, Q^2) = \frac{\int_{z_{min}}^1 dz \sum_q e_q^2 \Delta q(x, Q^2) \cdot D_q^h(z, Q^2)}{\int_{z_{min}}^1 dz \sum_q e_q^2 q(x, Q^2) \cdot D_q^h(z, Q^2)}$$

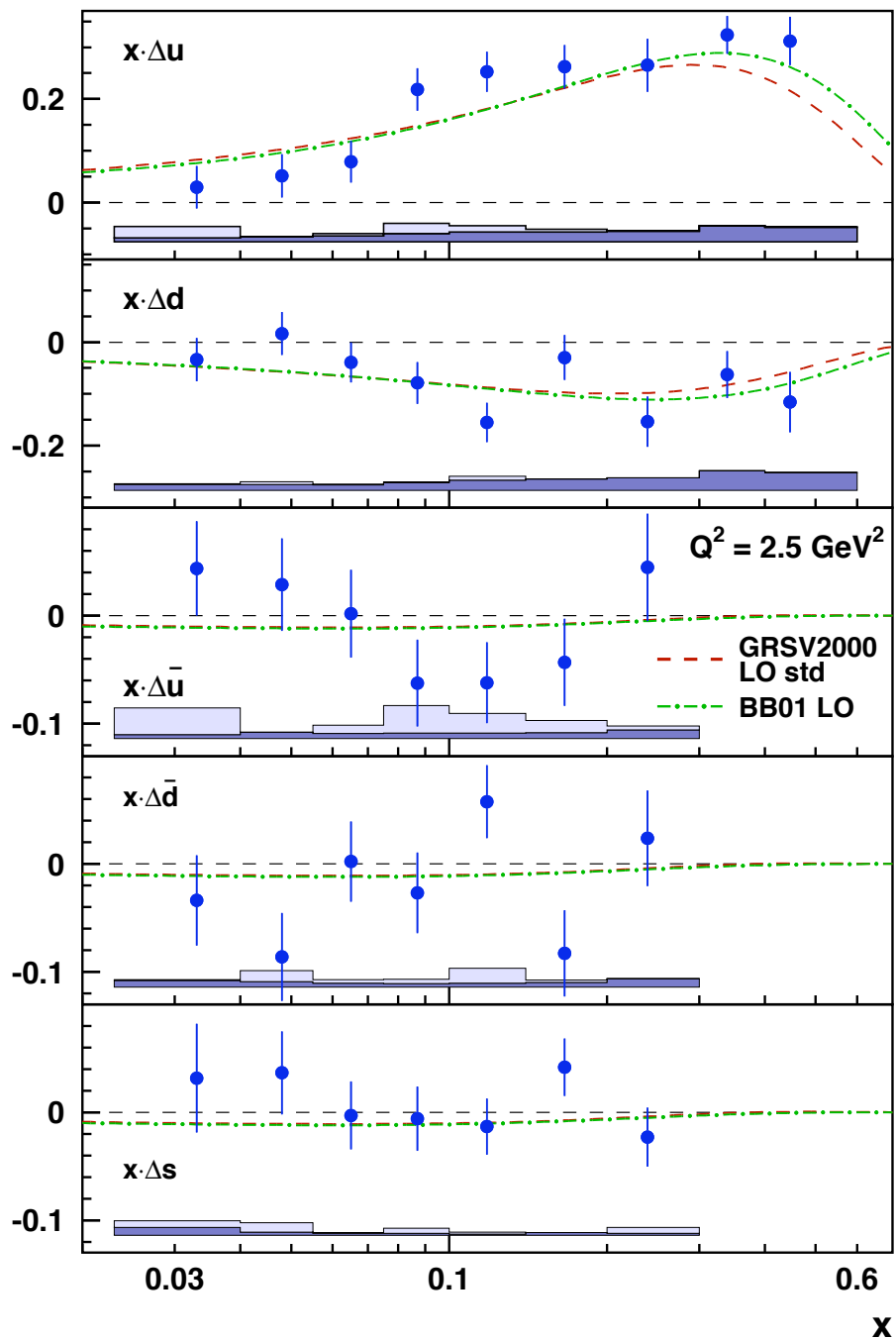
Rewriting ...

$$A_1^h(x, z) = \sum_q P_q^h(x, z) \frac{\Delta q(x)}{q(x)}$$

Purity matrix P_q^h = probability that hadron h came from struck quark q

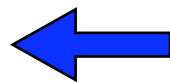
Purities are spin-independent ... compute using Monte Carlo tuned to *unpolarized* data

Final Δq Measurement from HERMES Run 1



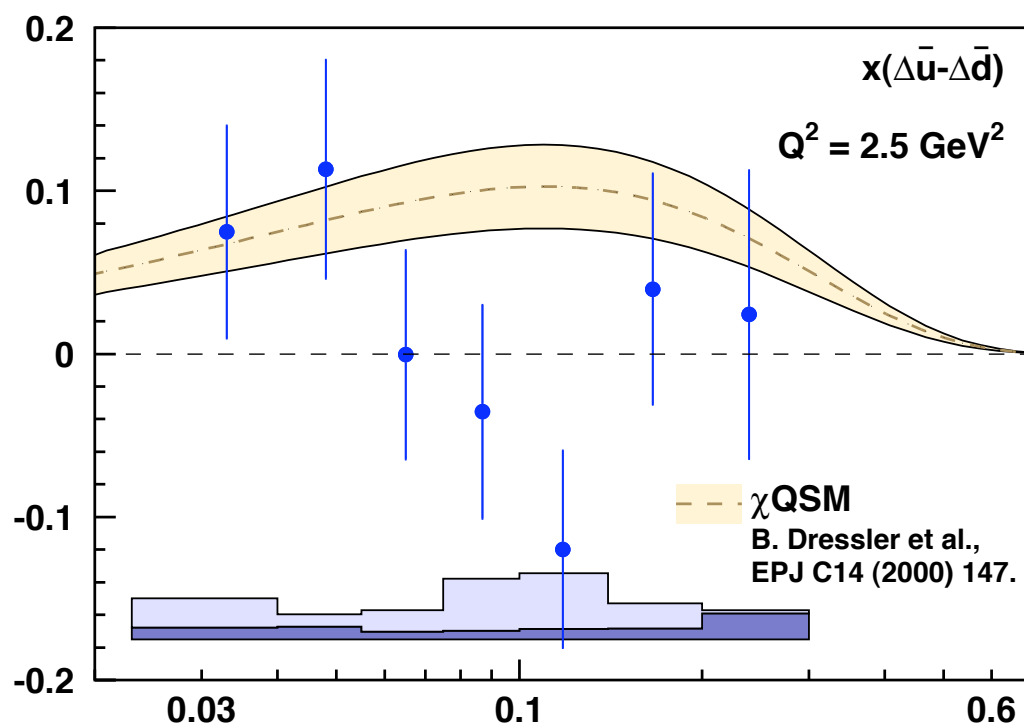
using polarized data 1996–2000:

$$A_{1,p}, A_{1,p}^{\pi^{\pm}}, A_{1,d}, A_{1,d}^{\pi^{\pm}}, A_{1,d}^{K^{\pm}}$$



First 5-flavor fit to $\Delta q(x)$

No evidence of anti-quark polarization, or flavor-asymmetry thereof



New Analysis: Isoscalar extraction of Δs

Extract isoscalar combinations of $\Delta q(x)$:

$$\Delta S(x) \equiv \Delta s(x) + \Delta \bar{s}(x)$$

$$\Delta Q(x) \equiv \Delta u(x) + \Delta \bar{u}(x) + \Delta d(x) + \Delta \bar{d}(x)$$

Asymmetries measured from isoscalar deuteron data:

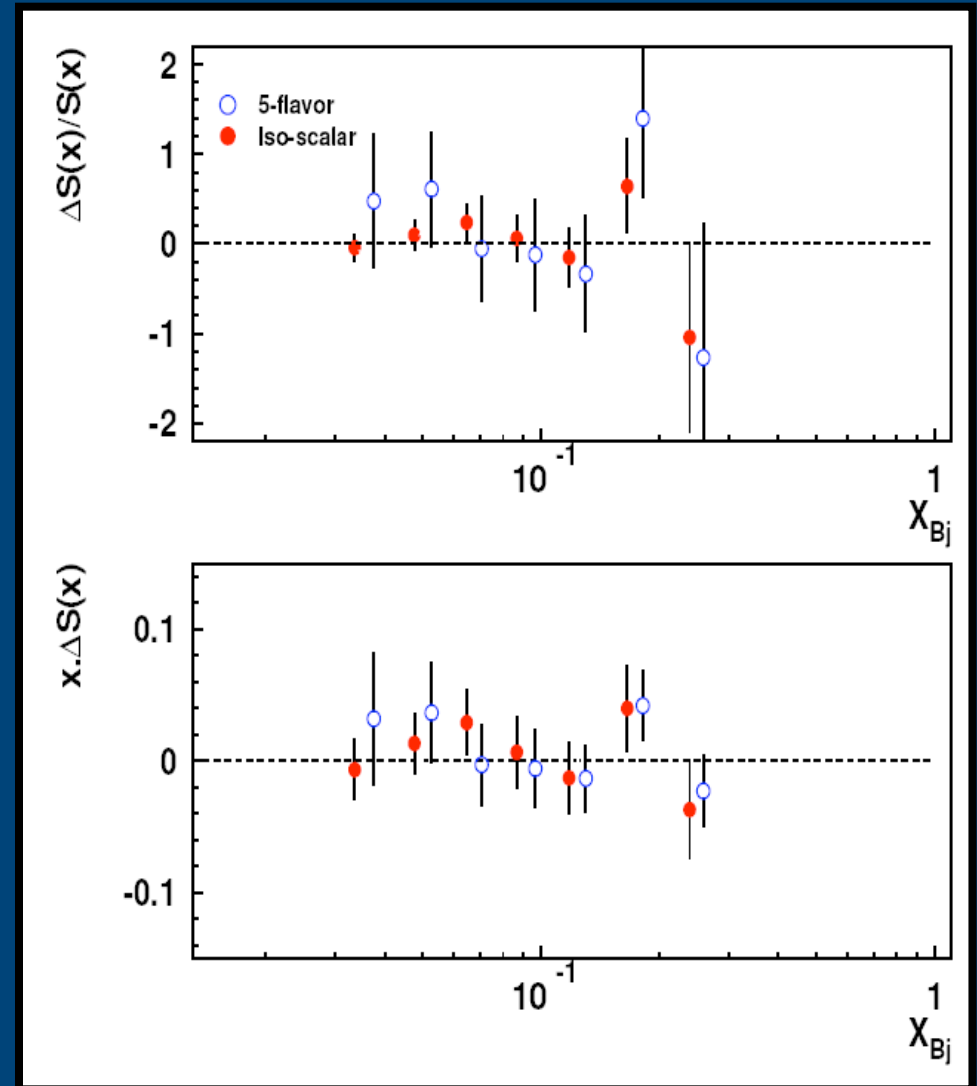
$$\begin{pmatrix} A_d(x) \\ A_d^{K^+ + K^-}(x) \end{pmatrix} = C_R \begin{pmatrix} P_q & P_s \\ P_q^{K^+ + K^-} & P_s^{K^+ + K^-} \end{pmatrix} \begin{pmatrix} \Delta Q(x)/Q(x) \\ \Delta S(x)/S(x) \end{pmatrix}$$

• Inclusive purities are simple combinations of unpolarized PDFs.

$$P_Q(x) = \frac{5Q(x)}{5Q(x) + 2S(x)}, P_S(x) = \frac{2S(x)}{5Q(x) + 2S(x)}$$

• Kaon purities can be computed from the unpolarized K multiplicity assuming only charge symmetry in fragmentation.

