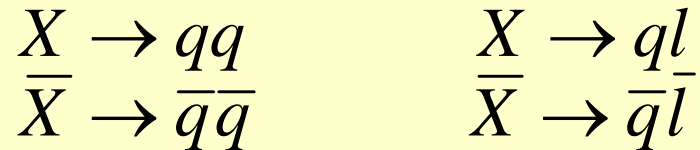


Baryogenesis

- **Plausibility Argument** (GUTS Baryogenesis)
 - Consider very heavy boson X ($M_X \sim 10^{19}$ GeV)

- Baryon number violation:



- C & CP violation

$$\begin{array}{ll} \Gamma_{X \rightarrow qq} = (1 + \Delta_q)\Gamma_q; & \Gamma_{X \rightarrow ql} = (1 - \Delta_l)\Gamma_l \\ \Gamma_{\bar{X} \rightarrow \bar{q}\bar{q}} = (1 - \Delta_q)\Gamma_q; & \Gamma_{\bar{X} \rightarrow \bar{q}\bar{l}} = (1 + \Delta_l)\Gamma_l \end{array}$$

but $\Gamma_X^{Tot} = \Gamma_{\bar{X}}^{Tot}$ (**CPT conservation**)

- Out of Thermal Equilibrium

If in Equilibrium then the reverse reactions

(e.g. $qq \rightarrow X$, $\bar{q}\bar{q} \rightarrow \bar{X}$) will smooth out any matter/antimatter excess

Origin of EDMs

- Standard Model EDMs are due to CP violation in the quark weak mixing matrix
 - CKM (e.g. the K^0/B^0 -system) but...
 - e^- and quark EDM's are zero at first order
 - Need at least two "loops" to get EDM's (electron actually requires 4 loops!)
 - Thus EDM's are VERY small in standard model

**Neutron EDM in Standard Model is
 $\sim 10^{-32}$ e-cm ($=10^{-19}$ e-fm)**

**Electron EDM in Standard Model is
 $< 10^{-40}$ e-cm**

Origin of EDMs

- Quark EDMs lead to hadronic EDMs (neutrons or diamagnetic atoms)



All e^- paired-up

- Electron EDMs lead to atomic EDMs in paramagnetic atoms



One unpaired e^-

- Nuclear searches are in hadrons

Origin of Hadronic EDMs

- Hadronic (strongly interacting particles) EDMs are from
 - θ_{QCD} (a special parameter in Quantum Chromodynamics - QCD)
 - or from the quarks themselves

EDM from θ_{QCD}

- θ_{QCD} results from CP-odd term is \mathcal{L}_{QCD}

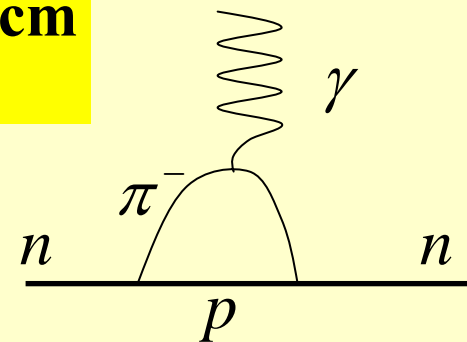
$$\mathcal{L}_{\text{QCD}} = -\theta \left(\frac{\alpha_s}{8\pi} \right) \tilde{G}_a^{\mu\nu} G_{\mu\nu}^a$$

- θ_{QCD} should be naturally about ~ 1

- This gives an "effective" neutron EDM of

$$d_n = \frac{g_{\pi NN}}{4\pi^2} \left(\frac{e}{m_p f_\pi} \right) \ln \left(\frac{m_\rho}{m_\pi} \right) \left(\frac{m_u m_d}{m_u + m_d} \right) \theta \approx (-10^{-15}) \theta \text{ e-cm}$$

but $d_n^{\text{exp}} < 10^{-25} \text{ e-cm}$
 $\therefore \theta < 10^{-10}$
 Why so small??

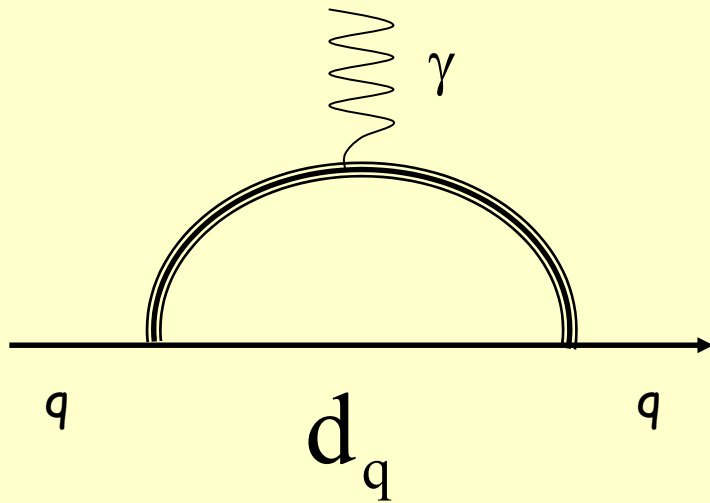


EDM from θ_{QCD}

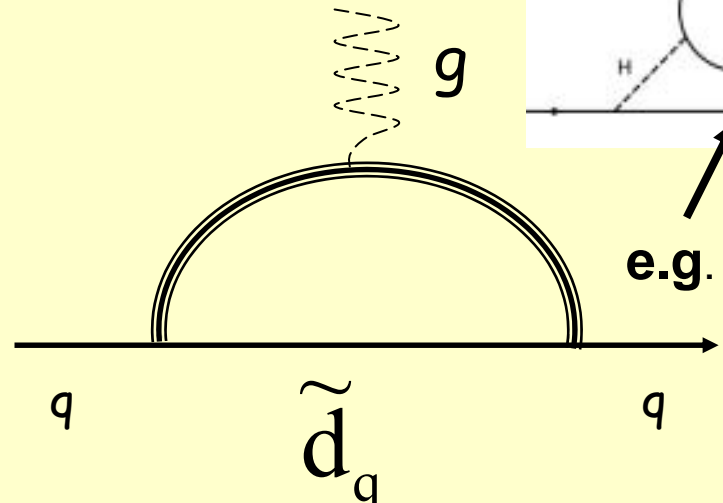
- This is the Strong CP problem in QCD
- Small θ_{QCD} does not provide any new symmetry for \mathcal{L}_{QCD}
 - Popular solution is "axions" (Peccei-Quinn symmetry) - new term in \mathcal{L}_{QCD}
 - No Axions observed yet
 - Extra dimensions might suppress θ_{QCD}
(Harnik et al arXiv:hep-ph/0411132)
 - Remains an unsolved theoretical "problem"

Hadronic EDM from Quarks

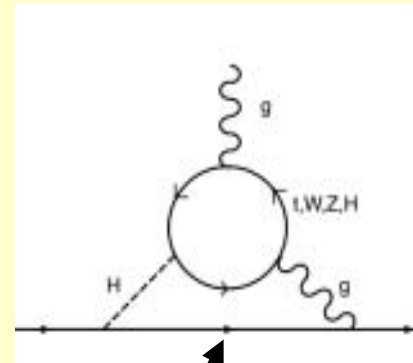
- Quark EDM contributes via



Quark EDM



Quark ChromoEDM



e.g.

but $d_n^{\text{Standard Model}} < 10^{-31} \text{ e-cm}$
 recall $d_n^{\text{exp}} < 10^{-25} \text{ e-cm}$

Note: $d_n \cong 0.7(d_d - \frac{1}{4}d_u) + 0.6(\tilde{d}_d + \frac{1}{2}\tilde{d}_u)$
 $d_{199\text{Hg}} \cong 0.007(d_d - \tilde{d}_u)$

Physics Beyond the Standard Model

- New physics (e.g. SuperSymmetry = SUSY) has additional CP violating phases in added couplings

- New phases: (ϕ_{CP}) should be ~ 1 (why not?)

- Contributions to EDMs depends on masses of new particles

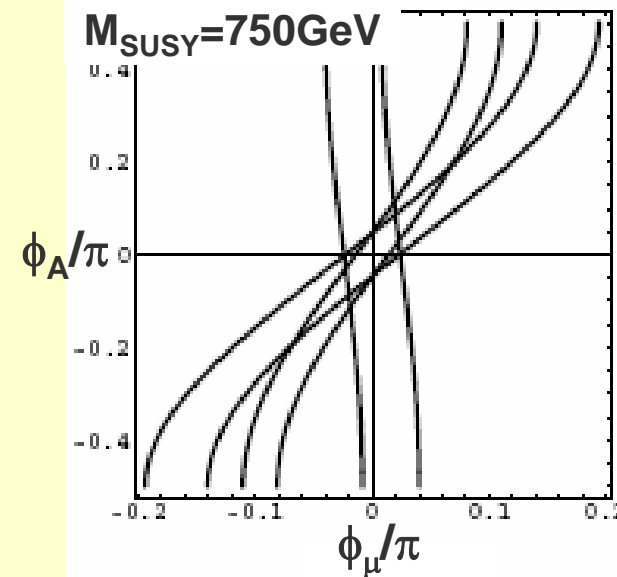
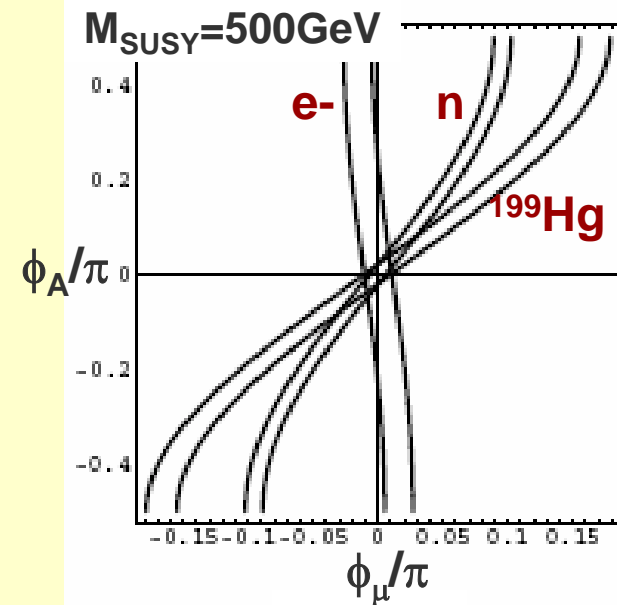
$$d_n \propto \frac{\sin \phi_{CP}}{M_{SUSY}^2}$$

- In MSSM (Minimal Supersymmetric Standard Model)

$$d_n \sim \left(\frac{200 \text{ GeV}}{M_{SUSY}} \right)^2 \times 10^{-25} \text{ e-cm}$$

Possible impacts of non-zero EDMs

- Must be new Physics (at proposed sensitivities)
- Sharply constrains models beyond the Standard Model (especially *with* LHC data)
Large Hadron Collider
- May account for matter-antimatter asymmetry of the universe (via ElectroWeak Baryogenesis)