$$D_q^{K^+ + K^-}(x) = D_{\bar{q}}^{K^+ + K^-}$$



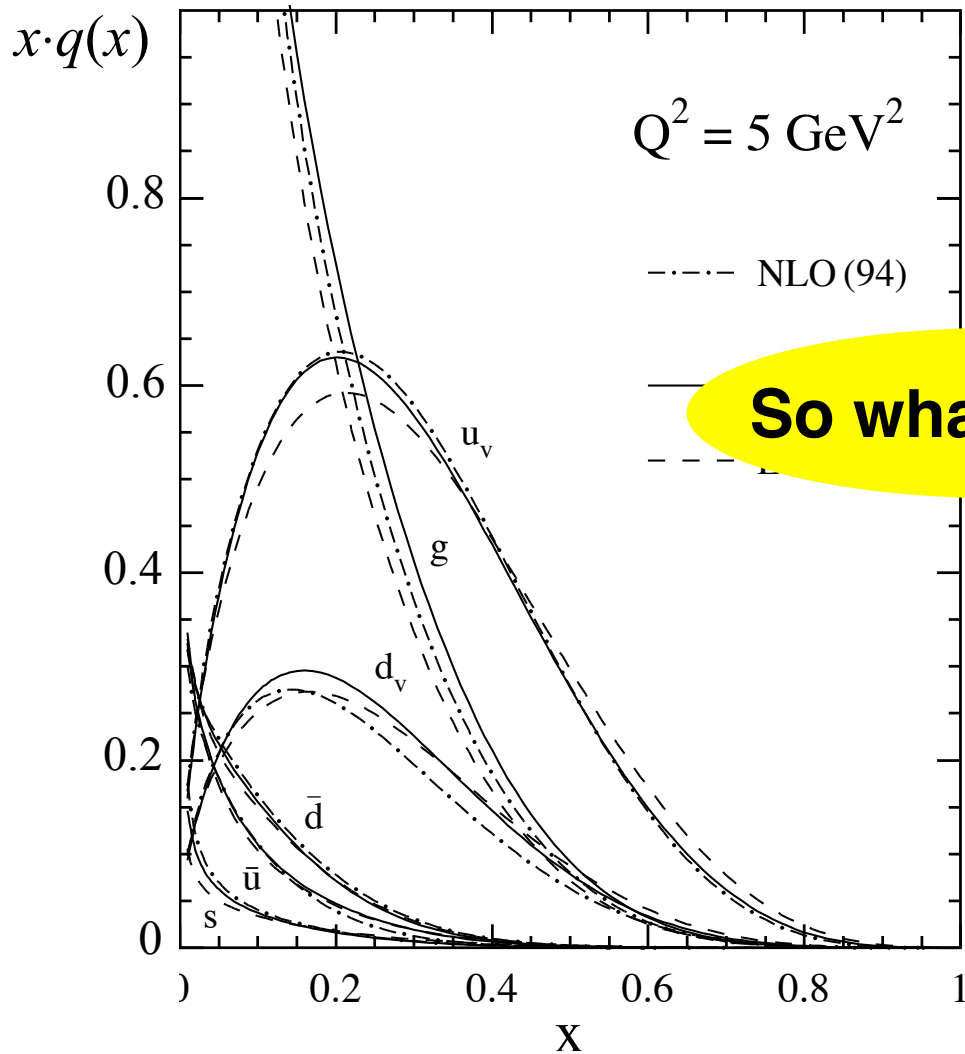
Excellent agreement -- No MC Dependence



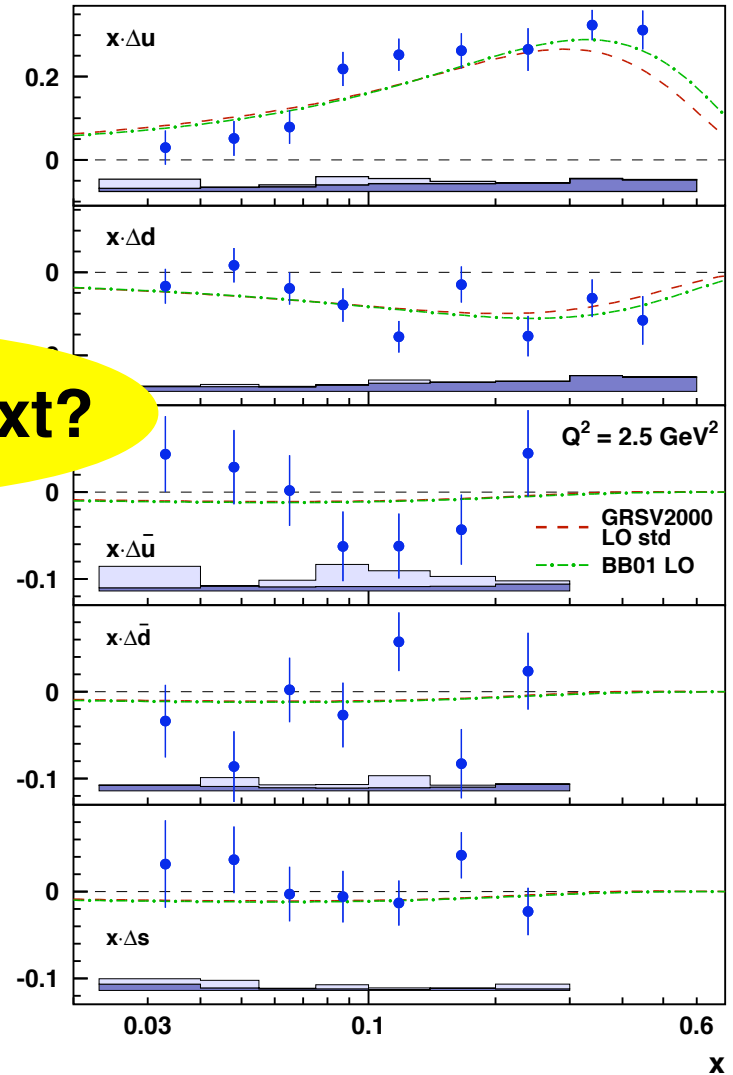
Single-Spin Asymmetries

Unpolarized PDF's

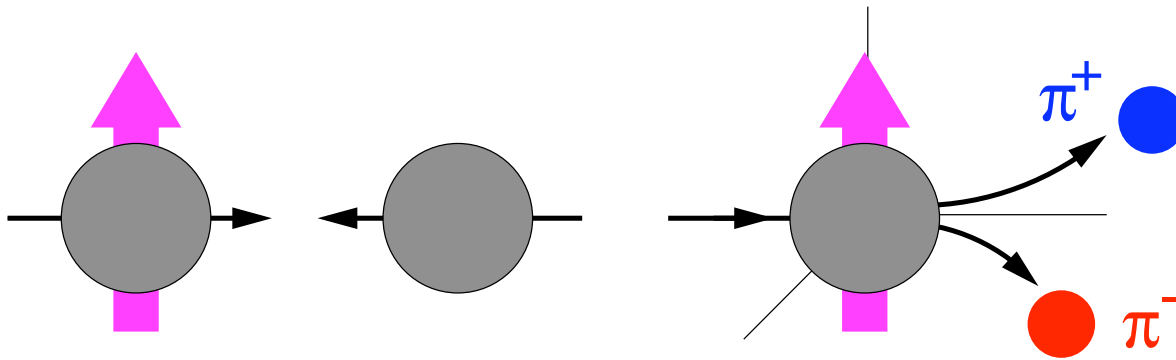
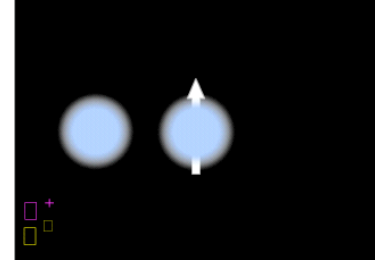
Polarized PDF's



So what's next?



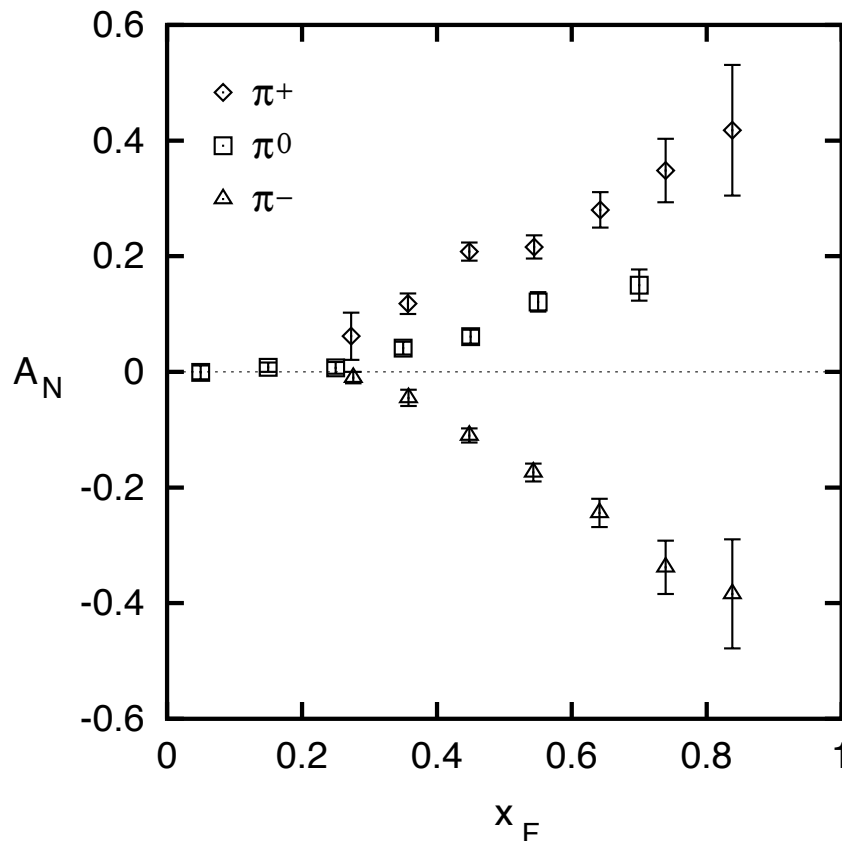
Fermilab E704: $p^\uparrow p \rightarrow \pi X$ at 400 GeV



Analyzing Power

$$A_N = \frac{1}{P_{\text{beam}}} \frac{N_{\text{left}}^\pi - N_{\text{right}}^\pi}{N_{\text{left}}^\pi + N_{\text{right}}^\pi}$$

Huge single-spin asymmetry !



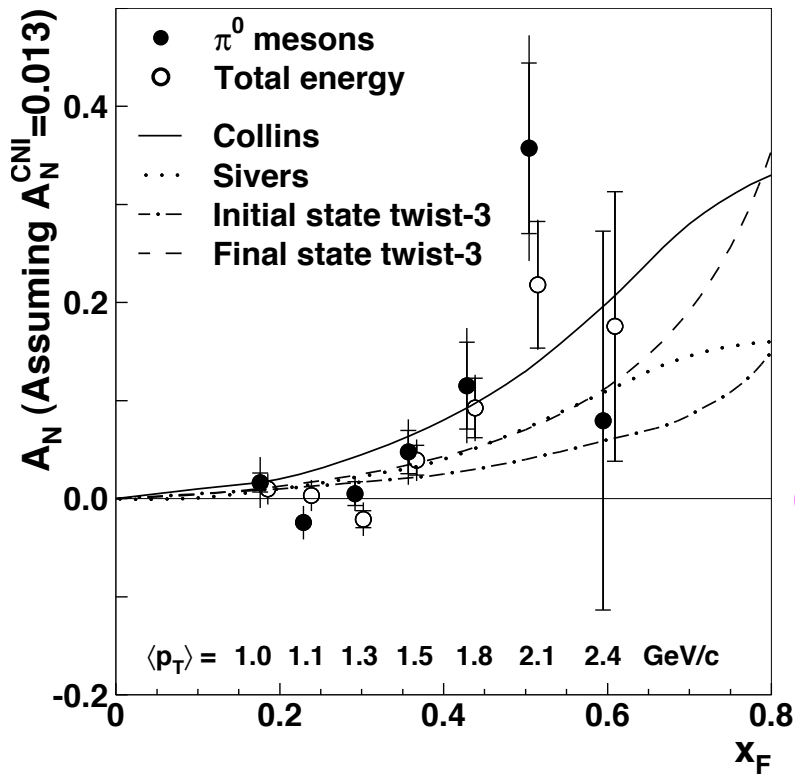
- Opposite sign for $\pi^+ = u\bar{d}$ than for $\pi^- = d\bar{u}$
- Effect larger for forward production
- Observable: $\vec{S}_{\text{beam}} \cdot (\vec{p}_{\text{beam}} \times \vec{p}_\pi)$
odd under naive Time-Reversal

Surprising observation! Why?

SSA's at high-energies

Now confirmed at STAR
at much higher energies

T-odd observables

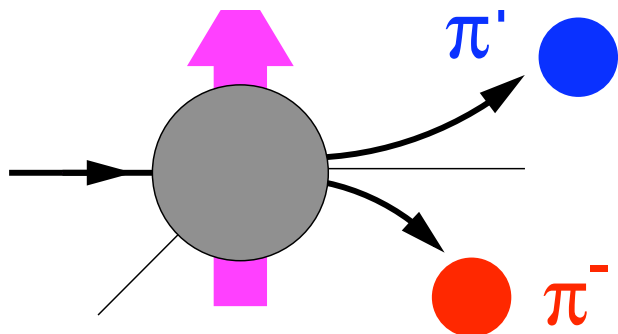


SSA observables $\sim \vec{J} \cdot (\vec{p}_1 \times \vec{p}_2)$
 \Rightarrow **odd** under naive **time-reversal**

Since QCD amplitudes are T-even, must arise from interference between spin-flip and non-flip amplitudes with different phases

Can't come from perturbative subprocess xsec:

- q helicity flip suppressed by m_q/\sqrt{s}
- need α_s -suppressed loop-diagram to generate necessary phase

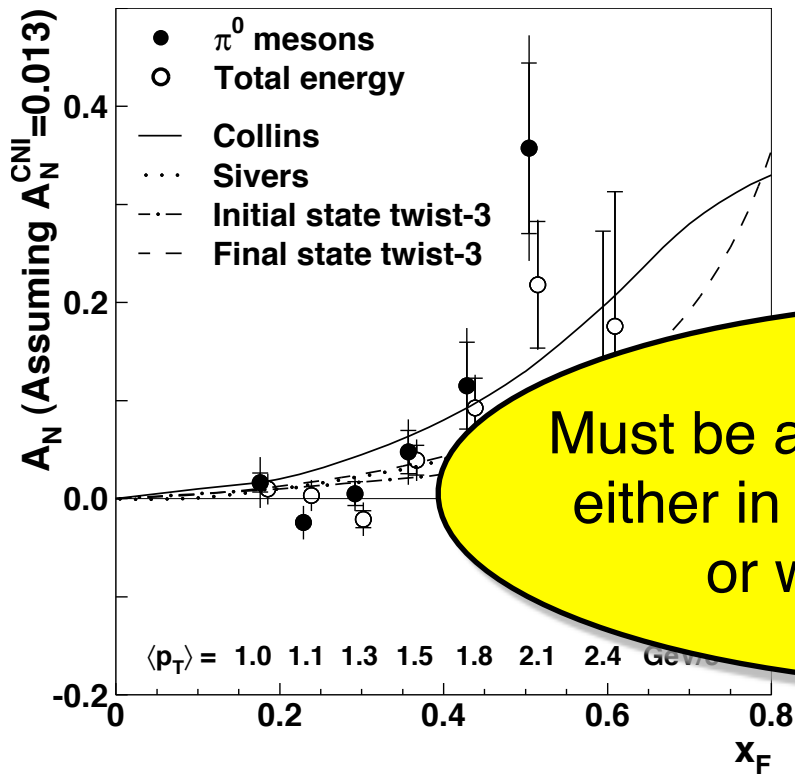


At hard (enough) scales, SSA's must arise from soft physics: **T-odd distribution / fragmentation functions**