From Pospelov et al for CMSSM

EDM Measurements

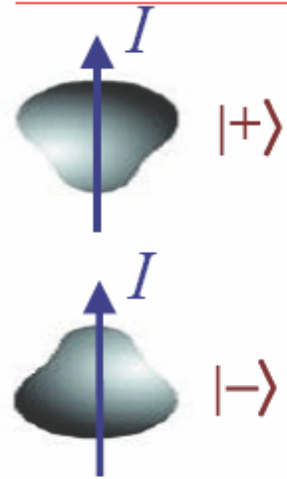
particle	Present Limit (90% CL) (e-cm)	Laboratory	Possible Sensitivity (e-cm)	Standard Model (e-cm)
e ⁻ (TI)	1.6×10^{-27}	Berkeley	10^{-29}	$<10^{-40}$
e ⁻ (PbO)		Yale		
e ⁻ (YbF)		Sussex		
e ⁻ (GGG)		LANL/Indiana		
μ	9.3×10^{-19}	CERN	$<10^{-24}$	$<10^{-36}$
μ		BNL		
n	6.3×10^{-26}	ILL	1.5×10^{-26}	$\sim 10^{-32}$
n		ILL	$\sim 2 \times 10^{-28}$	
n		PSI	$\sim 7 \times 10^{-28}$	
n		SNS	$< 1 \times 10^{-28}$	
¹⁹⁹ Hg	1.9×10^{-27}	Seattle	5×10^{-28}	$\sim 10^{-33}$
¹²⁹ Xe		Princeton	10^{-31}	$\sim 10^{-34}$
²²⁵ Ra		Argonne	10^{-28}	
²²³ Rn		TRIUMF	1×10^{-28}	
d		COSY/JPARC?	$<10^{-27}$	

New Nuclear EDM Techniques

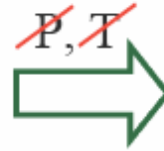
- Accelerator production of radioactive diamagnetic atoms
(probes hadronic EDMs)
- Charged particles in storage rings
- Superthermal sources of Ultra-Cold Neutrons

Enhanced Atomic EDM via Octupole deformations

Octupole deformations



$$\begin{aligned} \Psi^+ &= (|+\rangle + |-\rangle)/\sqrt{2} \\ \Psi^- &= (|+\rangle - |-\rangle)/\sqrt{2} \end{aligned}$$



$$\Psi^+ = ((1+\alpha)|+\rangle + (1-\alpha)|-\rangle)/\sqrt{2}$$

$$\Psi^- = ((1-\alpha)|+\rangle + (1+\alpha)|-\rangle)/\sqrt{2}$$

$$\alpha = \frac{\langle \Psi^- | V^{PT} | \Psi^+ \rangle}{\Delta E} \sim \frac{\beta_3 A^{-1/3}}{\Delta E}$$

$$S_{\text{intr}} \sim eZA\beta_2\beta_3$$

$$S_{\text{lab}} \sim eZA^{2/3}\beta_2\beta_3^2/\Delta E$$

$$\beta_2, \beta_3 \sim 0.1$$

Haxton & Henley; Auerbach, Flambaum & Spevak; Hayes, Friar & Engel

	^{223}Rn	^{223}Ra	^{225}Ra	^{225}Fr	^{225}Ac	^{229}Pa	^{199}Hg	^{129}Xe
$t_{1/2}$	23.2 m	11.4 d	14.9 d	22 m	10.0 d	1.5 d		
I	7/2	3/2	1/2	3/2	3/2	5/2	1/2	1/2
Δe_{th} (keV)	37	170	47	75	49	5		
ΔE_{exp} (keV)	--	50.2	55.2	160.5	40.1	0.22		
$10^5 S$ (efm ³)	1000	400	300	500	900	12000	-1.4	1.75
$10^{28} d_A$ (e cm)	2000	2700	2100	2800			-5.6	0.8

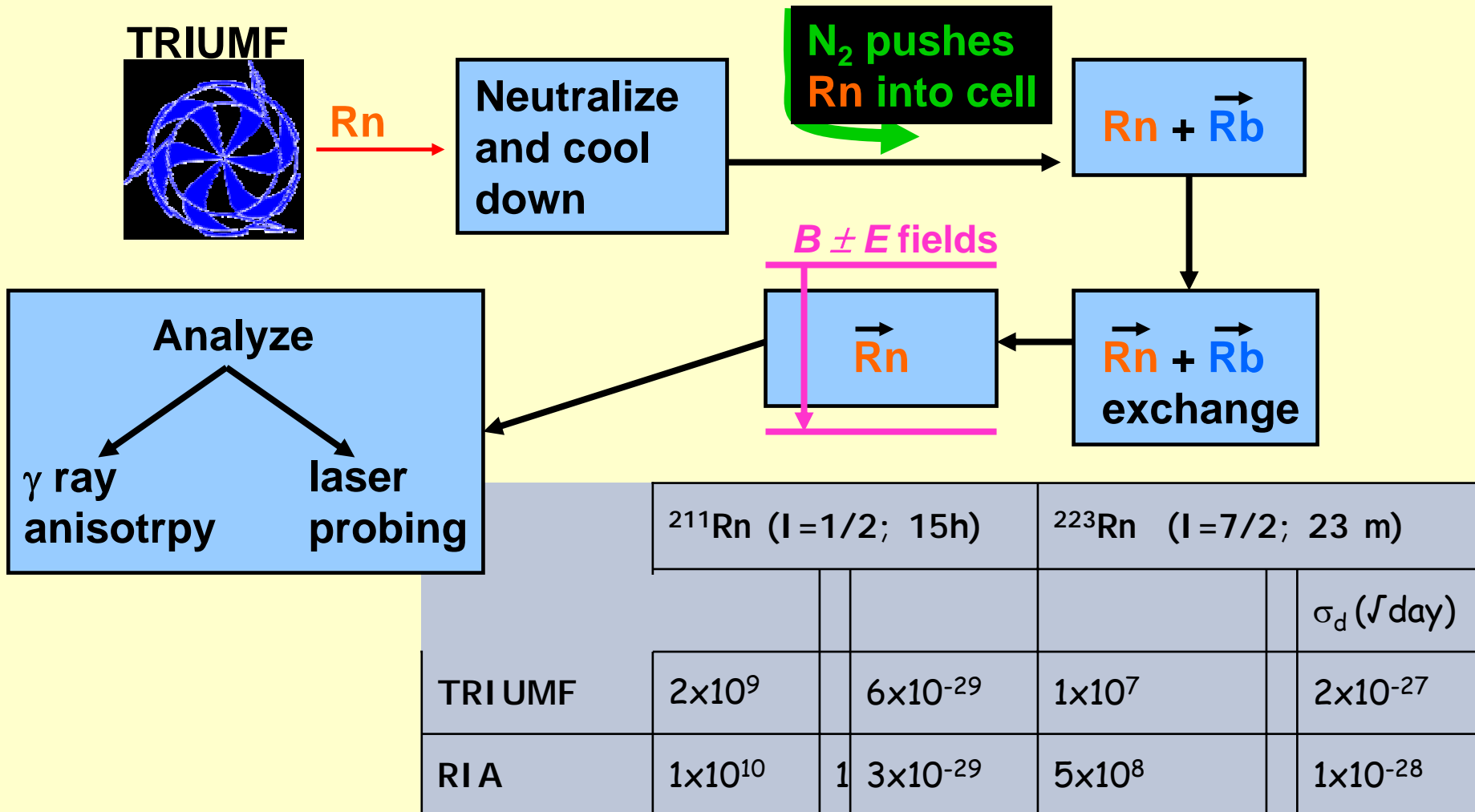
EDM in Rn

Spokesmen: Timothy Chupp² and Carl Svensson¹

Sarah Nuss-Warren², Eric Tardiff², Kevin Coulter², Wolfgang Lorenzon², Timothy Chupp²

John Behr⁴, Matt Pearson⁴, Peter Jackson⁴, Mike Hayden³, Carl Svensson¹

University of Guelph¹, University of Michigan², Simon Fraser University³, TRIUMF⁴



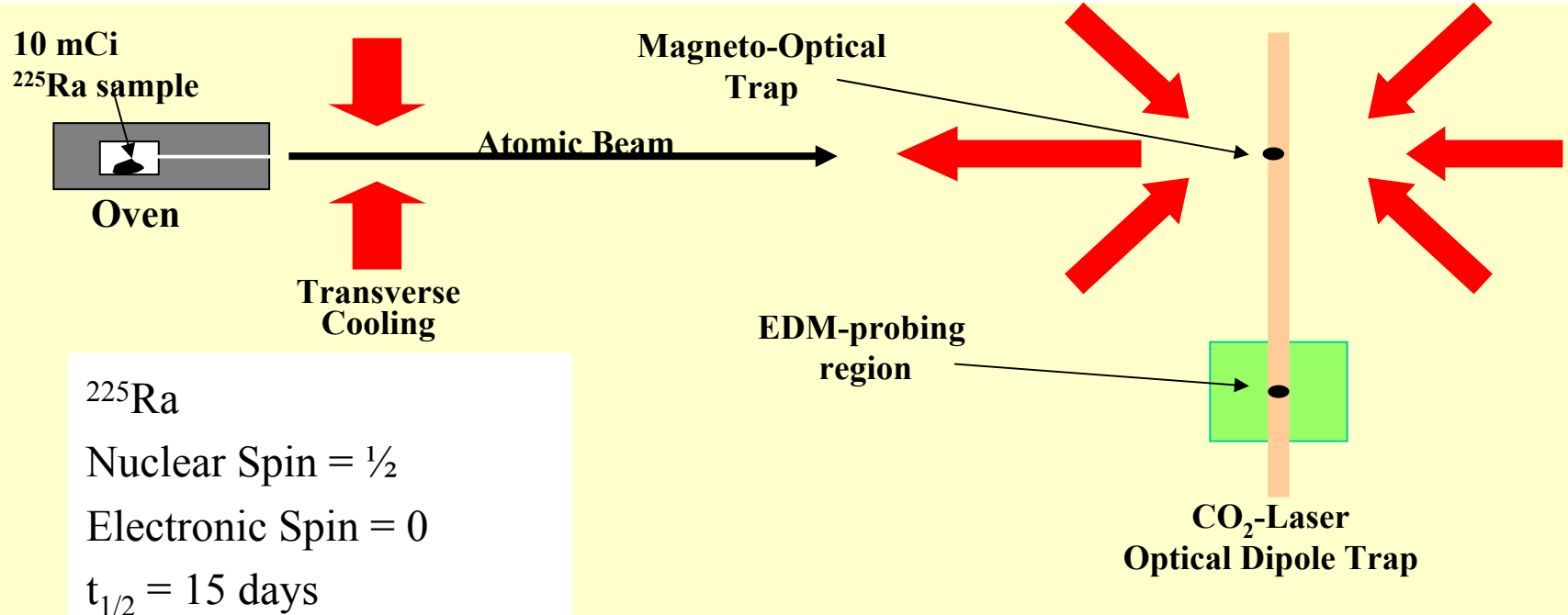
	²¹¹ Rn (I = 1/2; 15h)		²²³ Rn (I = 7/2; 23 m)		
					σ _d (J day)
TRIUMF	2×10 ⁹		6×10 ⁻²⁹	1×10 ⁷	2×10 ⁻²⁷
RIA	1×10 ¹⁰	1	3×10 ⁻²⁹	5×10 ⁸	1×10 ⁻²⁸

EDM with Trapped Radium Atoms

Irshad Ahmad, Roy J. Holt, Zheng-Tian Lu, Elaine C. Schulte, Physics Division, Argonne National Laboratory

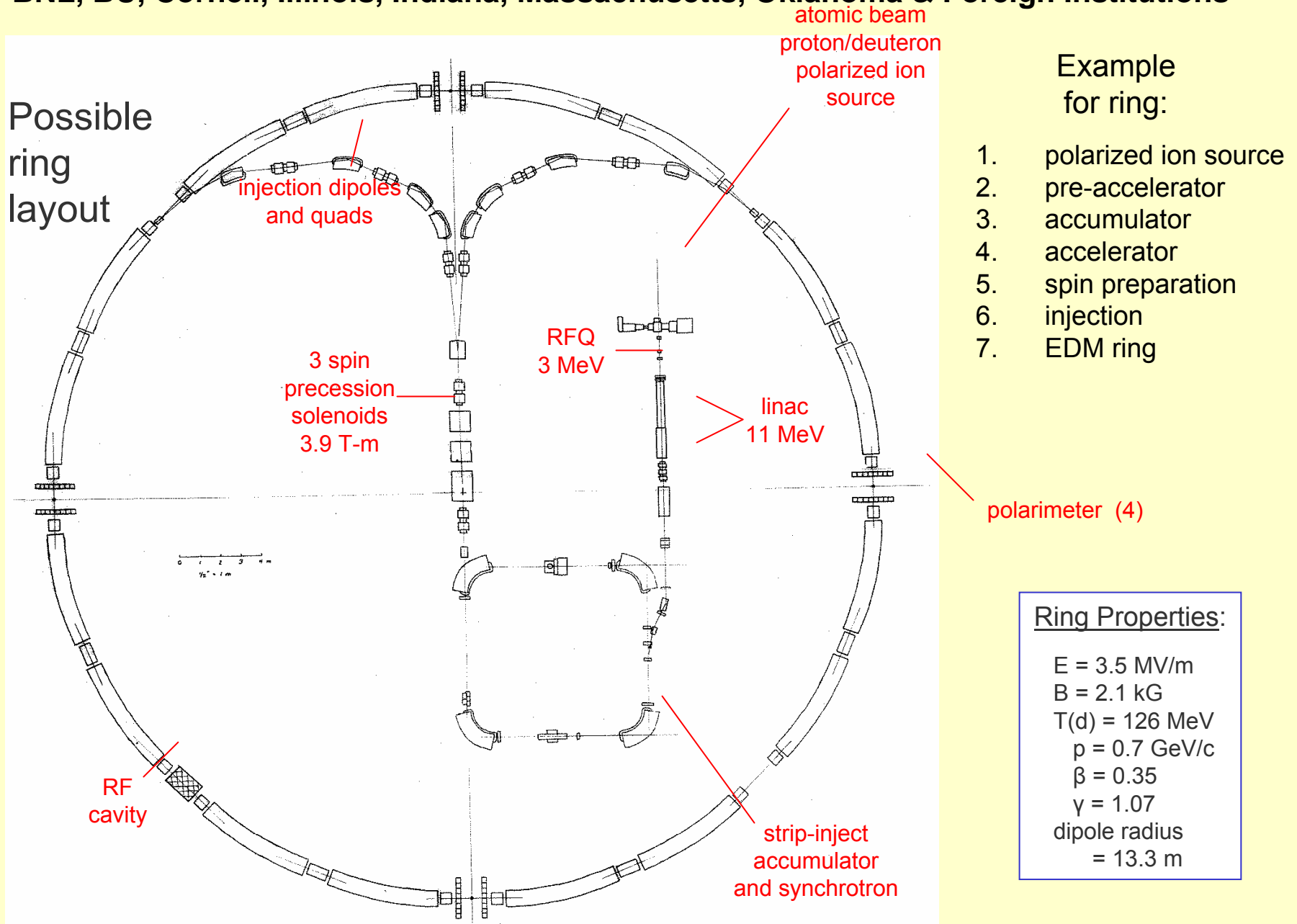
Advantages of an EDM measurement on ^{225}Ra atoms in a trap

- In ^{225}Ra the EDM effect is enhanced by two orders of magnitude due to nuclear quadrupole and octupole deformation.
- Trap allows a long coherence time (~ 300 s).
- Cold atoms result in a negligible “ $v \times E$ ” systematic effect.
- Trap allows the efficient use of the rare and radioactive ^{225}Ra atoms.
- Small sample in an UHV allows a high electric field (> 100 kV/cm).



Deuteron (and Muon) EDM in Storage Ring

BNL, BU, Cornell, Illinois, Indiana, Massachusetts, Oklahoma & Foreign Institutions

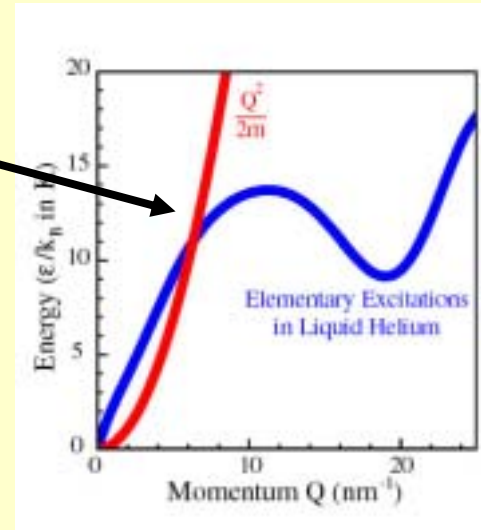


New Techniques for n-EDM:

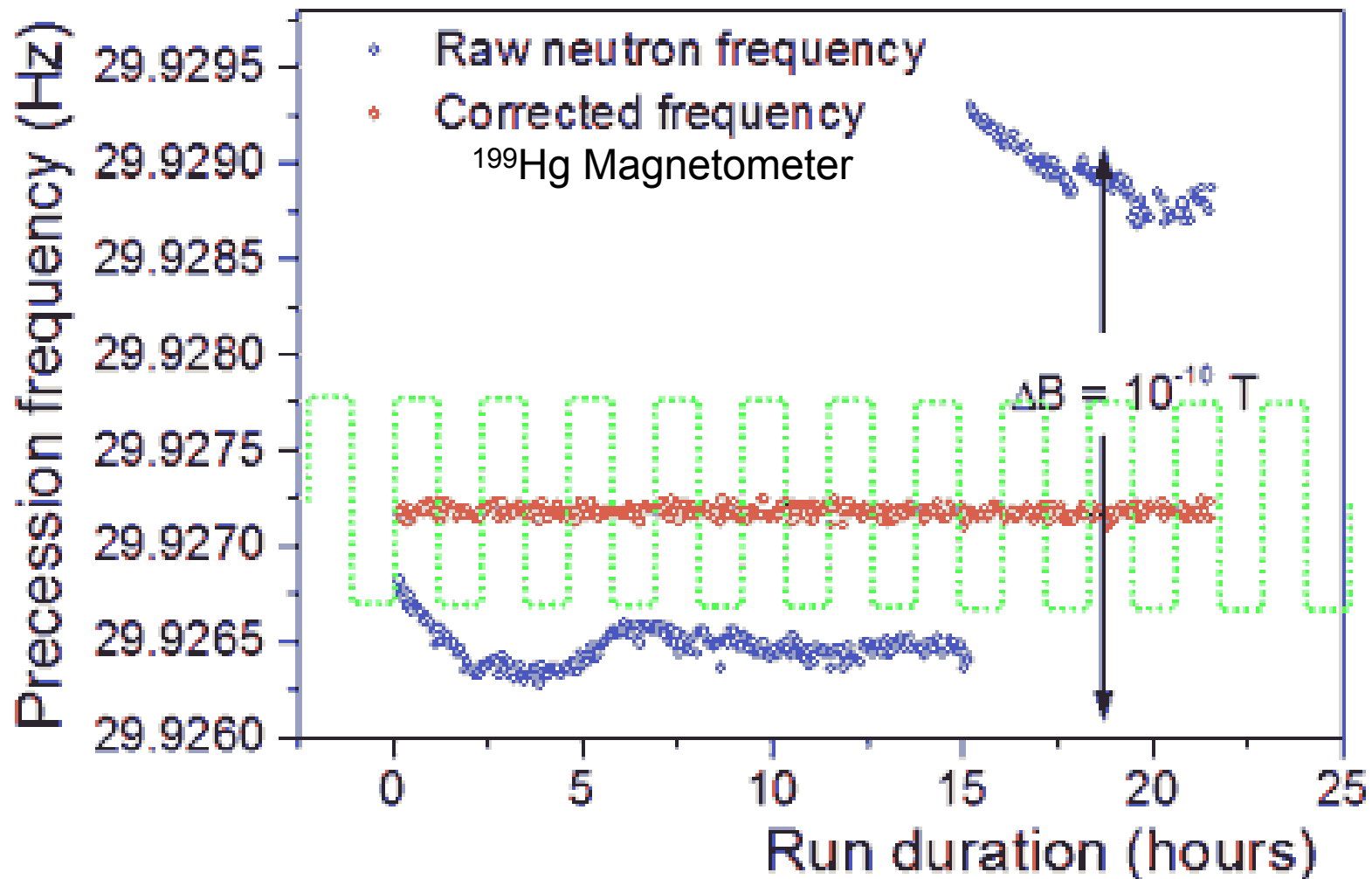
- Use Superthermal (non-equilibrium) system to produce UCN
 - Superfluid ^4He can yield ~ 1000 more UCN than conventional UCN source
- Higher Electric fields in ^4He
 - Breakdown voltage may be 10x vacuum breakdown
- ^3He comagnetometer measures B-field at same location as neutrons
 - *Very* small amount of ^3He in ^4He
 - Use SQUIDS to measure ^3He precession - calibrates B-field since

$$\omega_3 \propto |\vec{B}|$$
- $\vec{n} + ^3\vec{\text{He}} \Rightarrow t + p$ has $\sigma_{\uparrow\downarrow} \gg \sigma_{\uparrow\uparrow}$
 - Detect capture via scintillation of ^4He
 - Same technique as NIST LHe τ_n measurement (UV photons converted to visible in tetraphenyl butadiene - TPB)
 - Measures difference of ω_n and ω_3
- "Dressed" spin technique suppresses sensitivity to fluctuations in B-field