SSA's at high-energies

Now confirmed at STAR
at much higher energies

T-odd observables



Must be a new, **spin-orbit structure**
either in the fragmentation process
or within the proton itself

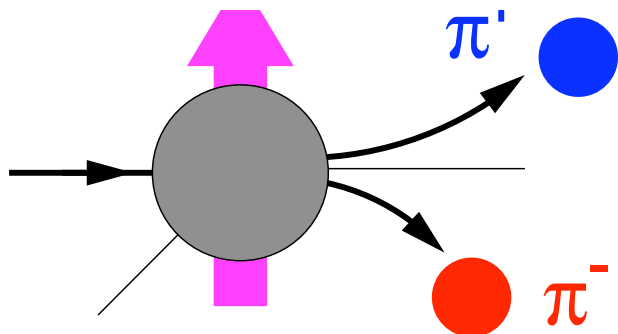
SSA observables $\sim \vec{J} \cdot (\vec{p}_1 \times \vec{p}_2)$
 \Rightarrow **odd** under naive **time-reversal**

Since QCD amplitudes are T-even, must arise
from interference between **spin-flip** and
different phases

subprocess xsec:

suppressed by m_q/\sqrt{s}

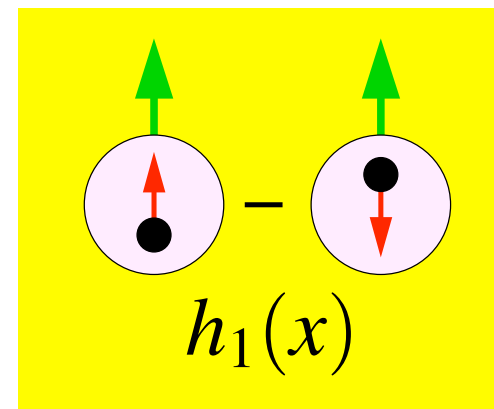
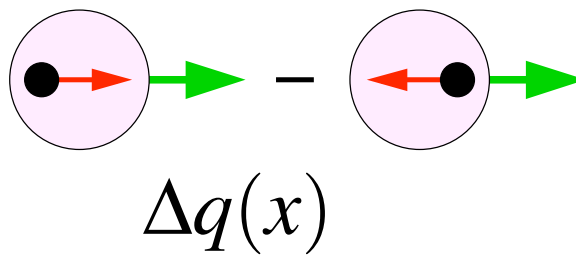
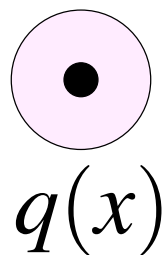
● need α_s -suppressed loop-diagram to
generate necessary phase



At hard (enough) scales, SSA's must
arise from soft physics: **T-odd distribution /**
fragmentation functions

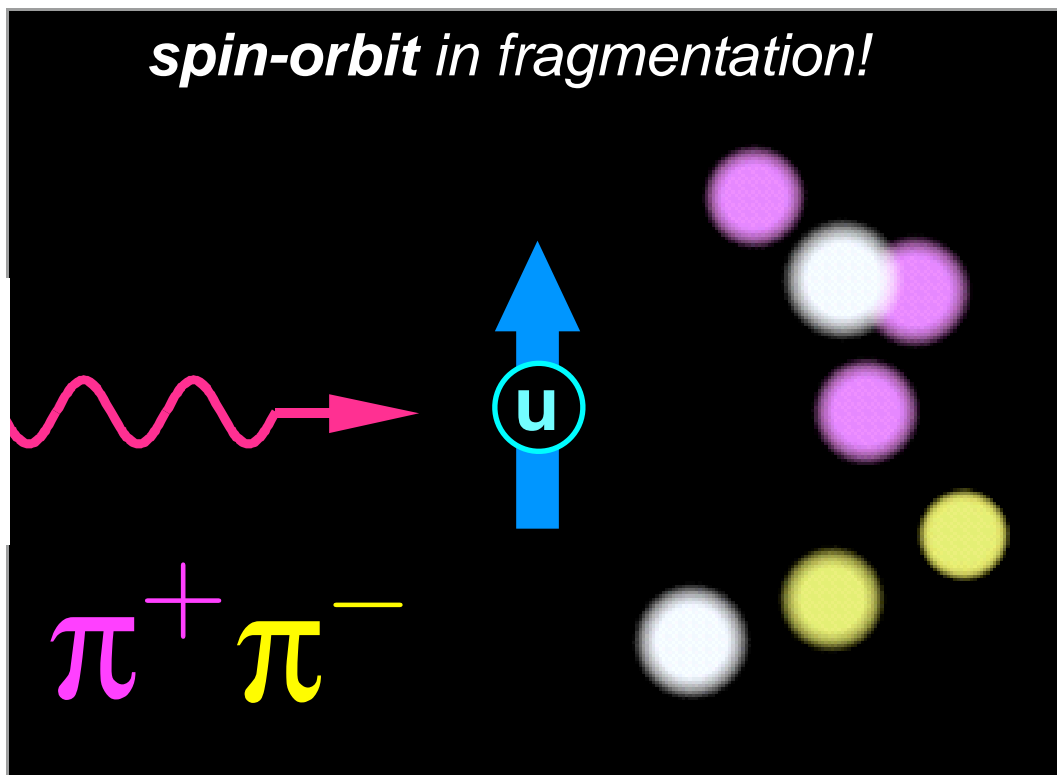
E704 Possible Mechanism #1: The "Collins Effect"

Need an ordinary distribution function ... **transversity**

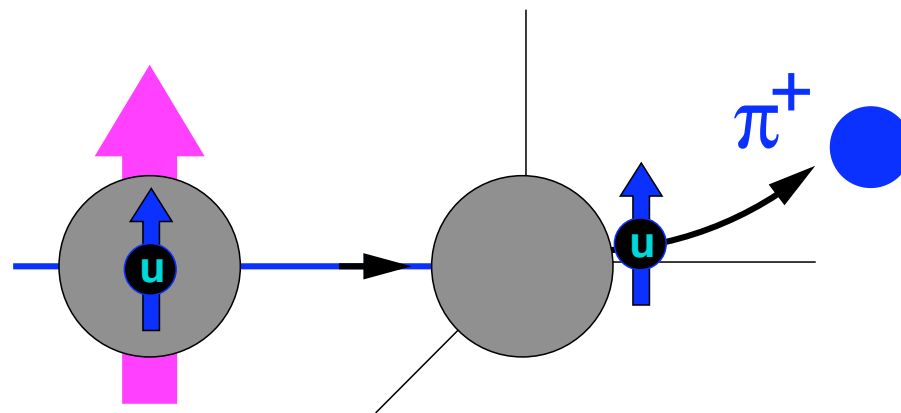


... with a new, **T-odd "Collins" fragmentation function**

$$H_1^\perp(z, p_T)$$



E704 effect:



$$h_1(x) \otimes H_1^\perp(z, p_T)$$

Transversity: The Third Structure Function

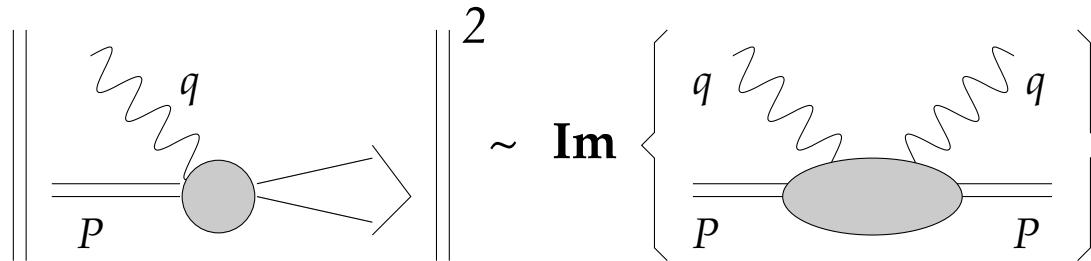
Proton Matrix Elements

vector charge $\langle PS | \bar{\psi} \gamma^\mu \psi | PS \rangle = \int_0^1 dx q(x) - \bar{q}(x) \rightarrow \# \text{ valence quarks}$

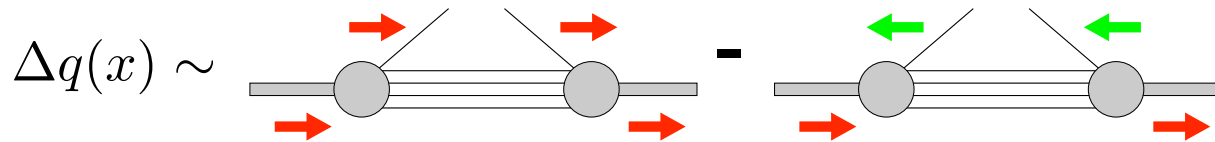
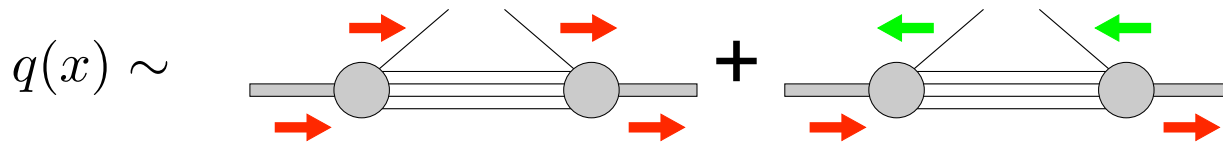
axial charge $\langle PS | \bar{\psi} \gamma^\mu \gamma_5 \psi | PS \rangle = \int_0^1 dx \Delta q(x) + \Delta \bar{q}(x) \rightarrow \text{net quark spin}$

tensor charge $\langle PS | \bar{\psi} \sigma^{\mu\nu} \gamma_5 \psi | PS \rangle = \int_0^1 dx \delta q(x) - \delta \bar{q}(x) \rightarrow ???$

Forward Helicity Amplitudes



(optical theorem applied to DIS)



Properties of Transversity

- **In Non-Relativistic Case**, boosts and rotations commute:
... but bound quarks are highly *relativistic* in nature

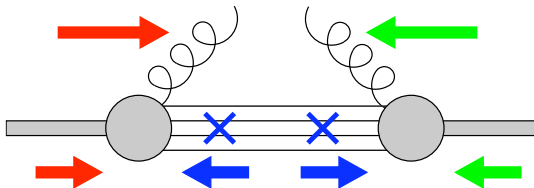
$$\delta q(x) = \Delta q(x)$$

- **No Gluons**

Angular momentum conservation: $\Lambda - \lambda = \Lambda' - \lambda'$

⇒ transversity has *no gluon* component

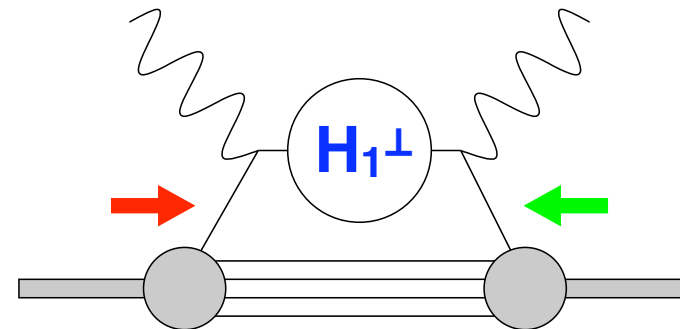
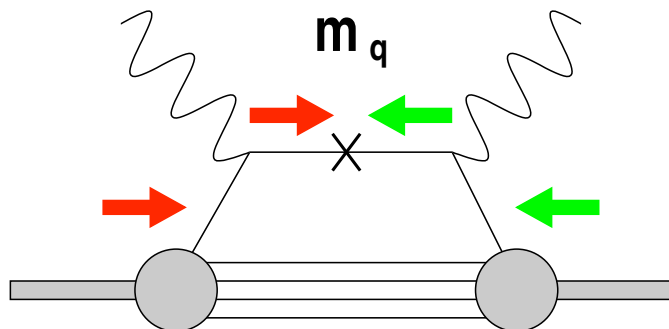
⇒ different Q^2 *evolution* than $\Delta q(x)$



- **Chiral Odd**

Helicity flip amplitudes occur only at $\mathcal{O}(m_q/Q)$ in inclusive DIS ...

but they are observable in e.g. **semi-inclusive** reactions



Properties of Transversity

- **In Non-Relativistic Case**, boosts and rotations commute:
... but bound quarks are highly *relativistic* in nature

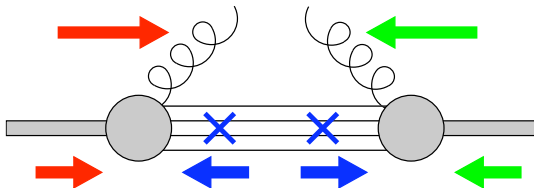
$$\delta q(x) = \Delta q(x)$$

- **No Gluons**

Angular momentum conservation: $\Lambda - \lambda = \Lambda' - \lambda'$

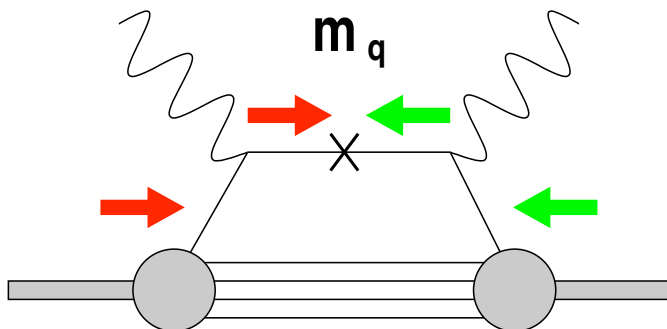
⇒ transversity has *no gluon* component

⇒ different Q^2 *evolution* than $\Delta q(x)$



- **Chiral Odd**

Helicity flip amplitudes occur only at $\mathcal{O}(m_q/Q)$ in inclusive DIS ...



tensor charge =
“*pure valence*” object

→ promising for LQCD
comparison?