8.9\AA incident n produces roton & becomes UCN

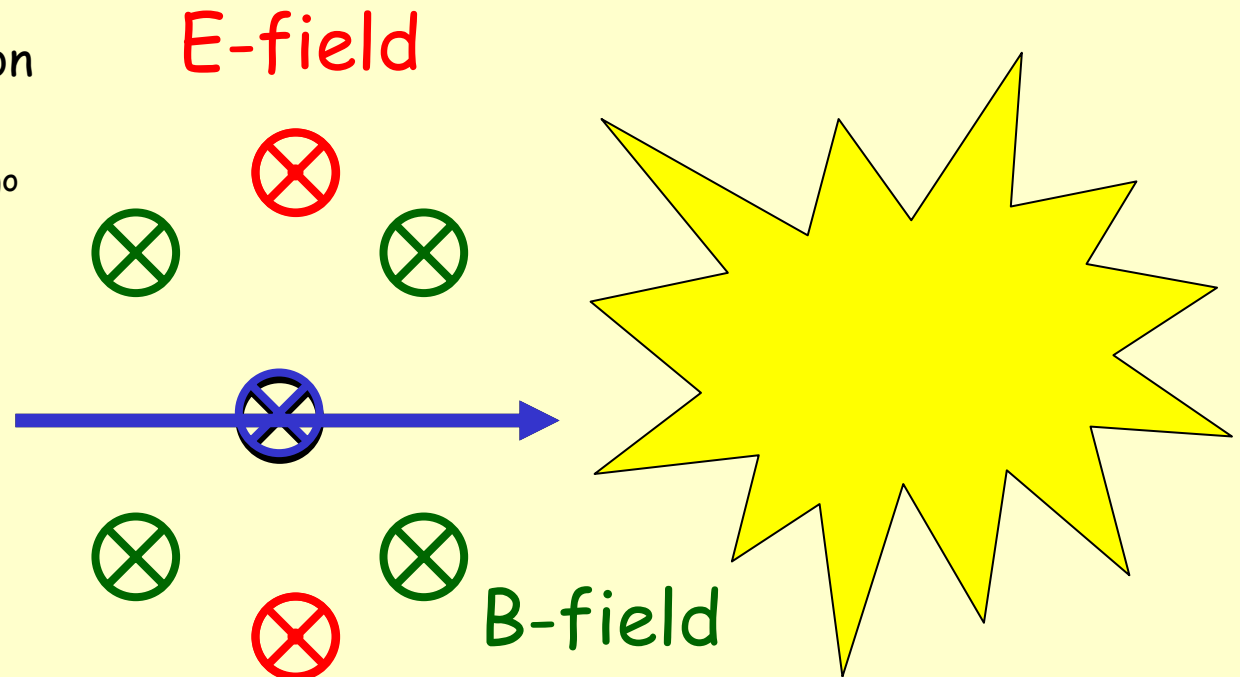


Careful magnetometry is essential !



New Technique for n-EDM

1. Inject polarized neutron & polarized ^3He
2. Rotate both spins by 90°
3. Measure $n+^3\text{He}$ capture vs. time
(note: $\sigma_{\downarrow\uparrow} \gg \sigma_{\uparrow\uparrow}$)
4. Flip E-field direction



^3He functions as "co-magnetometer"

EDM Statistical Sensitivity

	EDM @ ILL	EDM @ SNS
N_{UCN}	1.3×10^4	2×10^6
$ \vec{E} $	10 kV/cm	50 kV/cm
T_m	130 s	500 s
m (cycles/day)	270	50
σ_d (e-cm)/day	3×10^{-25}	3×10^{-27}
SNR (signal noise ratio)	1	1

$$\sigma_d \cong \frac{(1 + 1/SNR)\hbar}{2 |\vec{E}| T_m \sqrt{m N_{\text{UCN}}}}$$

Systematic Effects in EDM

- **Nonuniformity of B-field and E-field**
 - Comagnetometer monitors B-field variations
- **Leakage currents from Electric Field**
 - These produce B-fields that change with E-field (must be less than picoAmps)
- **Gravitational offset of n and ^3He ($\sim 10^{-29}$ e-cm)**
- **$\vec{v} \times \vec{E}$ effects are the largest sources of systematic error in present ILL exp.**
 - $\vec{B}_E = \vec{v} \times \vec{E} \rightarrow$ changes $\vec{\mu}$ precession frequency
 - Geometric phase due to \vec{B} gradients

Systematic Controls in new EDM experiment

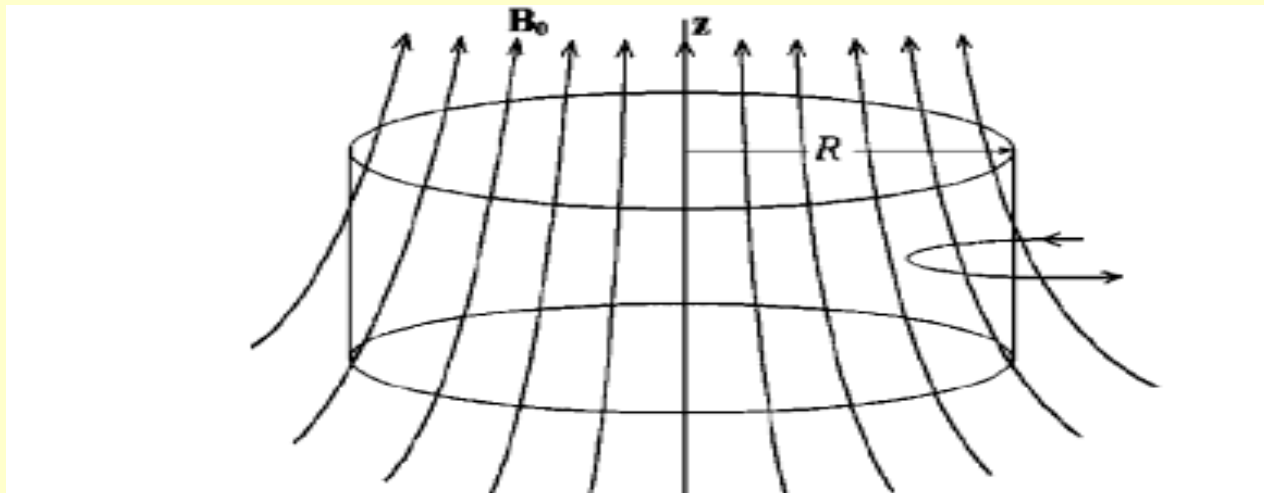
- Highly uniform E and B fields
 - $\cos\theta$ coil in Ferromagnetic shield
 - Kerr effect measurement of E-field
- Two cells with opposite E-field
- Ability to vary influence of B_0 field
 - via dressed spins
- Control of central temperature
 - Can vary ^3He diffusion

Geometric Phase

- **Path-dependent phase (no \hbar)**
 - E.g. Parallel transport of vector on sphere
- **In Quantum Mechanics often called Berry's phase**

False EDM from Geometric phase

- Pendlebury et al PRA 70 032102 (04)
- Lamoreaux and Golub nucl-ex/0407005
- Geometric phase gives false EDM's that depend strongly on radial B fields perpendicular to B_0
 - These result from dB_0/dz



Geometric phase contributions

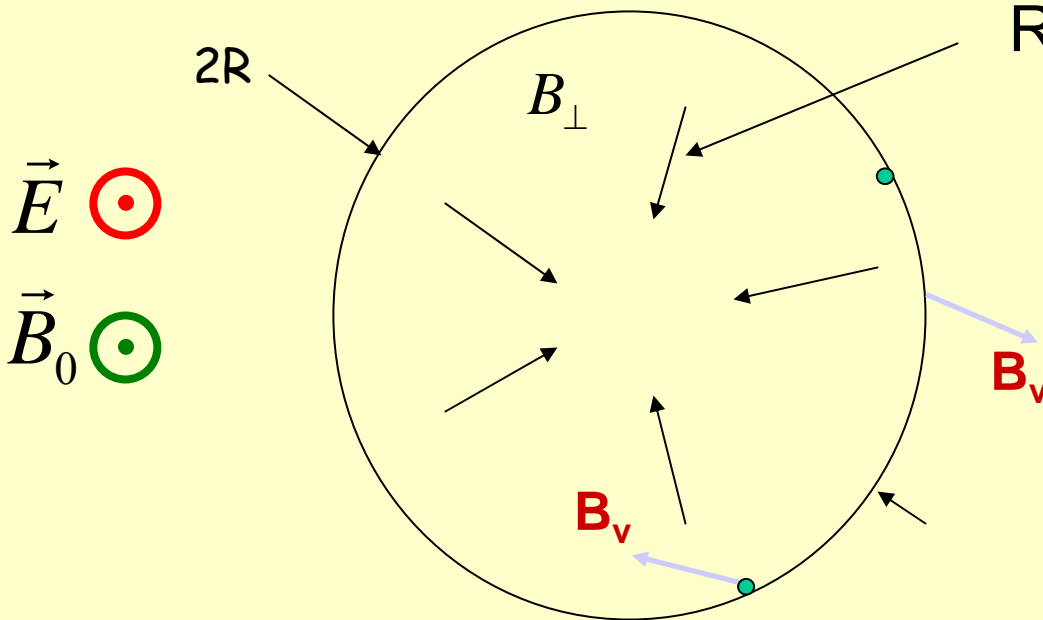
- Gradients and $v \times E$ field gives d^{false}
- Impacts neutron and ^3He
- Magnetometers (^{199}Hg & ^3He) pick up phase at different rate due to higher velocities
 - Can be reduced by frequent collisions with buffer gas or phonons

- if collision rate $\frac{v_{\text{rms}}}{\lambda_{\text{mfp}}} \gg \omega_L$

Then GP doesn't have time to build up

Geometric phase with

$B_E = v \times E$ field



Radial B-field due to gradient

- Motion in B – field shifts the precession frequency - ω_0 :

$$\Delta\omega \cong \frac{\gamma_n^2 \left(B_{\perp} \mp (\vec{v}_n \times \vec{E})/c^2 \right)^2}{4(\omega_0 \mp v_n/R)}$$

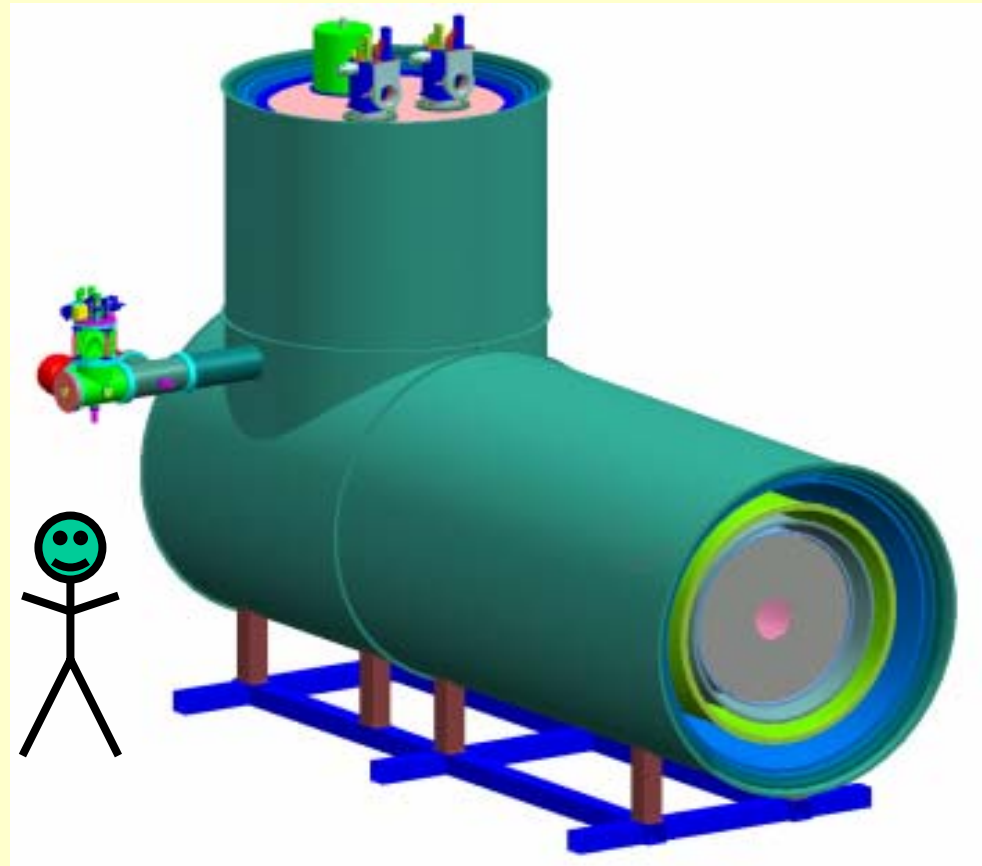
- \mp due to different trajectories
- Does NOT average to 0
- Gives $\Delta\omega$ that depends on direction of $\vec{E} \Rightarrow$ false EDM

$v \times E$ field
changes sign with
neutron direction

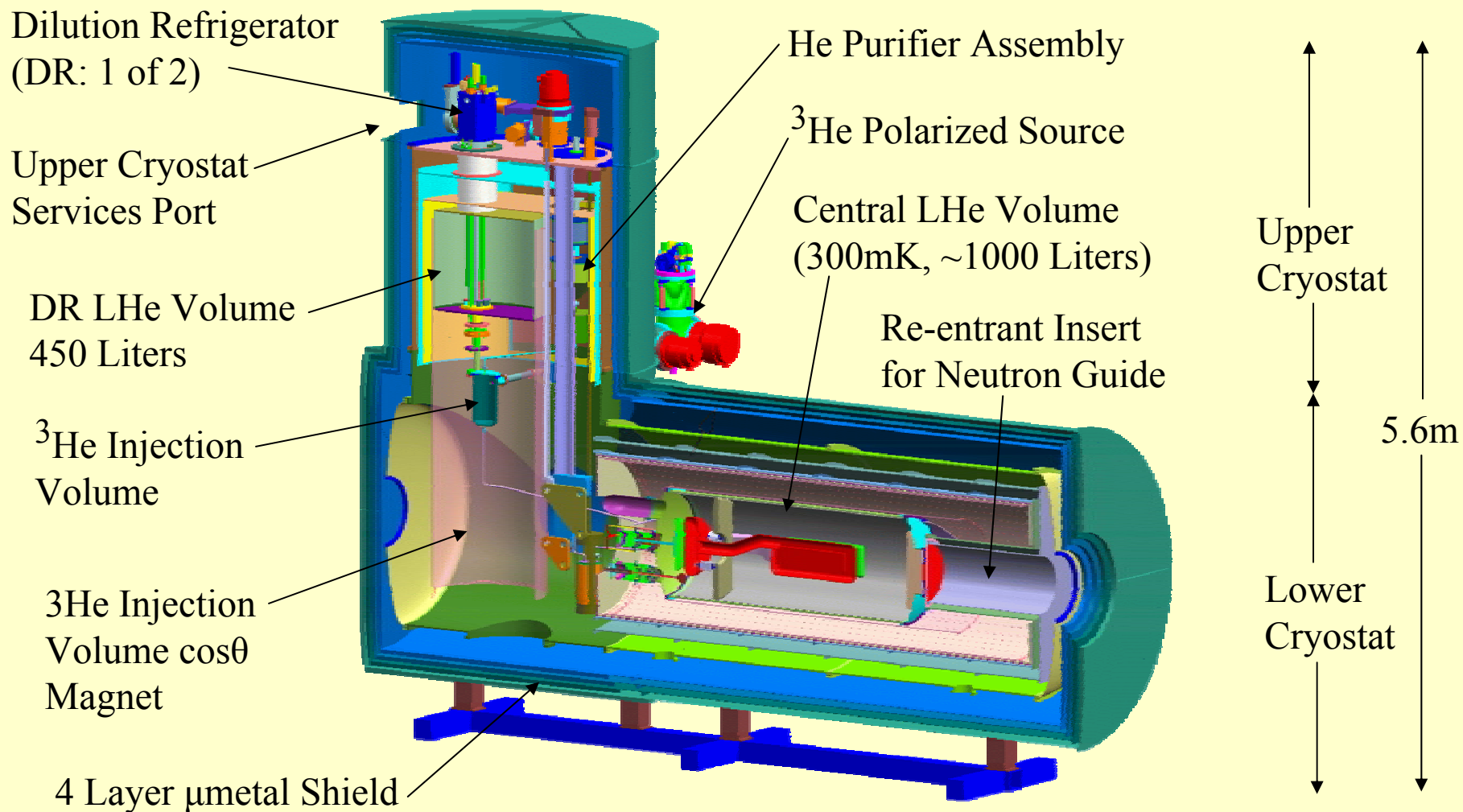
Dressed Spins

- By applying a strong non-resonant RF field, the gyromagnetic ratios can be modified or “dressed”
$$\gamma' = \gamma J_0 (\gamma B_{\text{rf}} / \omega_{\text{rf}})$$
- For a particular value of the dressing field, the neutron and ^3He magnetic moments are equal
- Can tune the dressing parameter until the relative precession is zero. Measure this parameter vs. direction of \vec{E}