E704 Possible Mechanism #2: The “Sivers Effect”

Need the ordinary fragmentation function

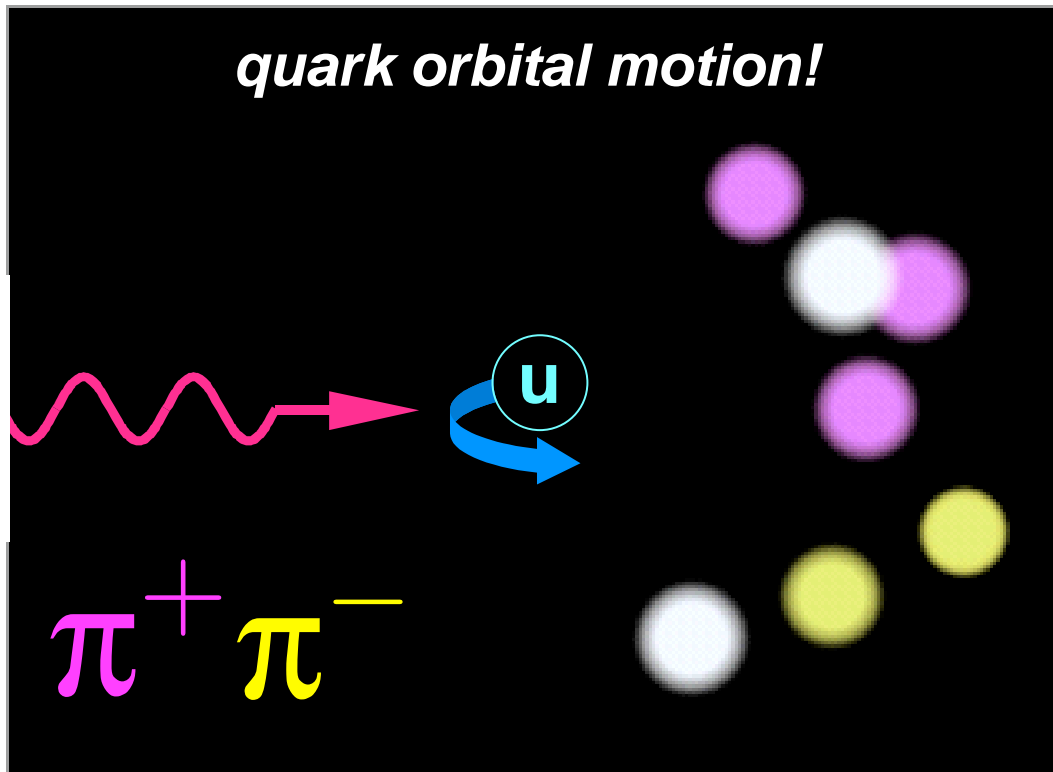
$$D_1(z)$$

... with a new, **T-odd “Sivers” distribution function**

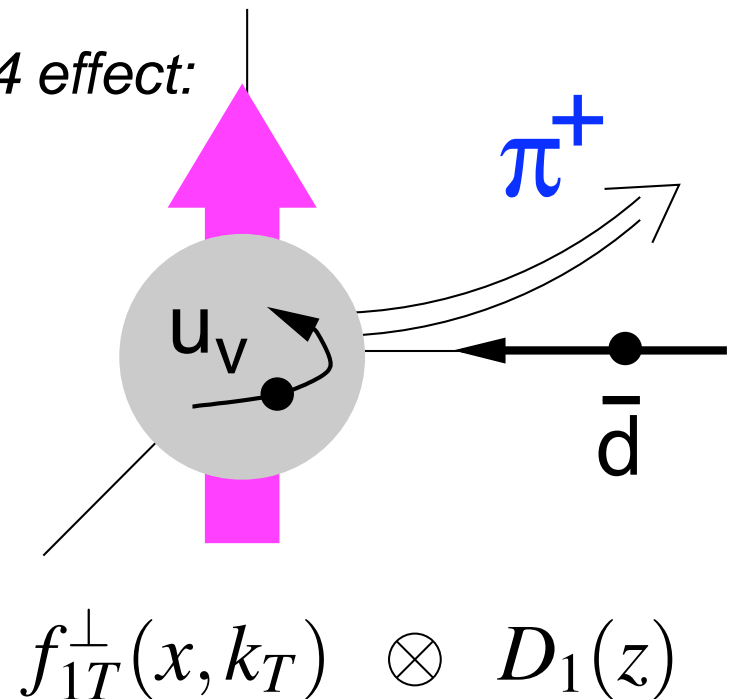
$$f_{1T}^\perp(x, k_T)$$

Phenomenological model of **Meng & Chou**:

Forward π^+ produced from **orbiting valence-u quark** by recombination at front surface of beam protons



E704 effect:



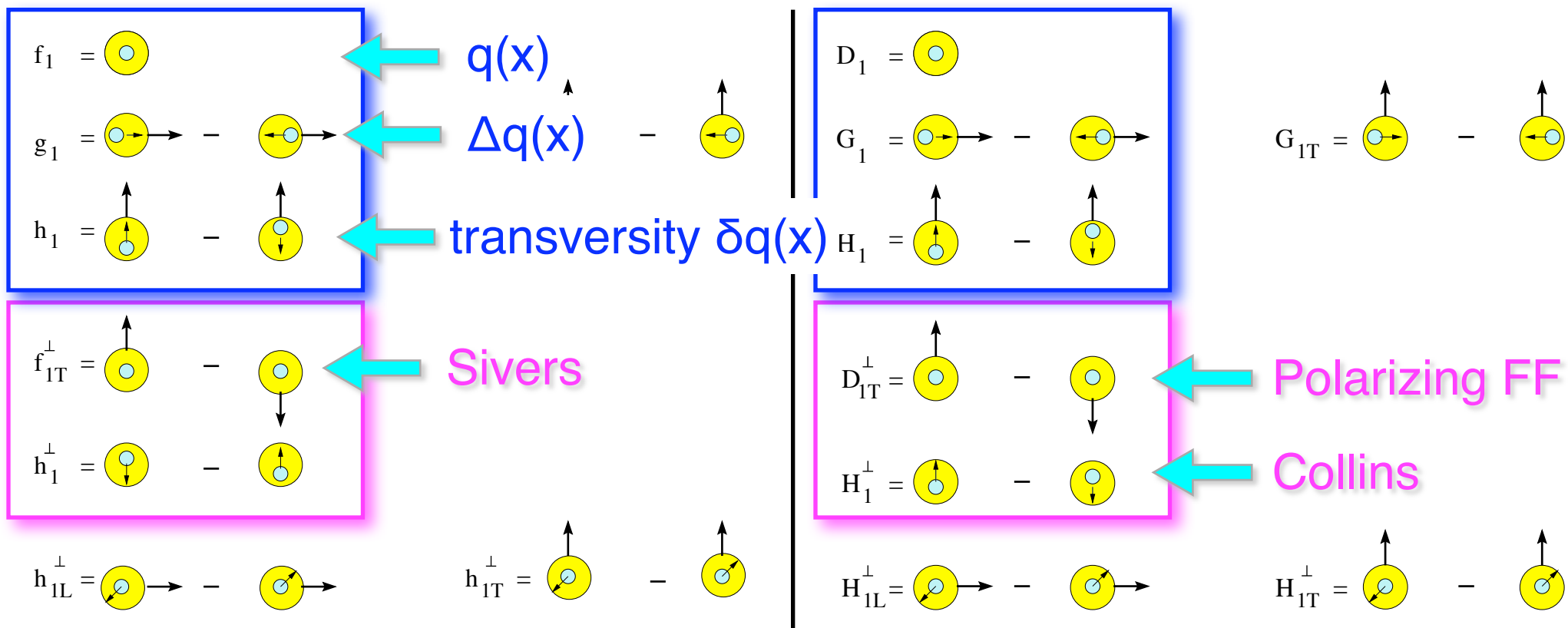
Functions surviving on integration over Transverse Momentum

The others are sensitive to *intrinsic* k_T in the nucleon & in the fragmentation process

Mulders & Tangerman, NPB 461 (1996) 197

Distribution Functions

Fragmentation Functions

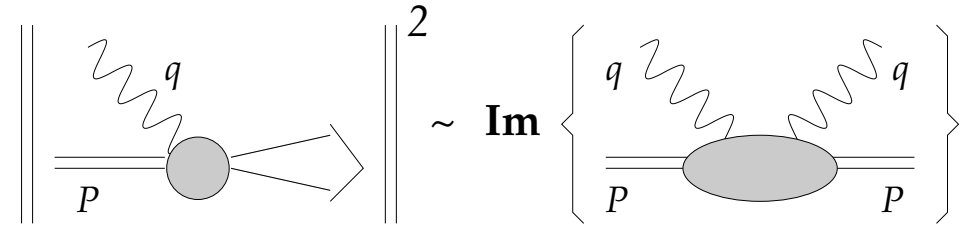


Functions Odd under naive Time Reversal

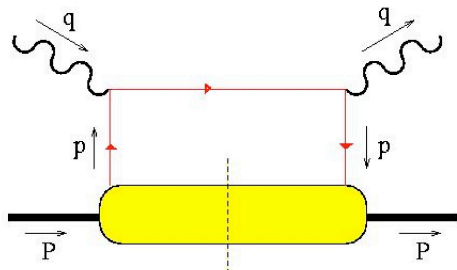
One T-odd function required to produce *single-spin asymmetries* in SIDIS

The Leading-Twist Sivers Function: Can it Exist in DIS?

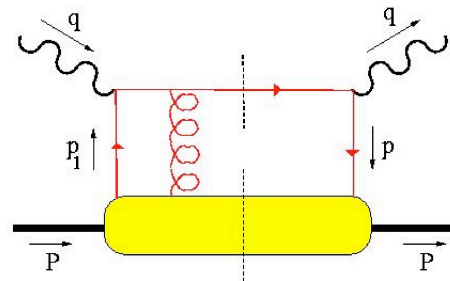
A T-odd function like f_{1T}^\perp **must** arise from **interference** ... but a distribution function is just a forward scattering amplitude, how can it contain an interference?



Brodsky, Hwang, & Schmidt 2002



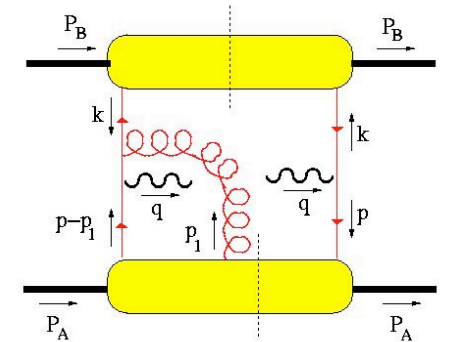
can interfere with



and produce a T-odd effect!
(also need $L_z \neq 0$)

It looks like higher-twist ... but no, these are soft gluons = “gauge links” required for color gauge invariance

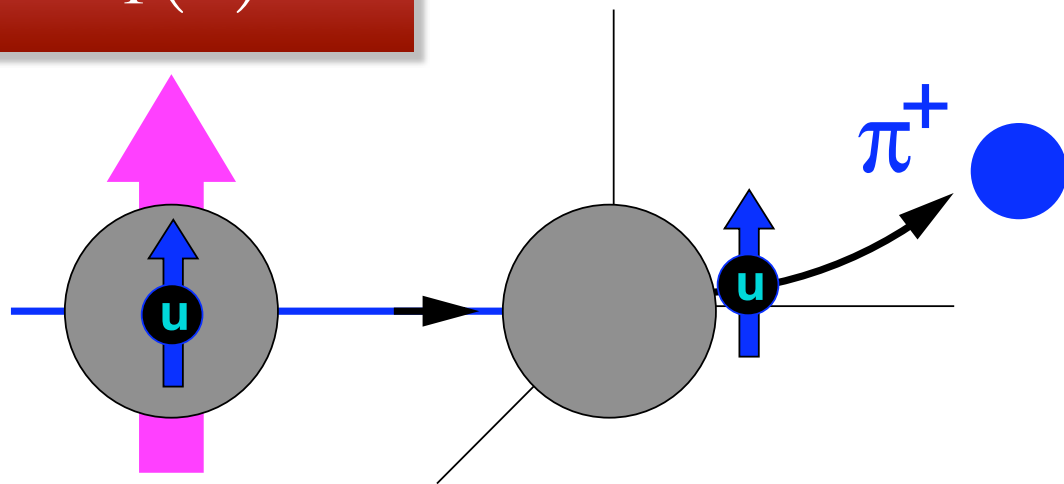
Such soft-gluon reinteractions with the soft wavefunction are **final (or initial) state interactions** ... and may be **process dependent!** → new **universality issues**



e.g. Drell-Yan

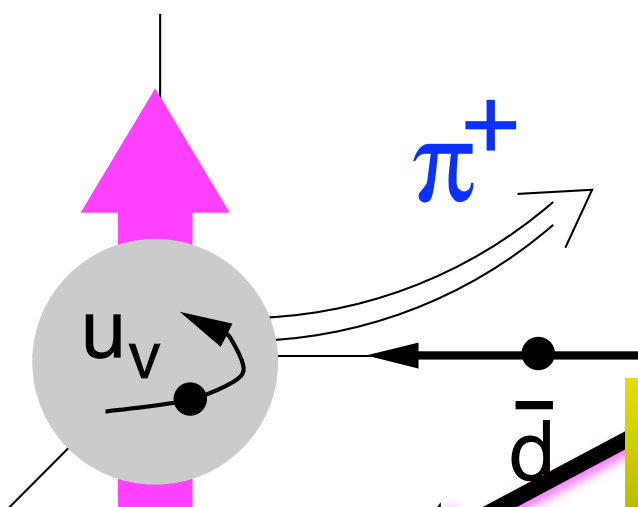
Transversity
 $h_1(x)$

Photo-Album of our New Friends!



Collins Function
 $H_1^\perp(z, p_T)$

T-Odd observables require interference between a spin-flip and a non-flip amplitude



Sivers Function
 $f_{1T}^\perp(x, k_T)$

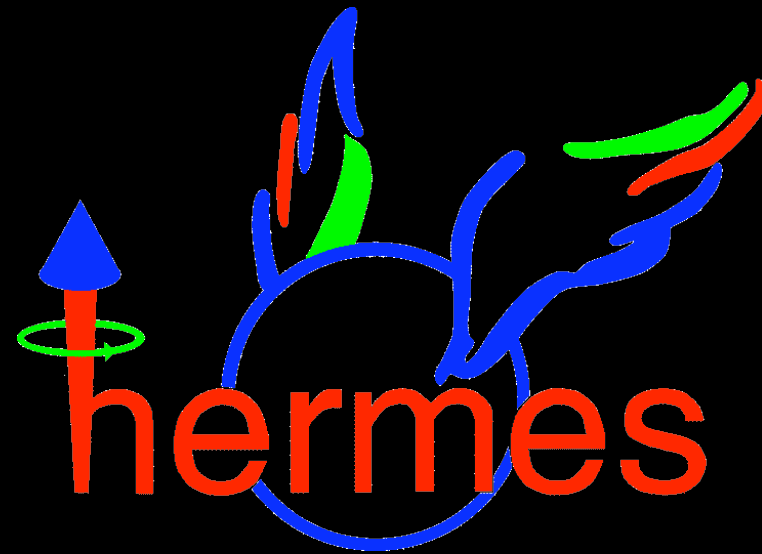
Favored / Disfavored Frag Functions

$$D_{\text{fav}} \equiv D^{u \rightarrow \pi^+} = D^{d \rightarrow \pi^-} = \dots$$

$$D_{\text{dis}} \equiv D^{u \rightarrow \pi^-} = D^{d \rightarrow \pi^+} = \dots$$

In Search of T-Odd Functions:

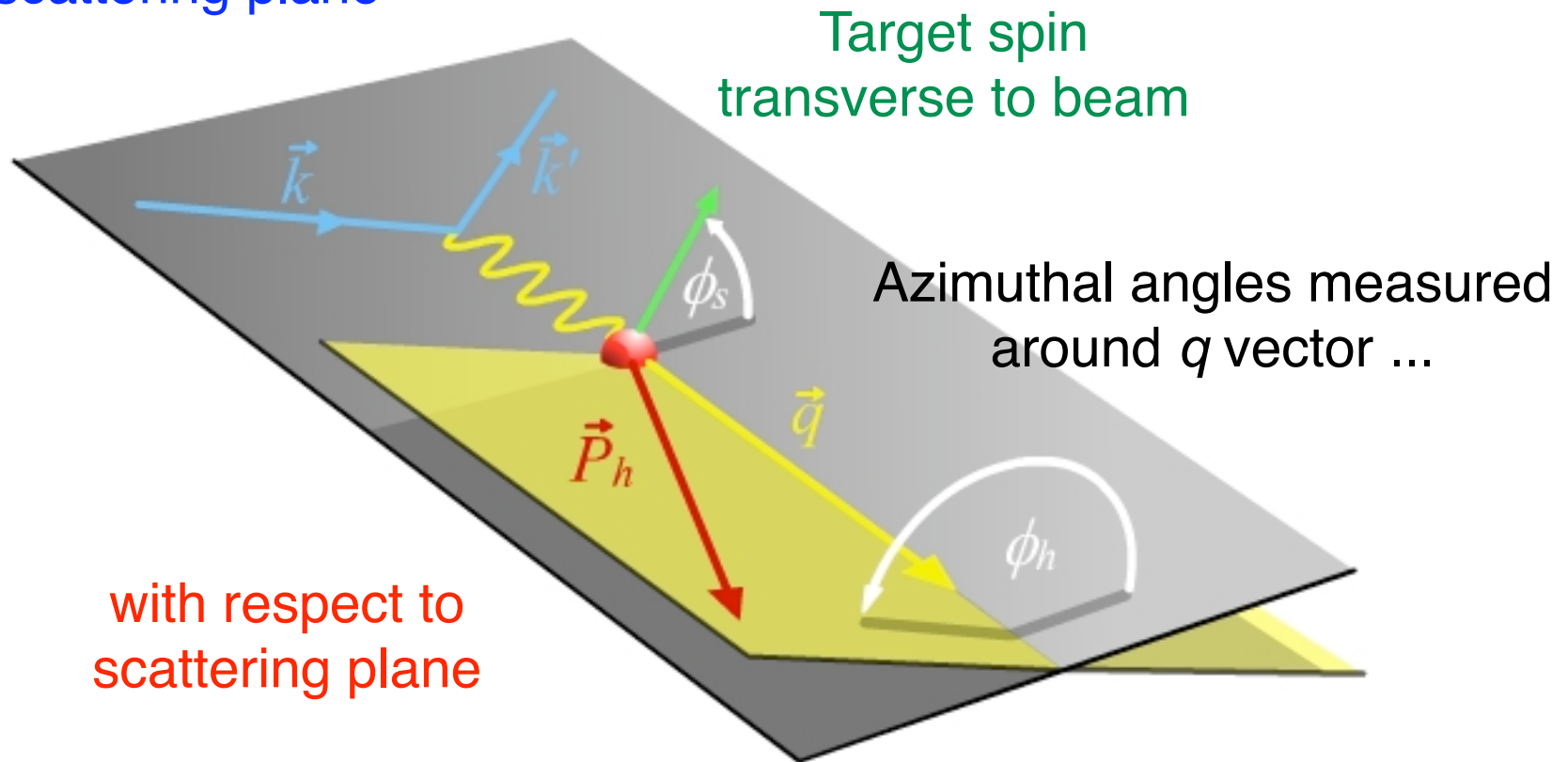
HERMES Run 2



Electro-Production of Pions with Transverse Target

Switched from longitudinal to **transverse** target polarization in 2002 ...
Measure dependence of pion production on **two azimuthal angles**

Electron beam defines
scattering plane



ϕ_S = target spin orientation

ϕ_h = pion ("hadron") direction

Separating the Collins & Sivers Mechanisms

Collins mechanism

$$\delta q(x) \otimes H_1^\perp(z, k_T) \Rightarrow \sin(\phi_h + \phi_s)$$

Sivers mechanism

$$f_{1T}^\perp(x, k_T) \otimes D_1(z) \Rightarrow \sin(\phi_h - \phi_s)$$

Measure **azimuthal moments** of SIDIS
xsec to **separate the mechanisms**

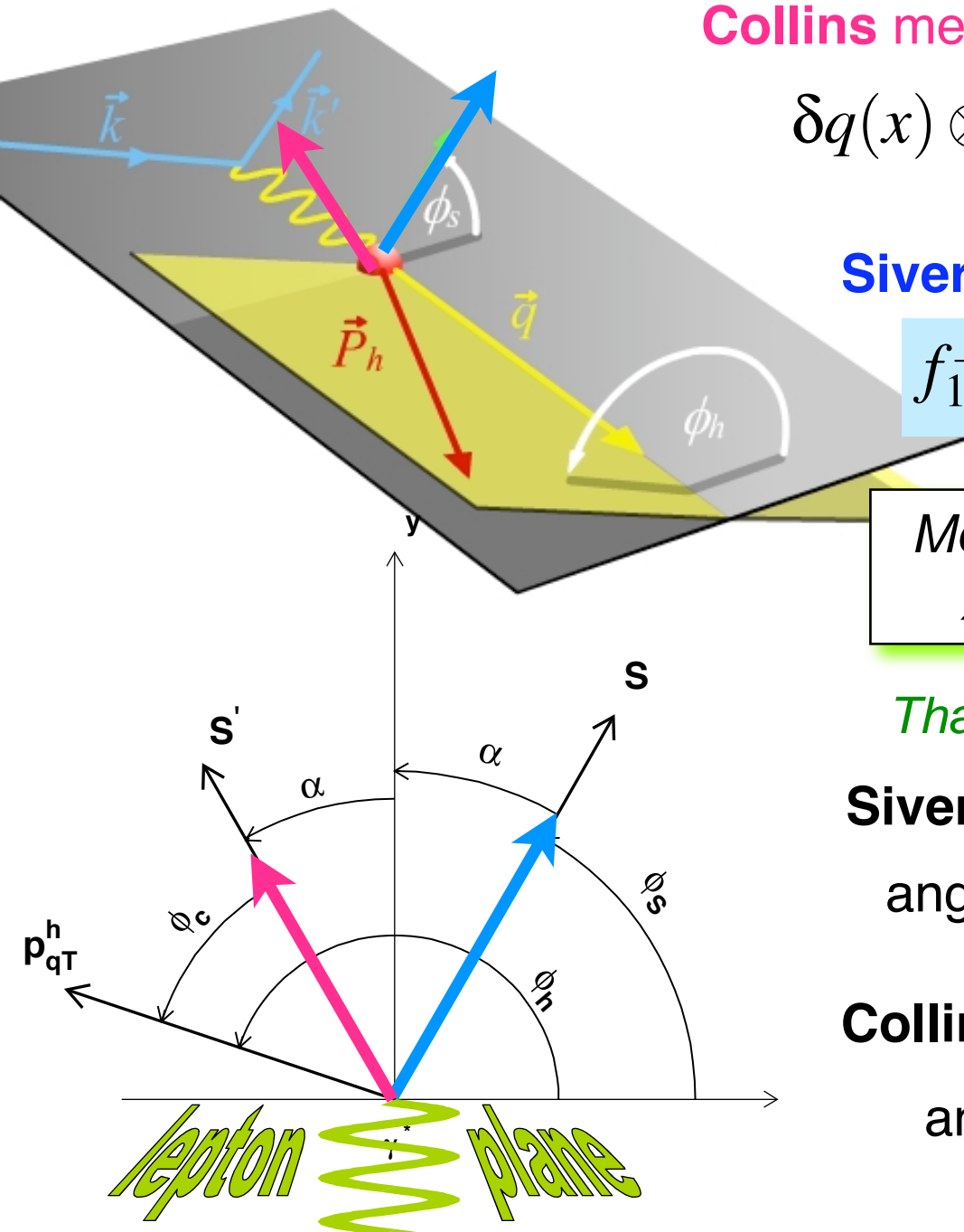
Thanks to **linear polarization** of photon ...