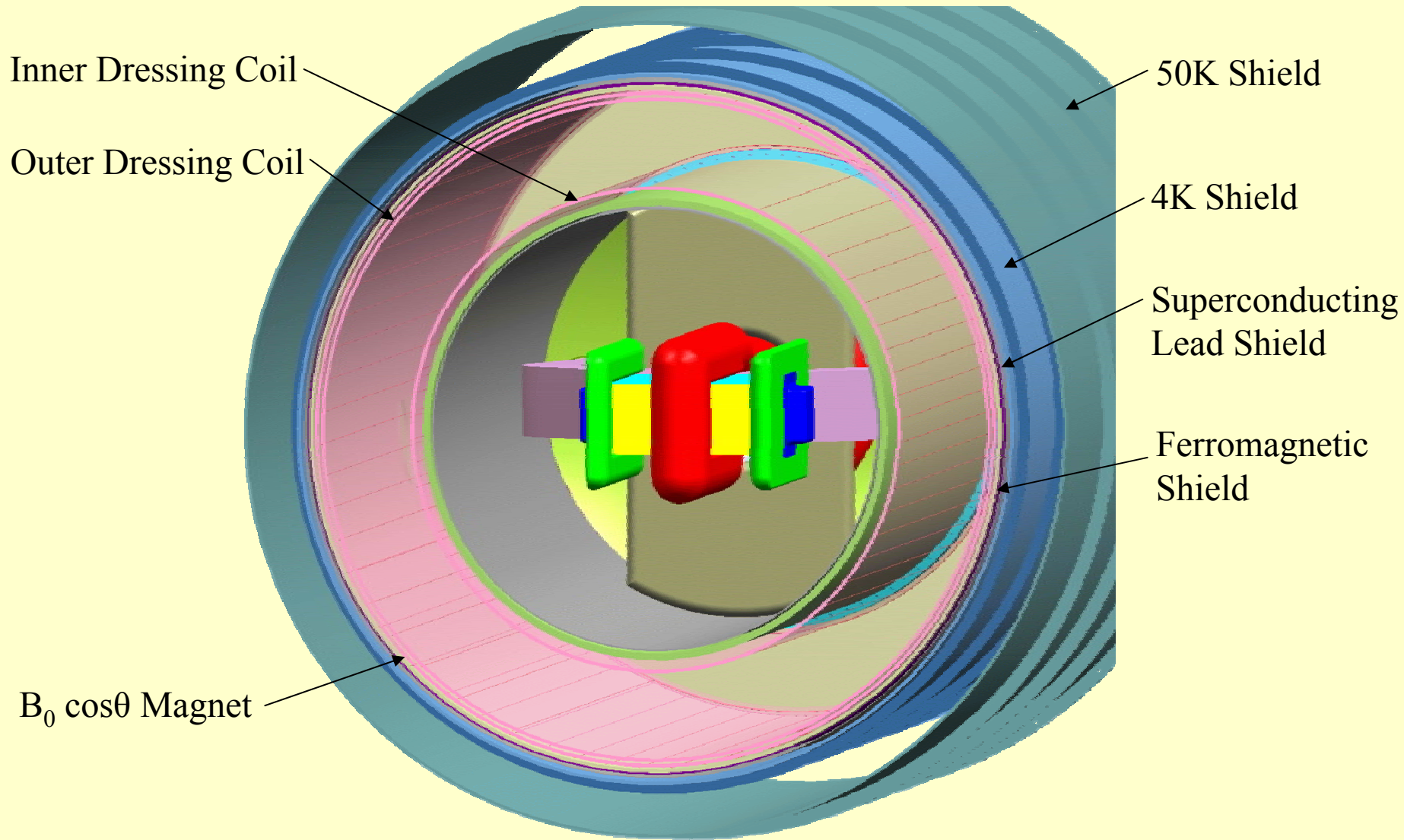
n-EDM Design Concept



EDM Experiment Section View



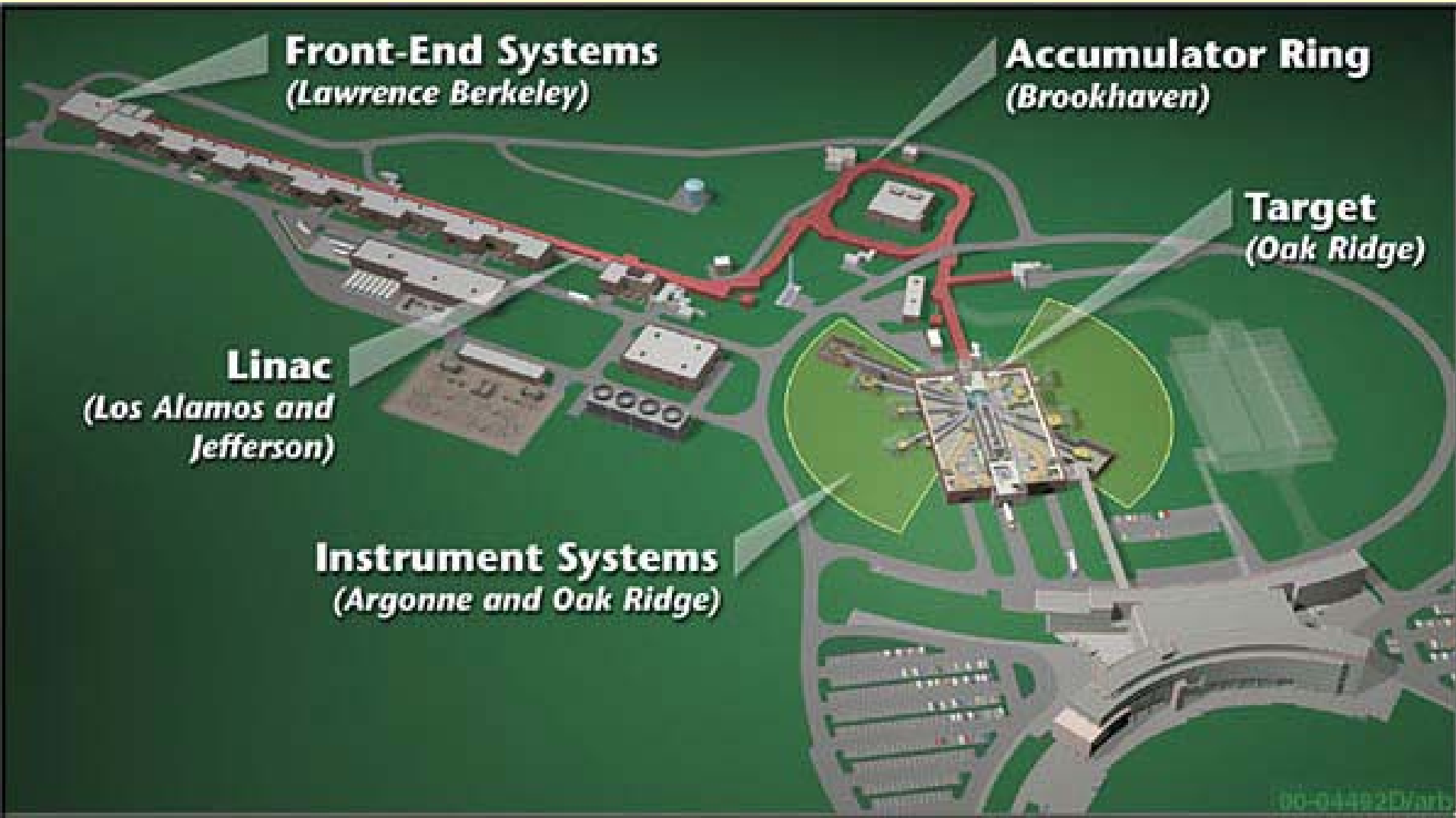
Measurement Cell



Neutrons come from at Oak Ridge National Laboratory

Spallation Neutron Source (SNS) at ORNL

1 GeV proton beam with 1.4 MW on spallation target



Spallation Neutron Source Primary Parameters

Proton beam power on target	1.4 MW
Proton beam kinetic energy on target	1.0 GeV
Average beam current on target	1.4 mA
Pulse repetition rate	60 Hz
Protons per pulse on target	1.5×10^{14} protons
Charge per pulse on target	24 μ C
Energy per pulse on target	24 kJ
Proton pulse length on target	695 ns
Ion type (Front end, Linac, HEBT)	H minus
Average linac macropulse H- current	26 mA
Linac beam macropulse duty factor	6 %
Front end length	7.5 m
Linac length	331 m
HEBT length	170 m
Ring circumference	248 m
RTBT length	150 m
Ion type (Ring, RTBT, Target)	proton
Ring filling time	1.0 ms
Ring revolution frequency	1.058 MHz
Number of injected turns	1060
Ring filling fraction	68 %
Ring extraction beam gap	250 ns
Maximum uncontrolled beam loss	1 W/m
Target material	Hg
Number of ambient / cold moderators	1/3
Number of neutron beam shutters	18
Initial number of instruments	5

SNS Status



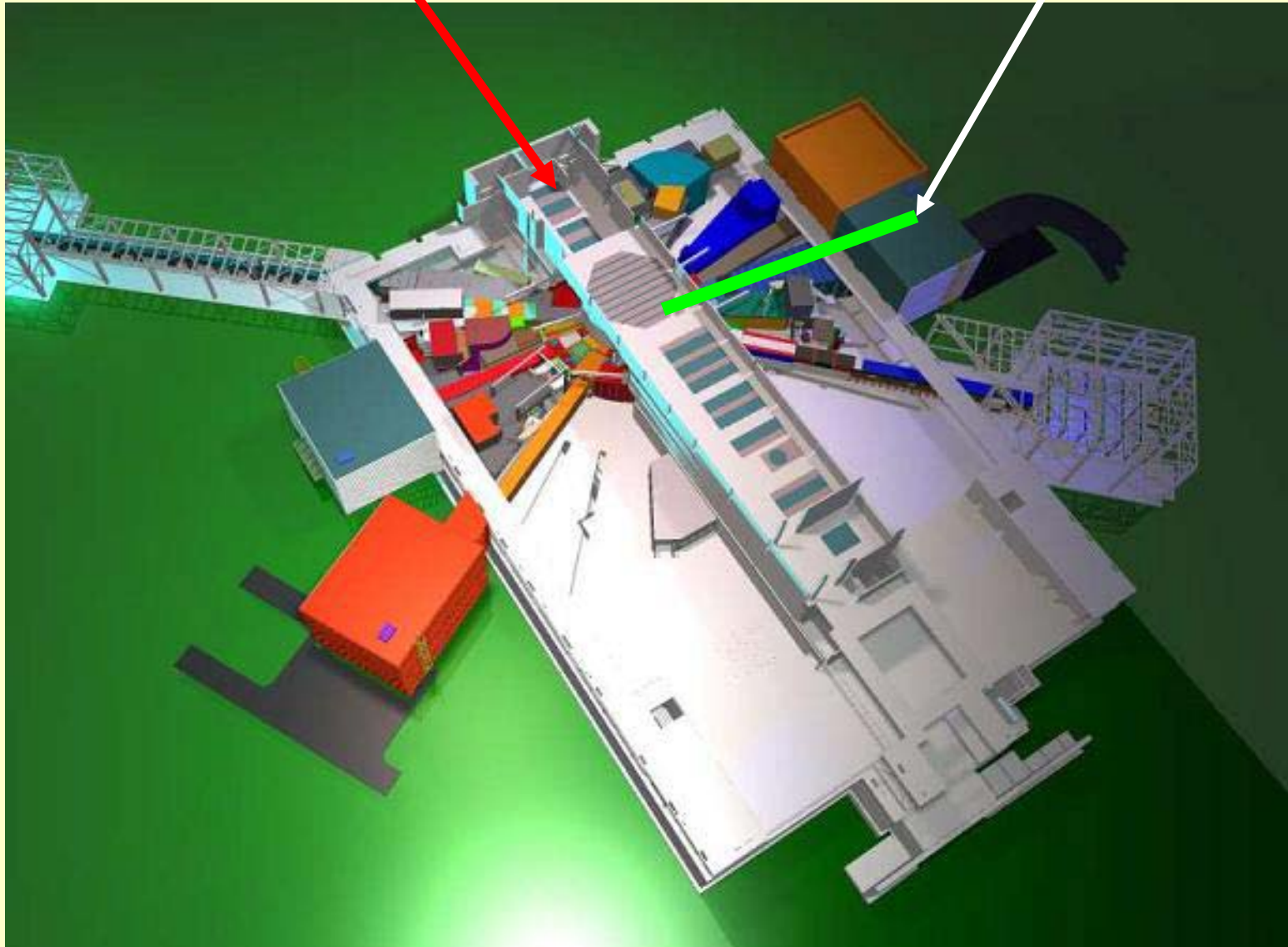
Photo courtesy of ORNL

- **SNS completed:** 2006
- **FNPB Beam line completed:** 2007
- **Full design flux:** 2008
- **SNS Total Project Cost:** 1.411B\$

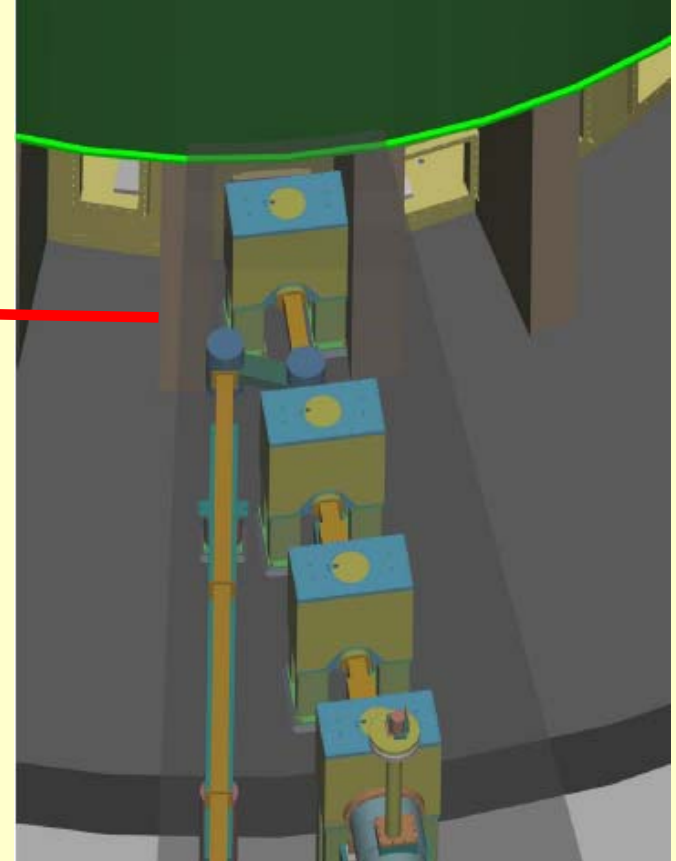
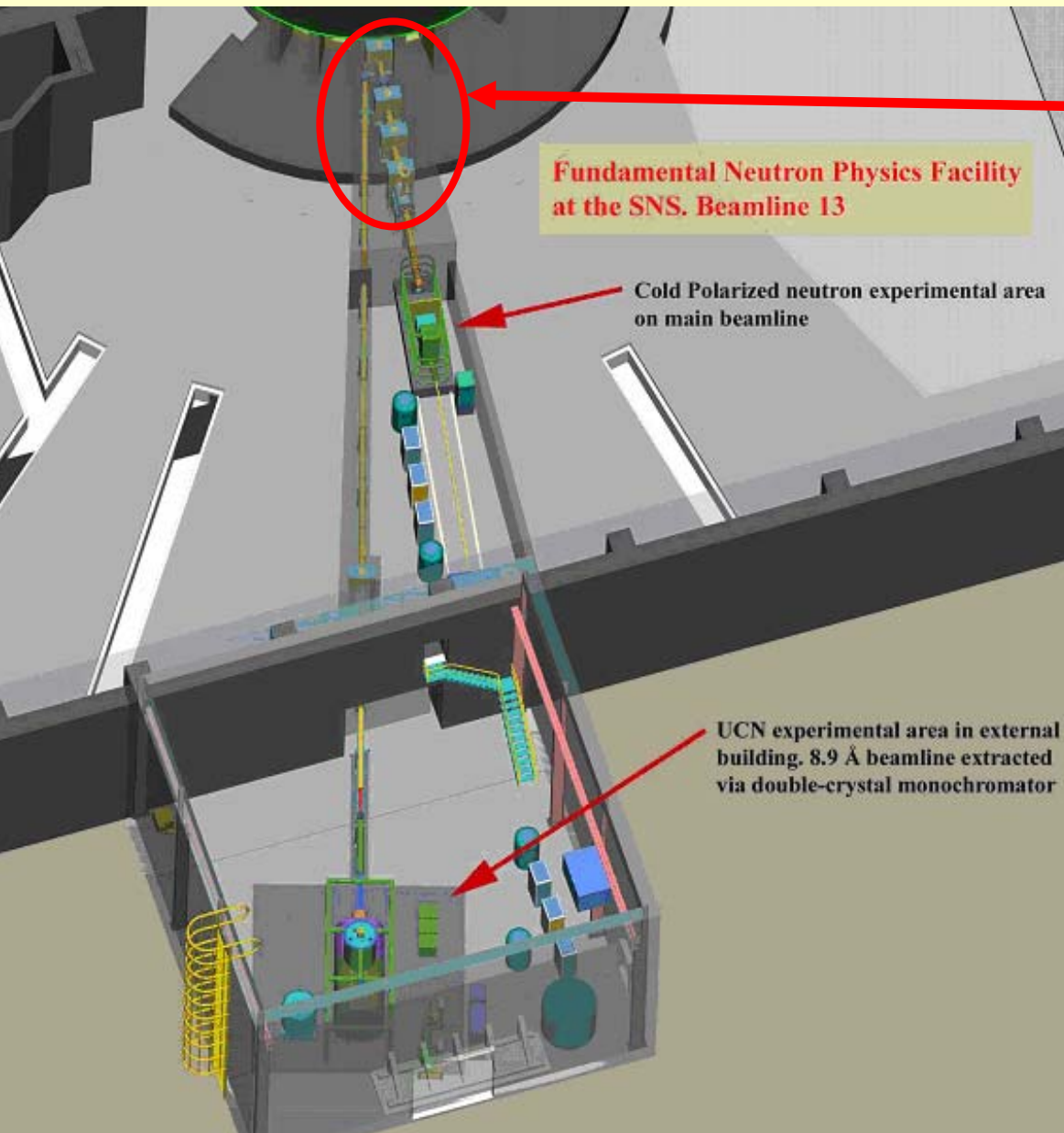
SNS Target Hall