Sivers: $(\phi_h - \phi_s)$

angle of pion relative to **initial** quark spin

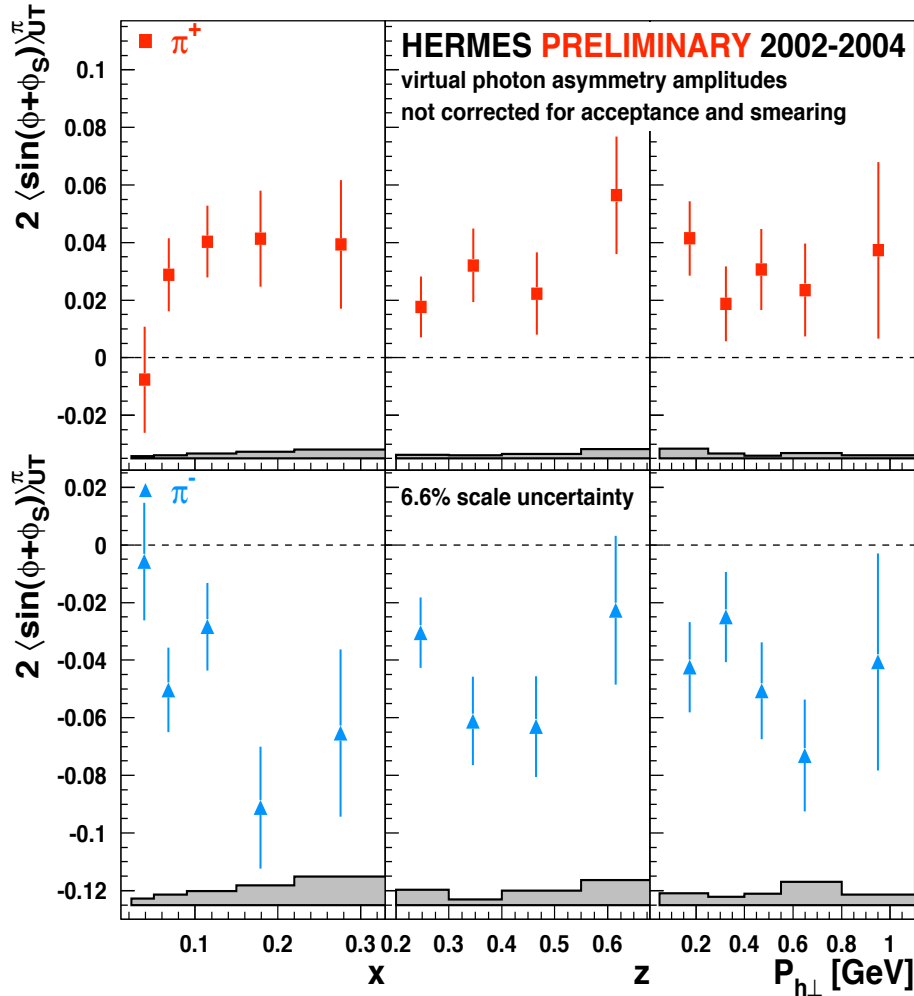
Collins: $(\phi_h + \phi_s) = \pi + (\phi_h - \phi_s)$

angle of pion relative to **final** quark spin



*SSA Results 1:
Collins Effect*

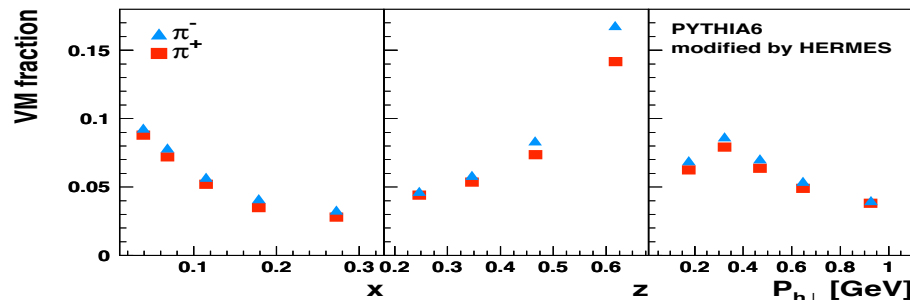
Collins Moments for $\pi^+ \pi^-$ from 2002–2004 H \uparrow Data



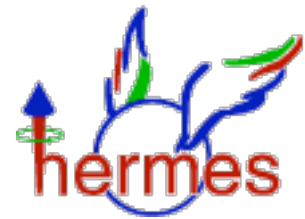
It exists!



- First evidence for **non-zero Collins function ... and transversity!**
- **Positive** for π^+ ...
Negative and ***larger*** for π^- ...
- Systematic error bands include acceptance and smearing effects, and contributions from unpolarized $\langle \cos(2\phi) \rangle$ and $\langle \cos(\phi) \rangle$ moments

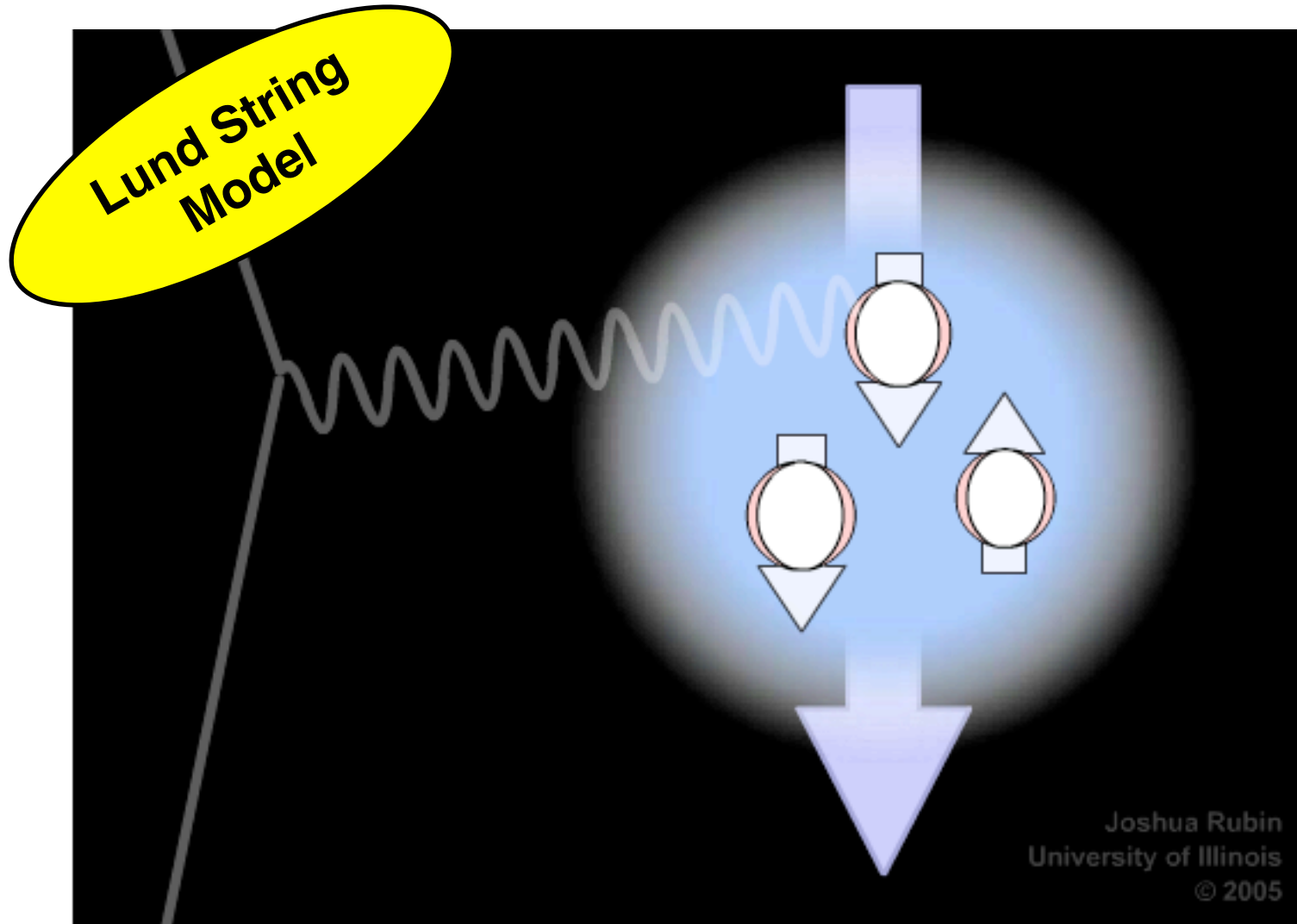


Understanding the Collins Effect



The Collins function exists! → **spin-orbit** correlations in π formation

Is the Artru mechanism responsible?



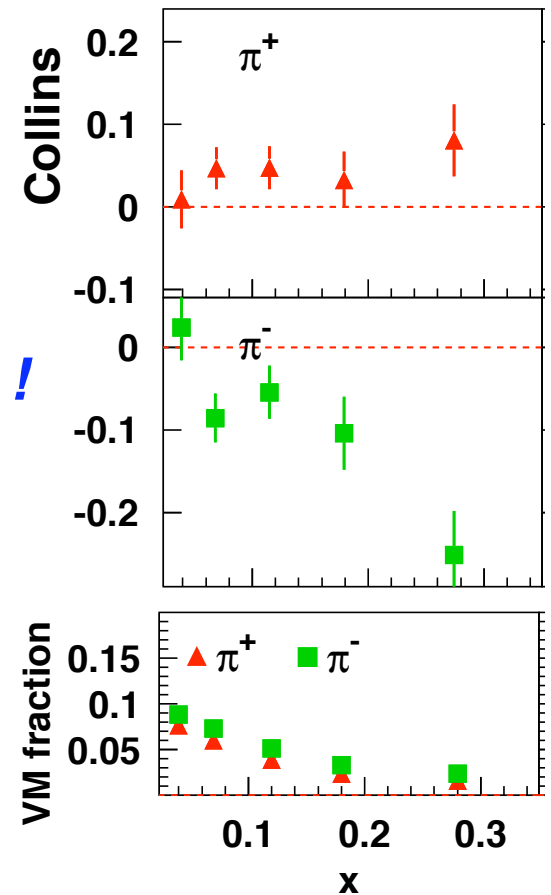
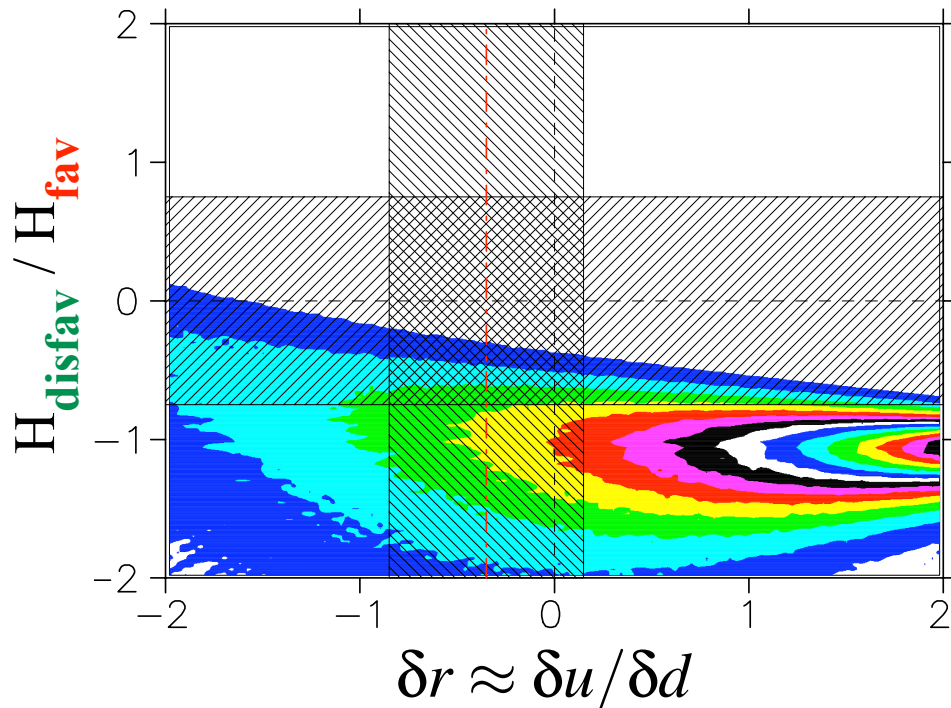
Why are the Collins π^- asymmetries so large?



DIS on proton target always dominated by *u-quark scattering*

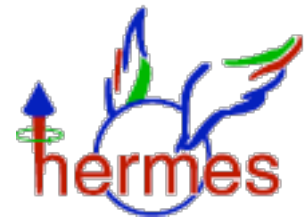
- $A_{\text{Col}}^{\pi^+} \sim \delta u H_{1,\text{fav}}^\perp$... expect: **positive**
- $A_{\text{Col}}^{\pi^-} \sim \delta u H_{1,\text{disfav}}^\perp$... expect: **~ zero**

Data indicate *disfavored* CollinsFF is **large & negative!**



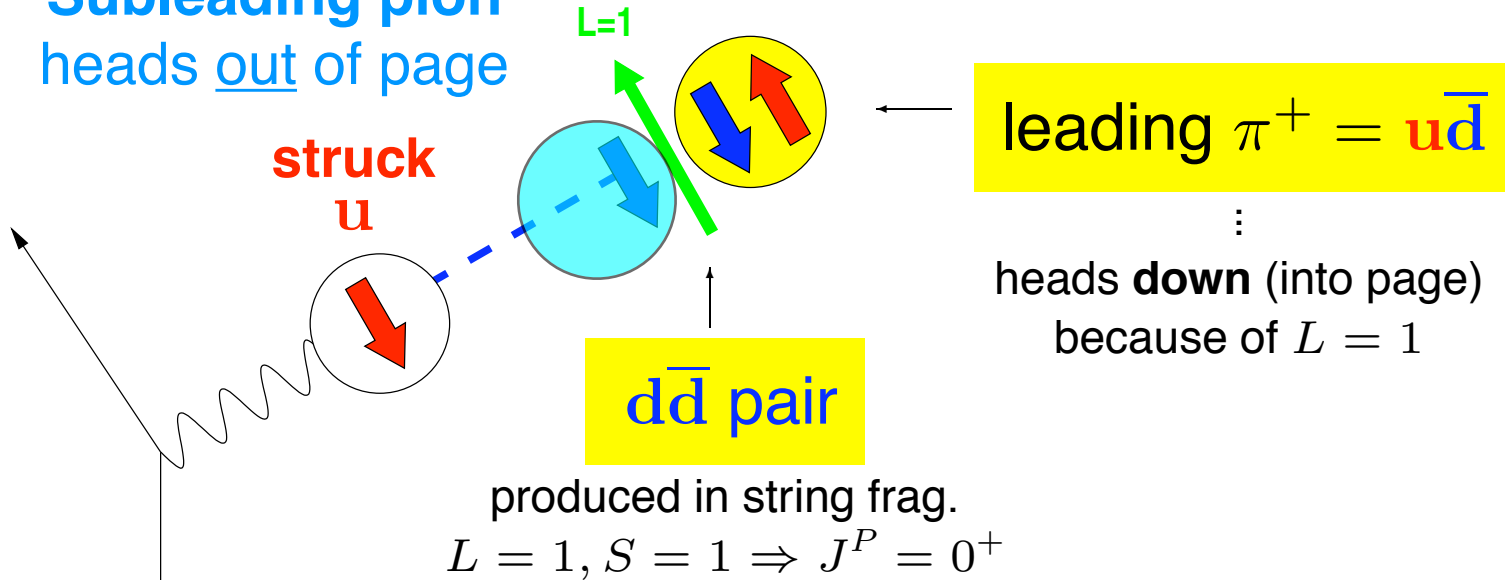
Map out solution space ...
find $H_{\text{disfav}} \approx -H_{\text{fav}}$