18 neutron beam ports with 1 for Nuclear Physics

p beam

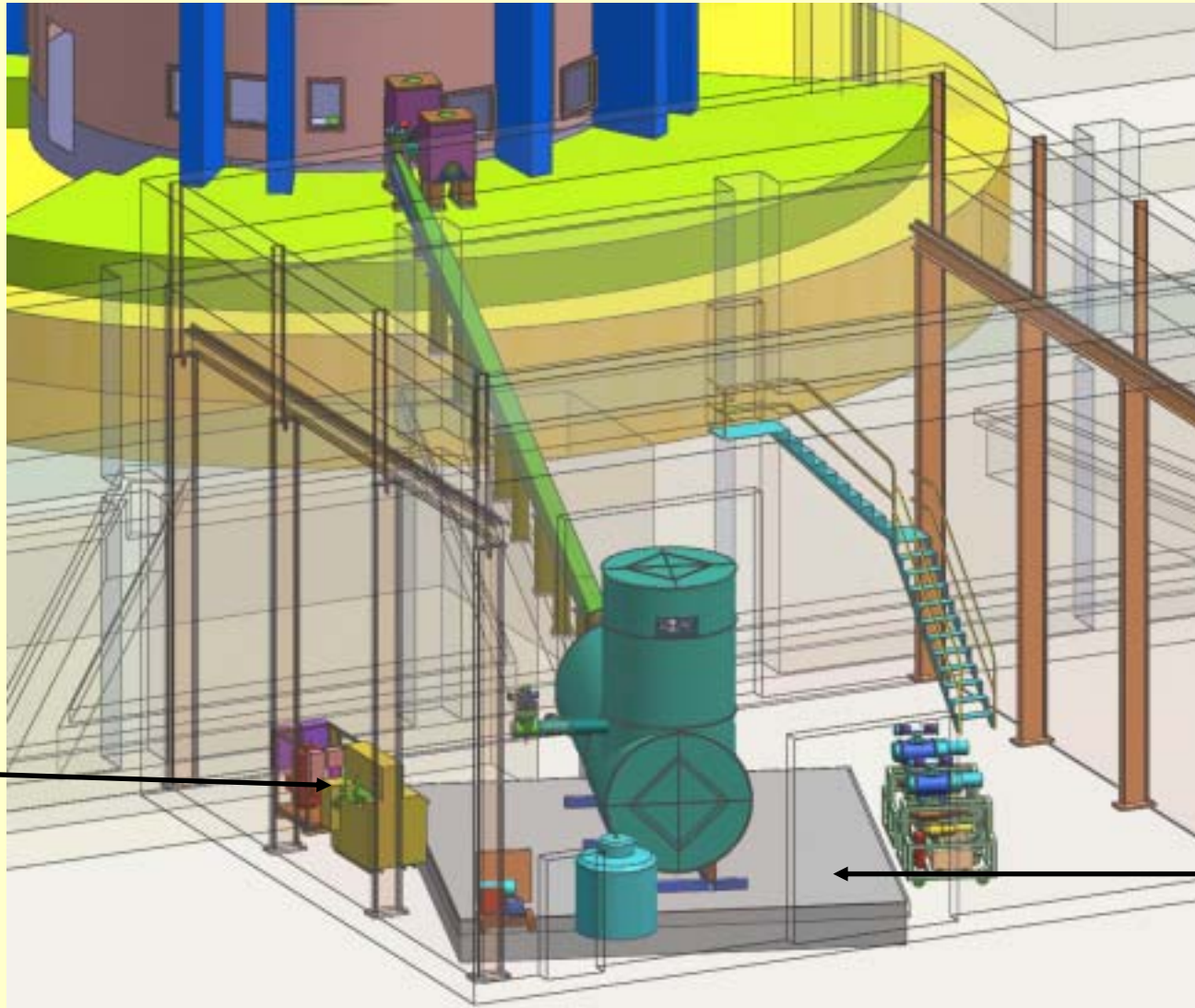


Fundamental Neutron Physics Beamline



Double monochromator
selects 8.9 Å neutrons

EDM Experiment at SNS



He
Liquifier

Isolated
floor

Funding Status of US n-EDM

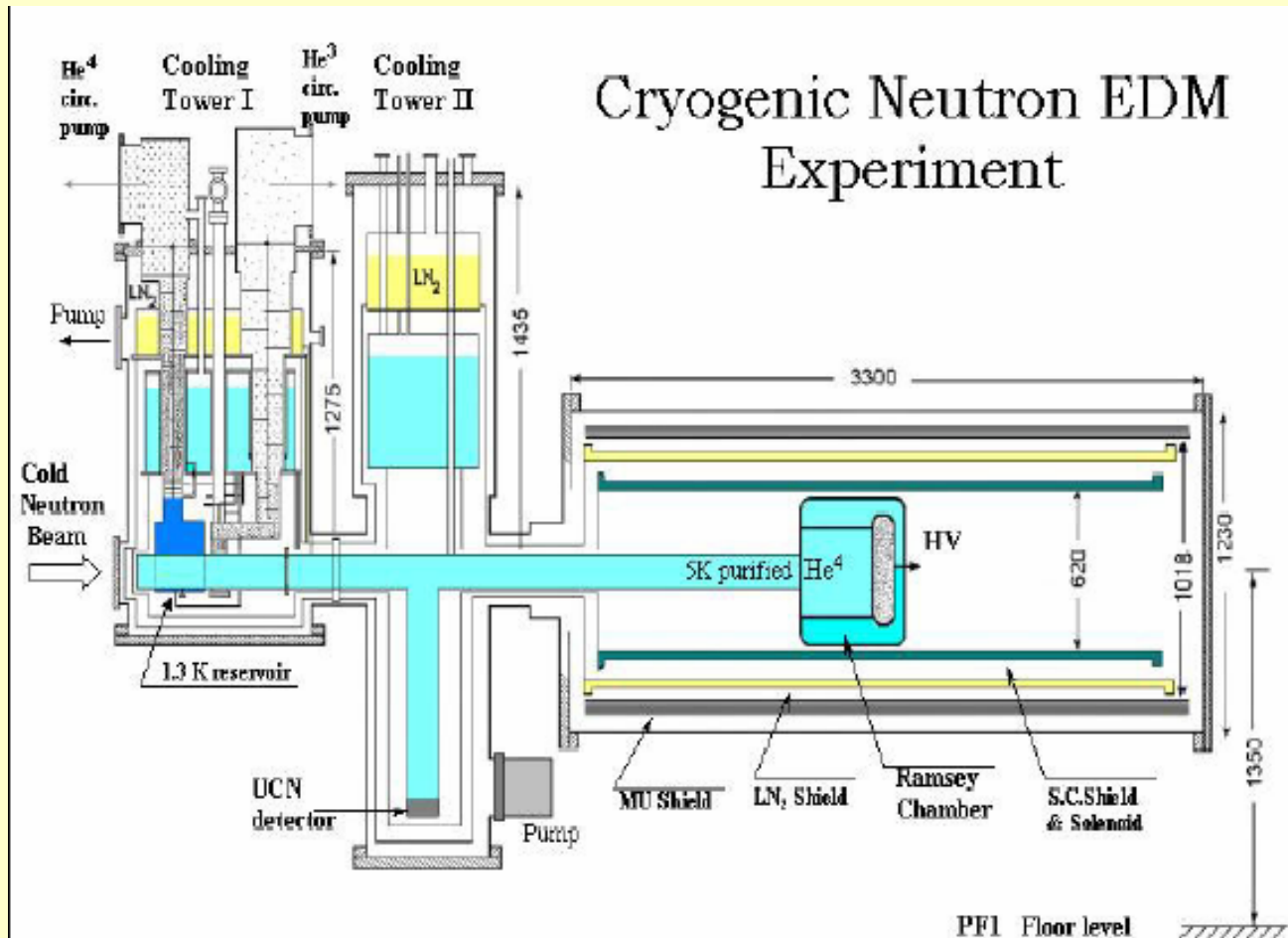
- 2002: Preproposal submitted to DOE: 2003
DOE/NSF subcommittee review Department of Energy
 - Established DOE "mission need"
- 2003: Requested R&D funding for FY03-04
 - Received 650k\$ - resolved key feasibility issues
- 2004: Requested R&D funds for FY05-06
 - Request = 1.65M\$ - "... we'll get back to you"
- 2005: "Internal" Cost and Schedule review
 - Total Cost of experiment = 16+/-2 M\$ (~5M\$ NSF)
 - Schedule of experiment: commissioning in 2012
- 11/05: n-EDM receives "Critical Decision 0" (CDO)

Other New EDM experiments

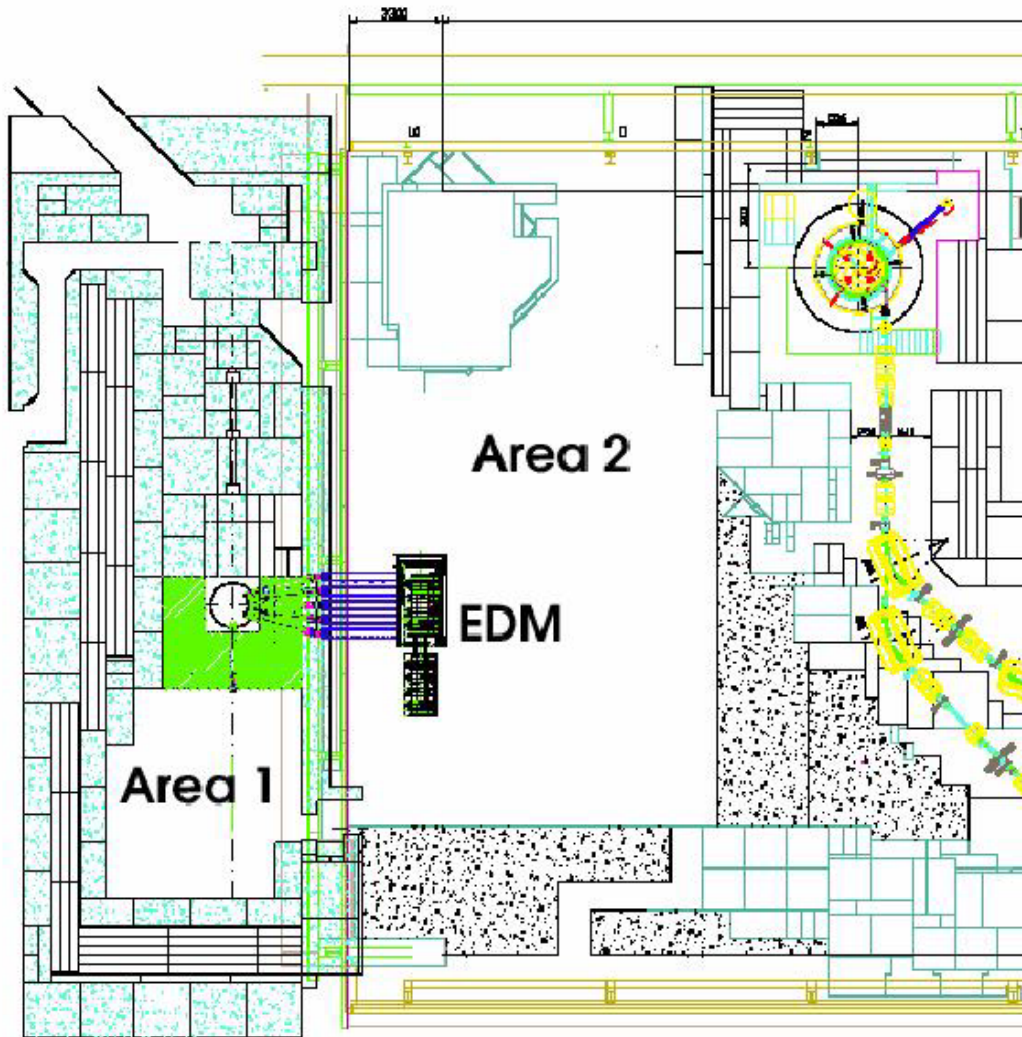
- **CryoEDM at ILL**
 - Similar to new US experiment
 - Does not use polarized ^3He
 - Sensitivity of 2×10^{-28} e-cm
 - First phase of experiment underway
- **EDM at PSI based on SD_2**
 - Similar technique to present ILL exp.
 - But no comagnetometer
 - Sensitivity of 7×10^{-28} e-cm
 - SD_2 source being constructed

New ILL EDM Experiment

- Similar to new US experiment
- No ^3He (only SQUIDS)



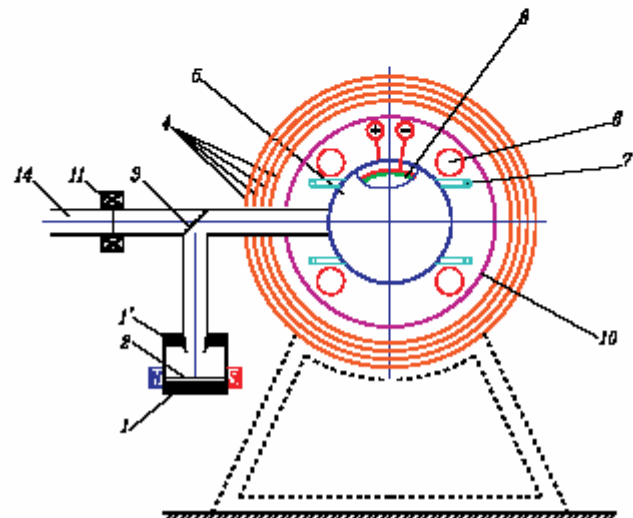
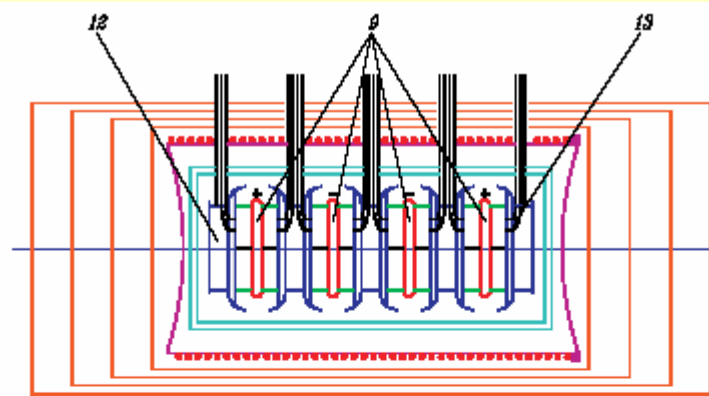
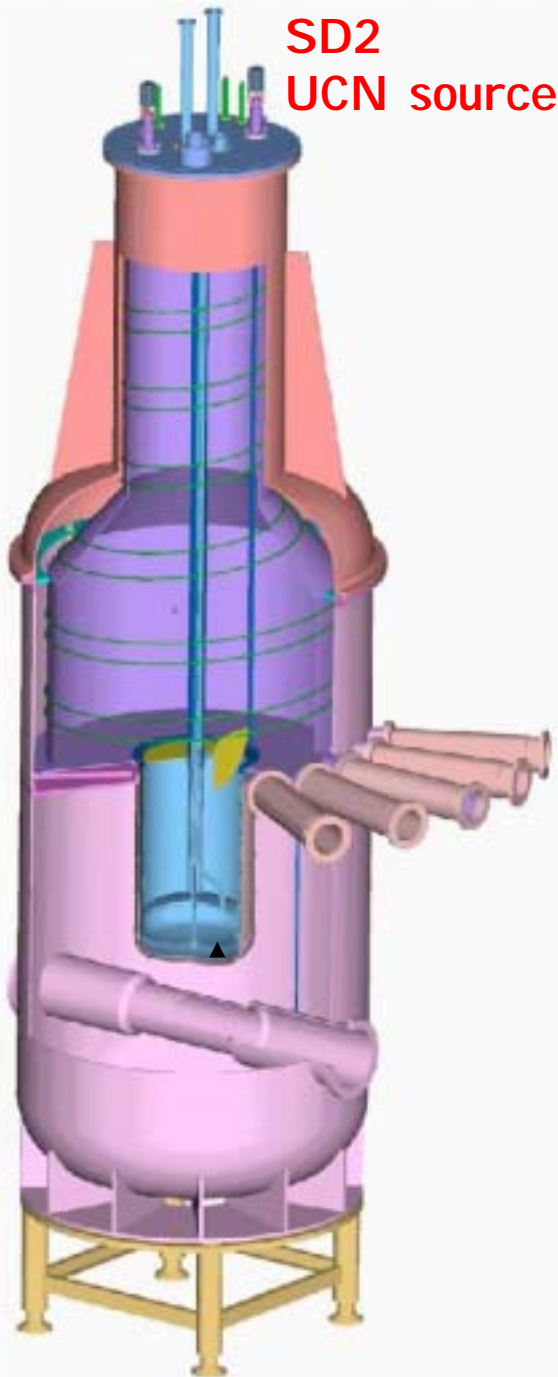
New EDM experiment at PSI



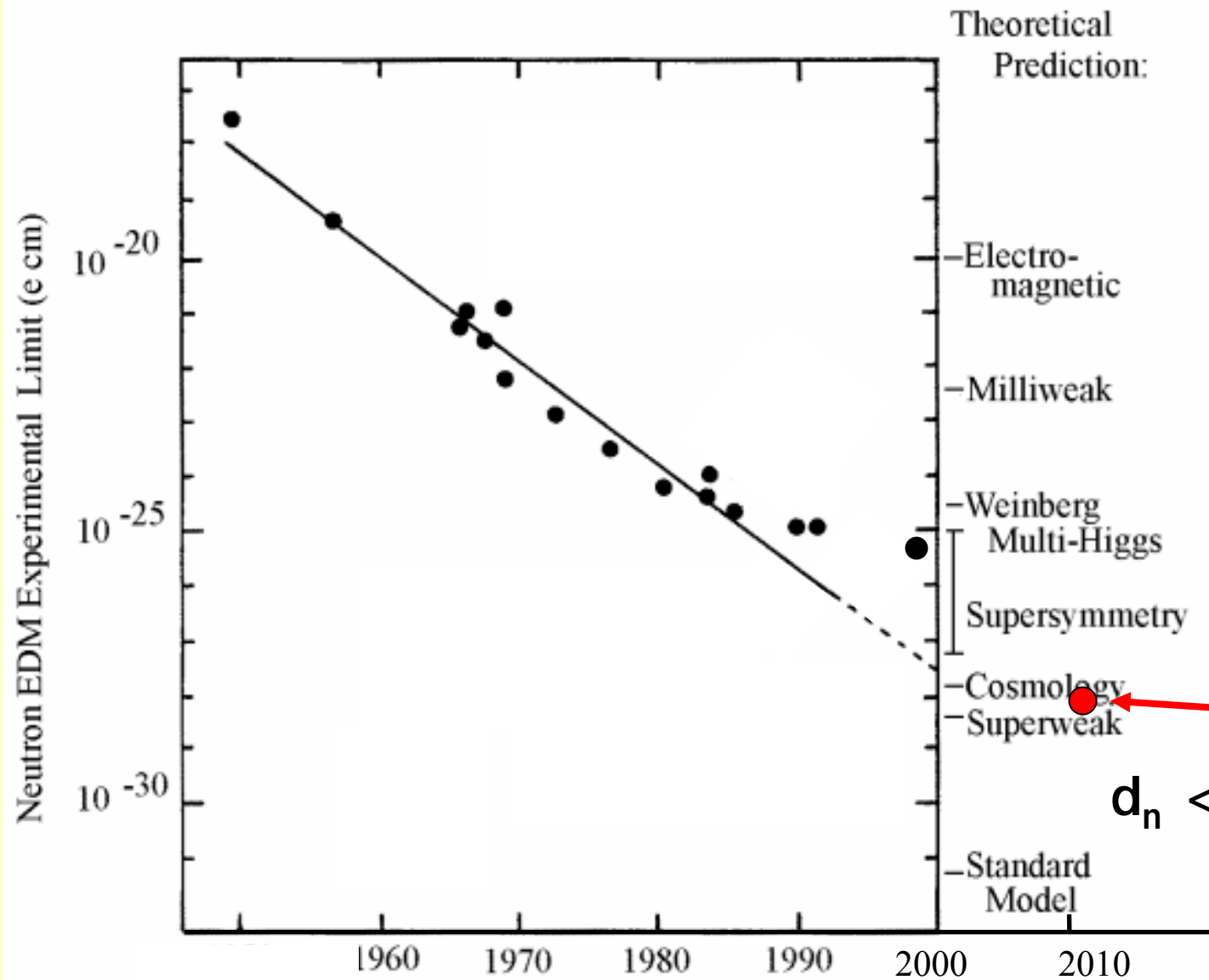
PSI UCN Facility
(using Solid D_2
from Los Alamos-
Caltech-...
collaboration)

PSI EDM

- Based on present ILL experiment
- No "Comagnetometer"
- Uses UCN in adjacent cell with $|\vec{E}| = 0$



New n-EDM Sensitivity



EDM @ SNS

$$d_n < 1 \times 10^{-28} \text{ e-cm}$$

Summary

- Thanks for Invite!
- Searches for physics beyond the standard model are key part of Nuclear Physics
- Precision measurements at low energy can access very high energy physics
- Physics reach of EDM measurements remains strong (even after Large Hadron Collider)
 - New sources of CP violation possible in SUSY