Interpretation of Collins Results



Lund model + 3P_0 hypothesis once more:

Subleading pion
heads out of page



➡ leading π^+ = **favored** transition, heads into page

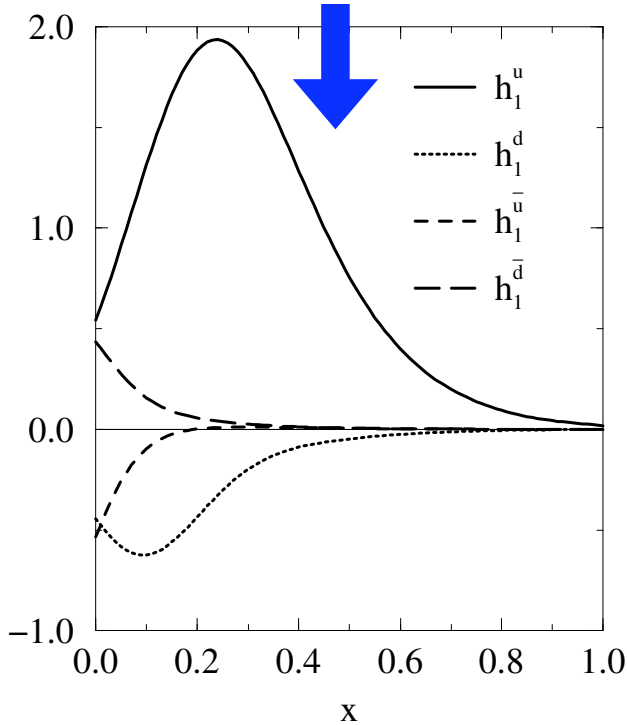
➡ subleading particle (prob π^-) = **disfavored** transition, heads out of page

Perhaps $H_{\text{dis}} \approx -H_{\text{fav}}$ is not only reasonable, but likely ?

Collins Global Fit: HERMES (H target) & COMPASS (D target)

Efremov, Goeke, Schweitzer, hep-ph/0603054

Take $h_1(x)$ from Chiral-Quark Soliton Model:



Fit K_T -integrated favored and unfavored Collins FF to HERMES data:

$$H_1^{\text{fav}} = H_1^{u/\pi^+} = H_1^{d/\pi^-} = \dots$$

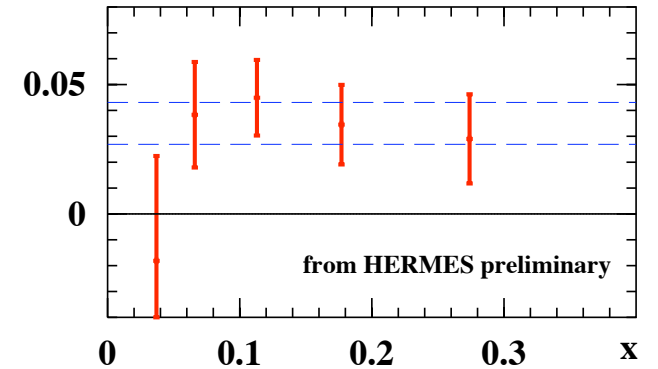
$$H_1^{\text{unf}} = H_1^{u/\pi^-} = H_1^{d/\pi^+} = \dots$$

$$B_{\text{Gauss}}(z) \equiv 1/\sqrt{1+z^2\langle \mathbf{p}_{h_1}^2 \rangle / \langle \mathbf{K}_{H_1}^2 \rangle}$$

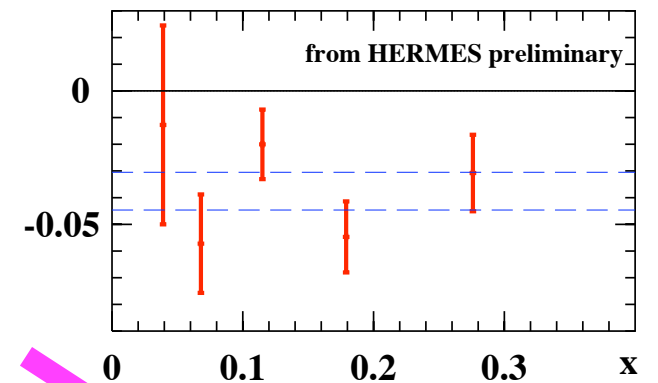
Also find $H_1^{\text{fav}} \approx -H_1^{\text{unf}}$

Gives good fit to COMPASS!

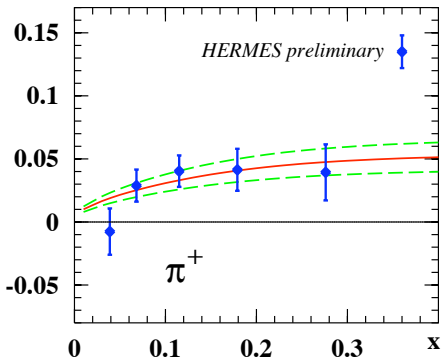
$2\langle B_{\text{Gauss}} H_1^{\perp (1/2) \text{fav}} \rangle$ (a)



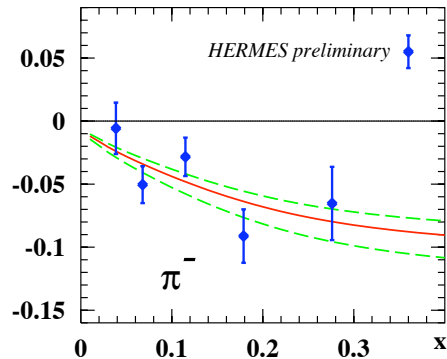
$2\langle B_{\text{Gauss}} H_1^{\perp (1/2) \text{unf}} \rangle$ (b)



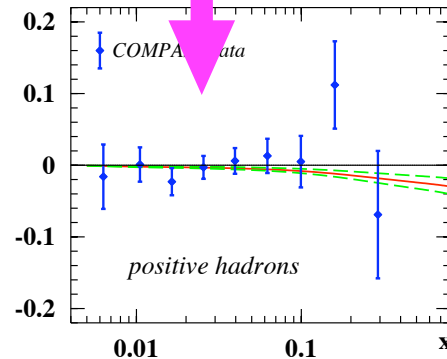
$A_{\text{UT}}^{\sin(\phi+\phi_S)}(x)$ for proton (a)



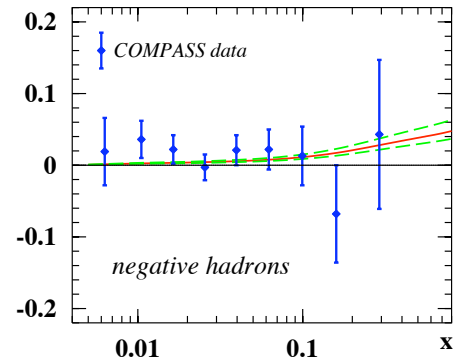
$A_{\text{UT}}^{\sin(\phi+\phi_S)}(x)$ for proton (b)



$A_{\text{UT}}^{\sin\phi_C}(x)$ for deuteron (c)



$A_{\text{UT}}^{\sin\phi_C}(x)$ for deuteron (d)



Collins Global Fit: HERMES (ep) & BELLE (e⁺e⁻)

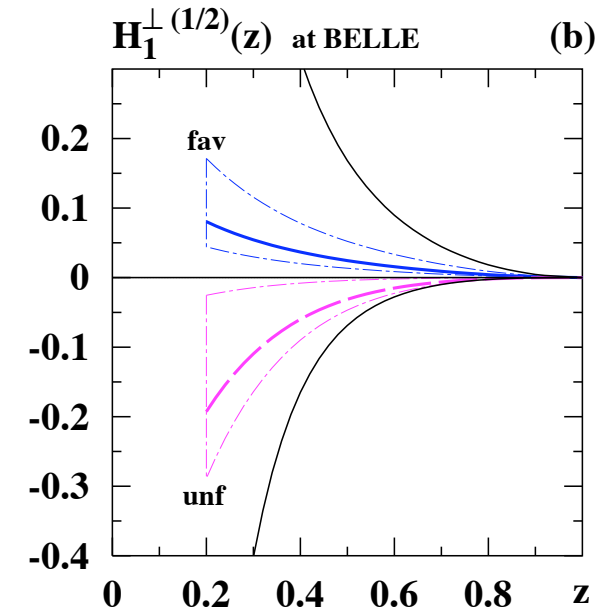
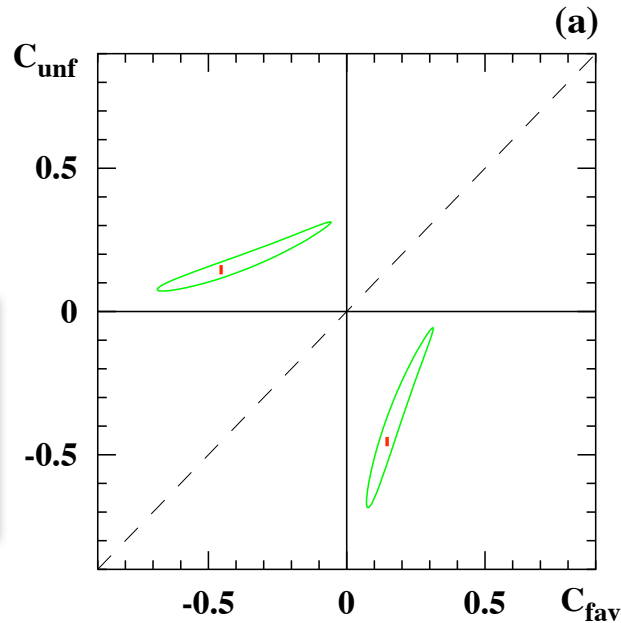
Efremov, Goeke, Schweitzer, hep-ph/0603054

Fit **BELLE** z-dependent results to

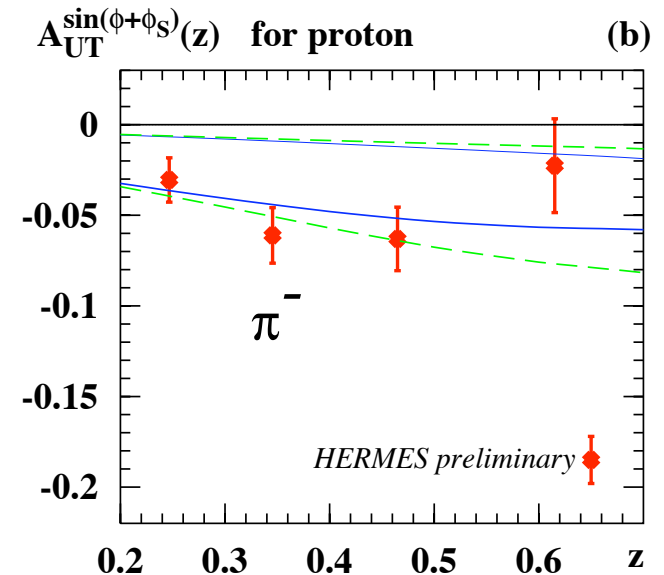
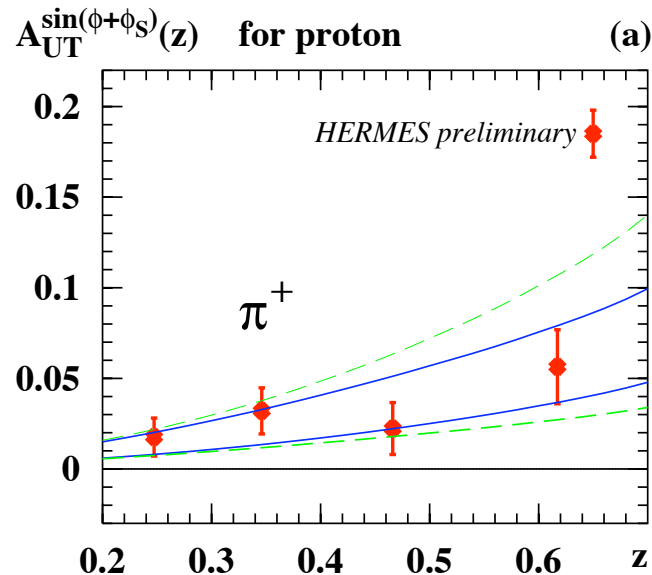
$$H_1^{\perp(1/2)a}(z) = C_a z D_1^a(z)$$

$$C_{\text{fav}} = 0.15, C_{\text{unf}} = -0.45$$

$$\text{and so } H_1^{\text{fav}} \approx -H_1^{\text{unf}}$$

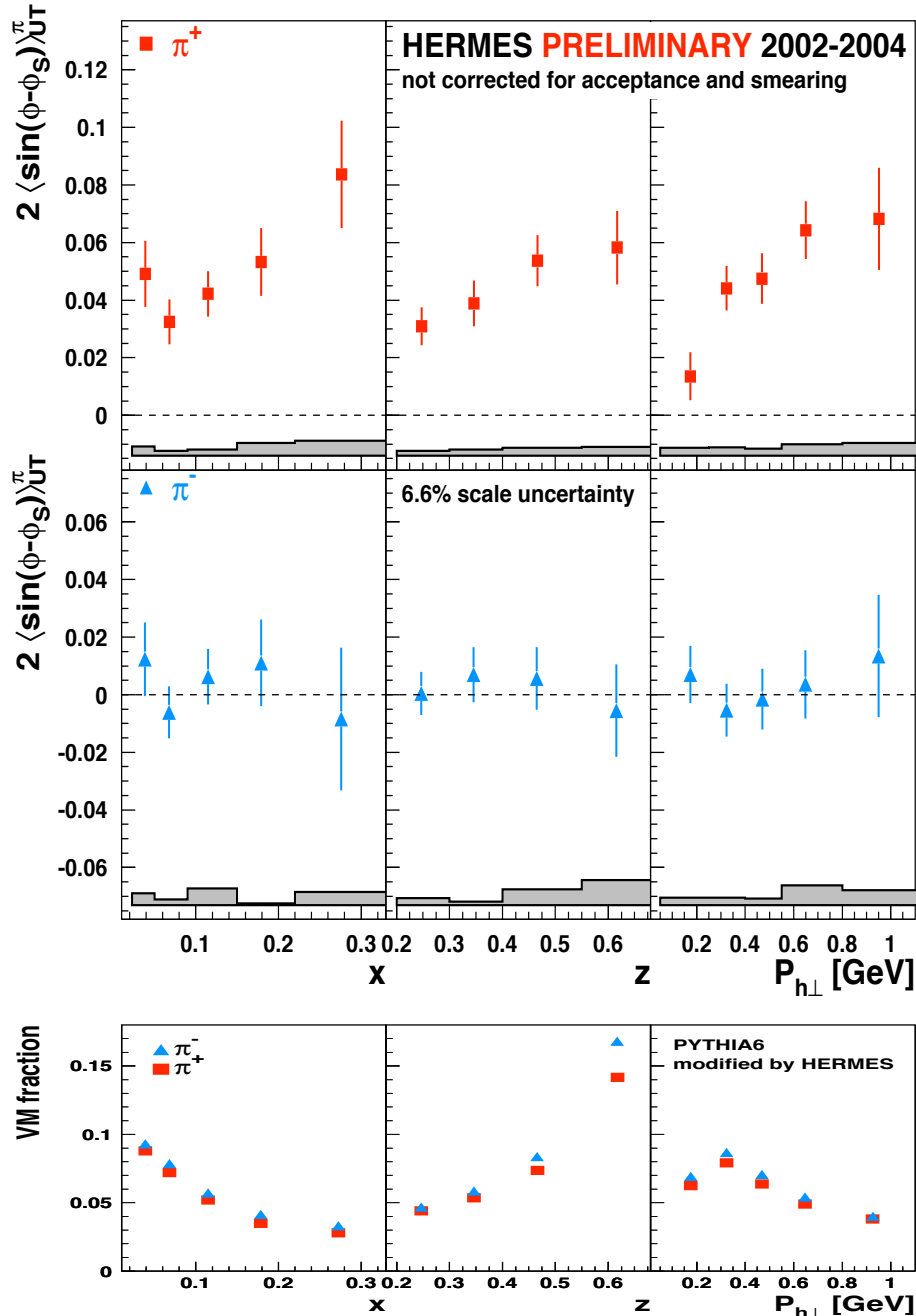


Resulting Collins FF also fit **HERMES** data well



*SSA Results 2:
Sivers Effect*

Sivers Moments for $\pi^+ \pi^-$ from 2002–2004 H \uparrow Data



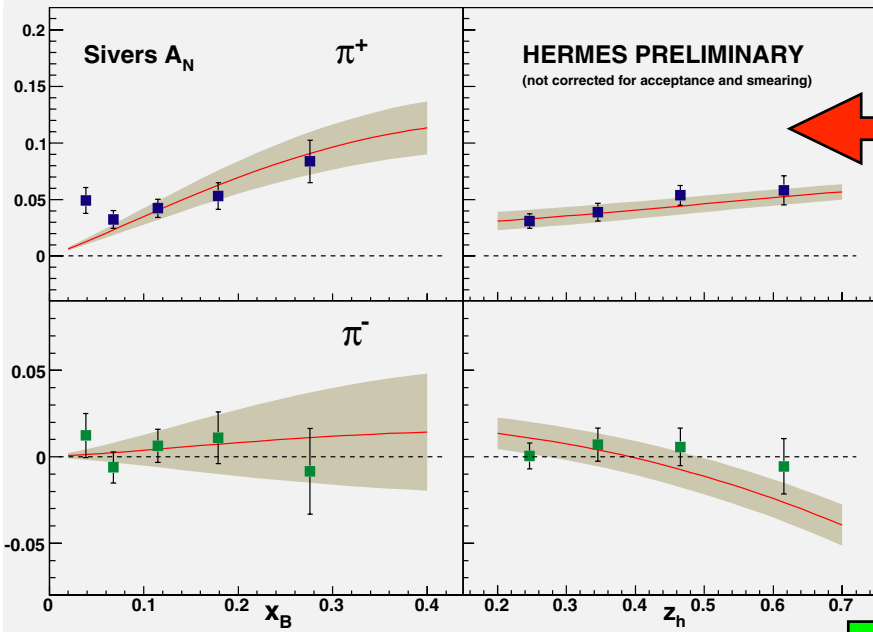
It exists too!



- First evidence for non-zero Sivers function!
- \Rightarrow presence of non-zero **quark orbital angular momentum!**
- **Positive** for π^+ ...
Consistent with zero for π^- ...
- Systematic error bands include acceptance and smearing effects, and contributions from unpolarized $\langle \cos(2\phi) \rangle$ and $\langle \cos(\phi) \rangle$ moments

Sivers Global Fit: HERMES & COMPASS

Vogelsang & Yuan,
PRD 72 (2005) 054028



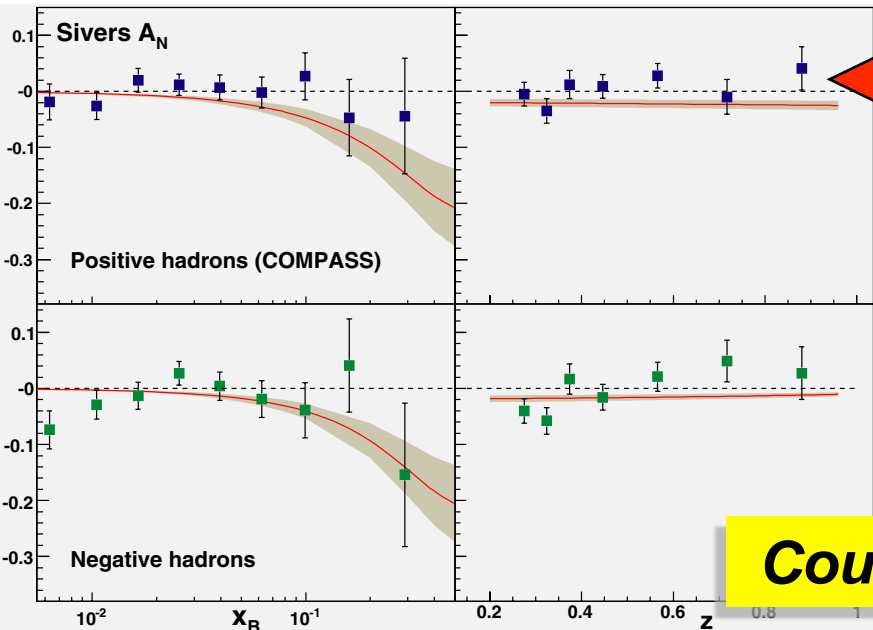
For convenience: $q_T(x) \equiv f_{1T}^{\perp,q}(x)$

Fit HERMES A_{UT} to Sivers funcⁿ of form:

$$\frac{u_T^{(1/2)}(x)}{u(x)} = S_u x(1-x), \quad \frac{d_T^{(1/2)}(x)}{u(x)} = S_d x(1-x)$$

- assume no antiquark Sivers func: $\bar{q}_T(x) = 0$
- unpol PDFs = GRV-LO, unpol FFs = Kretzer

$$S_u = -0.81 \pm 0.07, \quad S_d = 1.86 \pm 0.28$$



Fits COMPASS deuterium data well!

But a surprise! $S_d \gg S_u$!

e.g., large- N_C expectation: $u_T(x) \approx -d_T(x)$

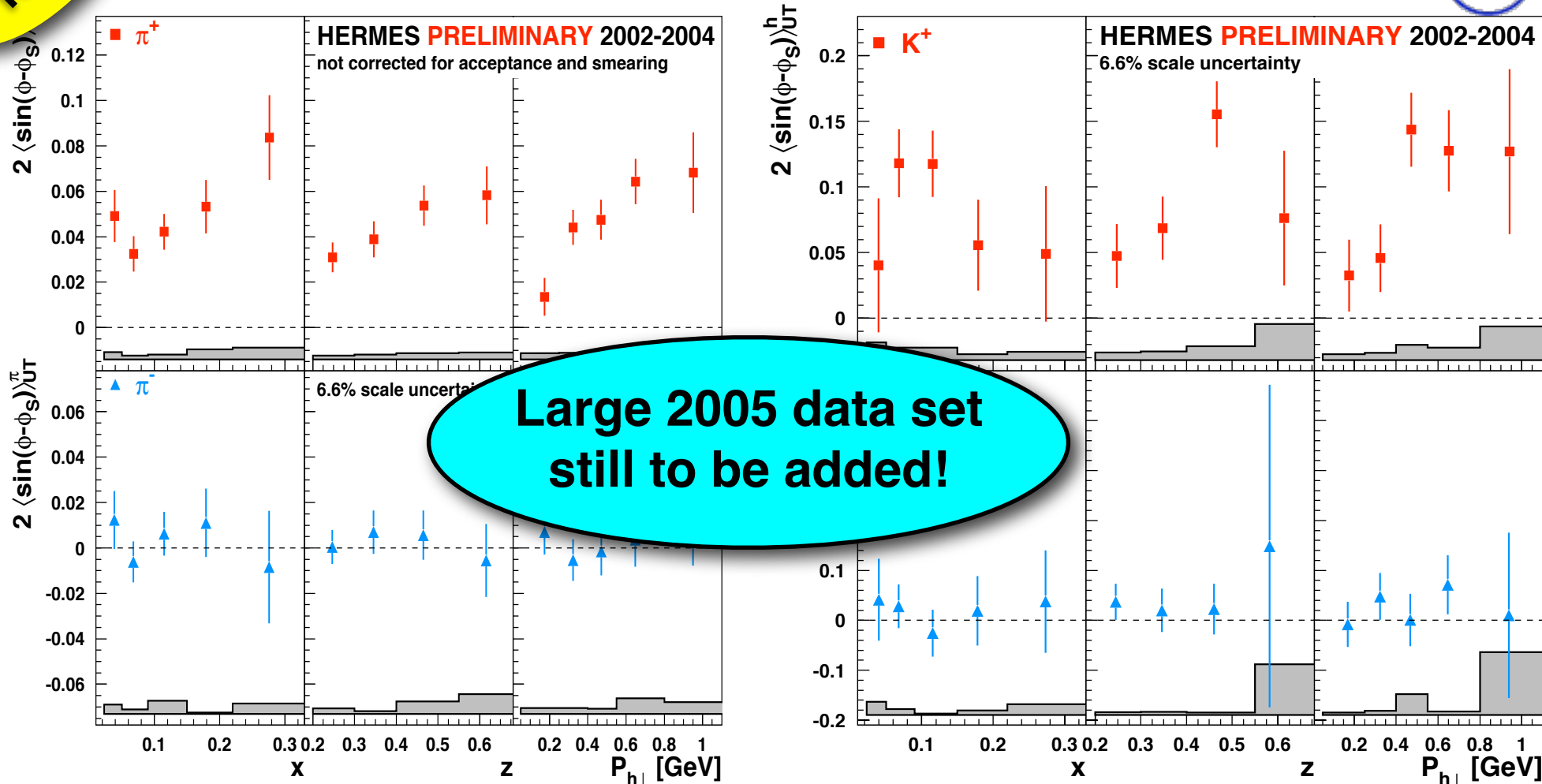
Hmm ... S_u actually reflects $u_T + \bar{d}_T/4$

... S_d actually reflects $d_T + 4\bar{u}_T$

Could Sivers (and L) be large for antiquarks?

NEW!

Sivers Moments for Kaons from 2002–2004 Data



Large 2005 data set still to be added!

- Effect about **equal** for $K^- = s\bar{u}$ and $\pi^- = d\bar{u}$ → note: same antiquark ...
- + Effect seems larger for $K^+ = u\bar{s}$ than $\pi^+ = u\bar{d}$ at $x \approx 0.1$... !

→ significant **antiquark** Sivers functions? and strongly flavor-dependent?

Conclusions

● Quark and gluon polarization

- ➔ **quark polarization** is **positive**, but much lower than CQM / bag model expectations
- ➔ **anti-quark** polarization consistent with zero within measured range, including improved verification of $\Delta s \approx 0$
- ➔ data coming in from COMPASS and RHIC-Spin on ΔG ... so far a **modest**, positive value favoured ...

● Collins fragmentation function

- ➔ **opposite sign** and similar magnitude to **favoured** function
sign of effect supports **3P_0 picture** of color string breaking
- ➔ result now **confirmed** by new data from **BELLE**,
+ successful global analyses including **COMPASS** data

● **Sivers** effect is **non-zero** in DIS!

- ➔ successful global analysis of **HERMES** (H) & **COMPASS** (D)
- ➔ ... and suggests large **antiquark** contributions to **orbital L**
- ➔ latest HERMES data on **Kaon producⁿ** seem to support this ...